

Méthodes intégrales pour l'évaluation en champ proche

Camille Carvalho

ENSTA(2009-2012)

32 Blvd Victor, Paris



MA102
MA201
MA206
MAE21
C7-4



Doctorat (2012-2015)

921 Blvd des Maréchaux, Palaiseau



La famille "changement de signe"

(K. Ramdani, C.M. Zwölf, L. Chesnel, C.
Carvalho, M. Rihani, F. Chaaban, ...)

JO des Poètes 2024 - Relais 4x60



The close evaluation problem

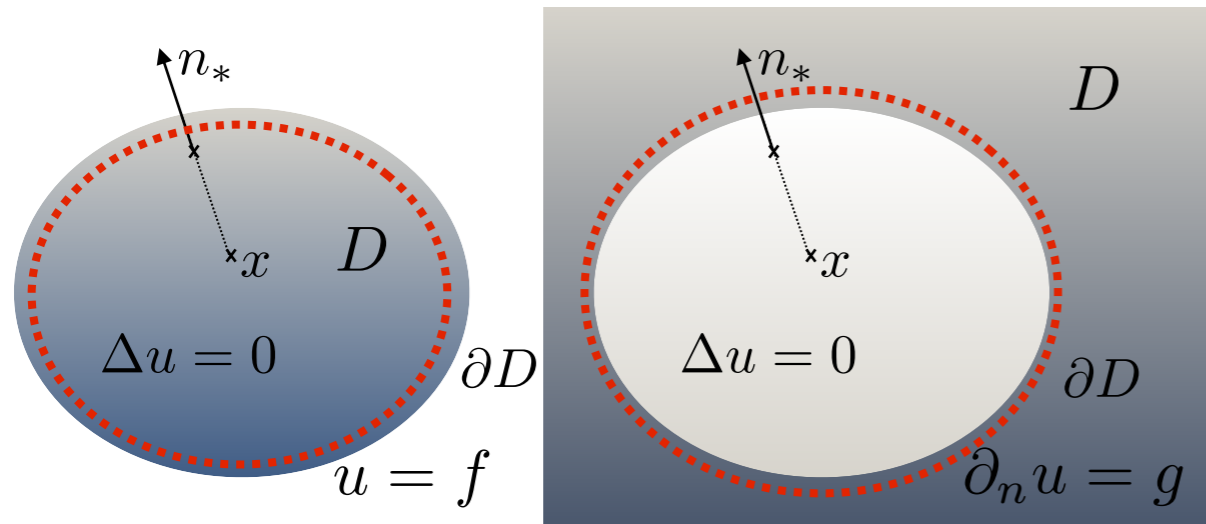
Goal: accurately evaluate the solution **near the boundary**.

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Examples:

Laplace problems

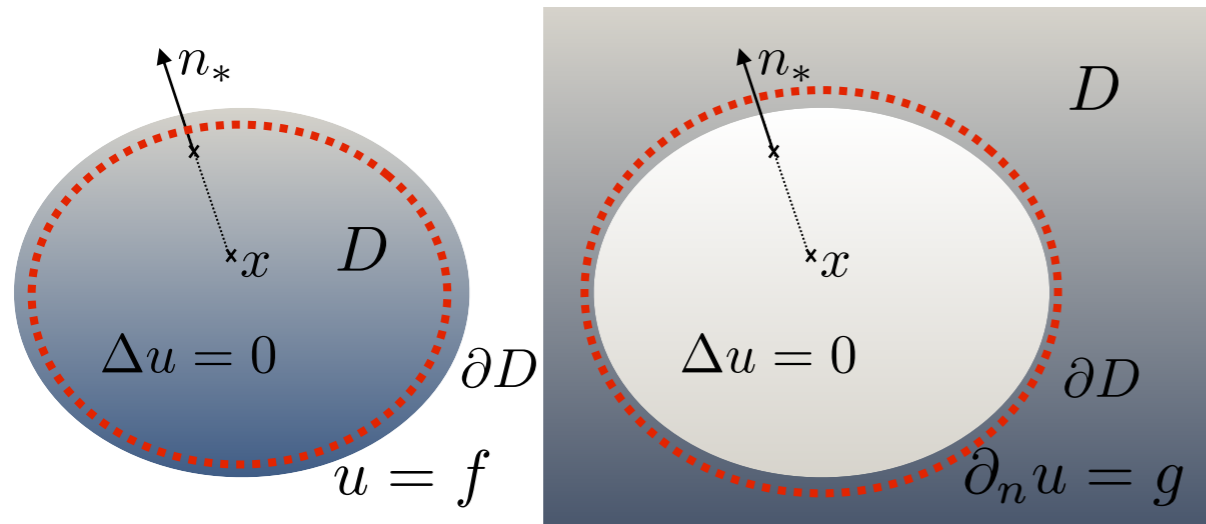


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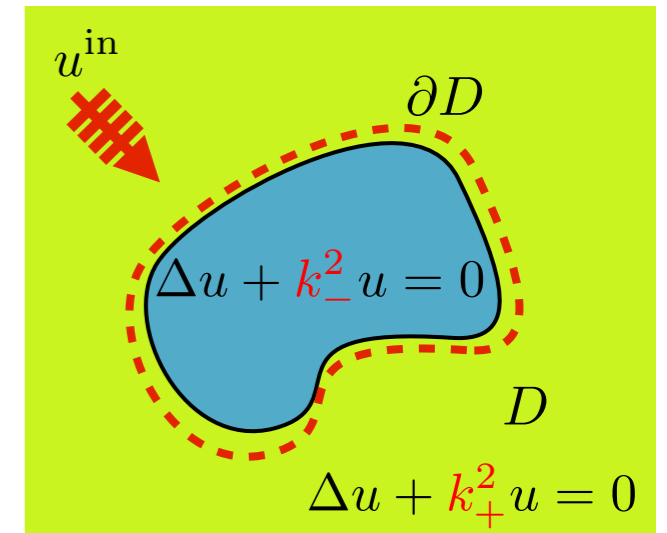
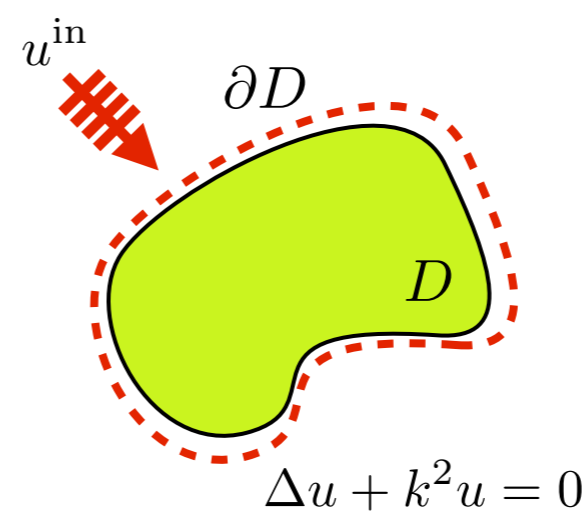
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Scattering problems

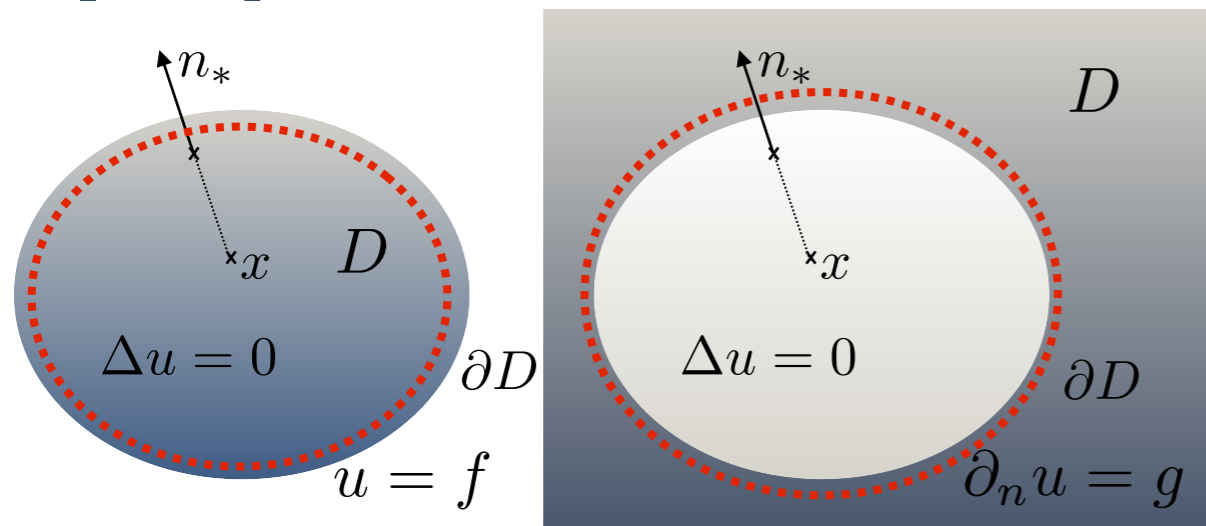


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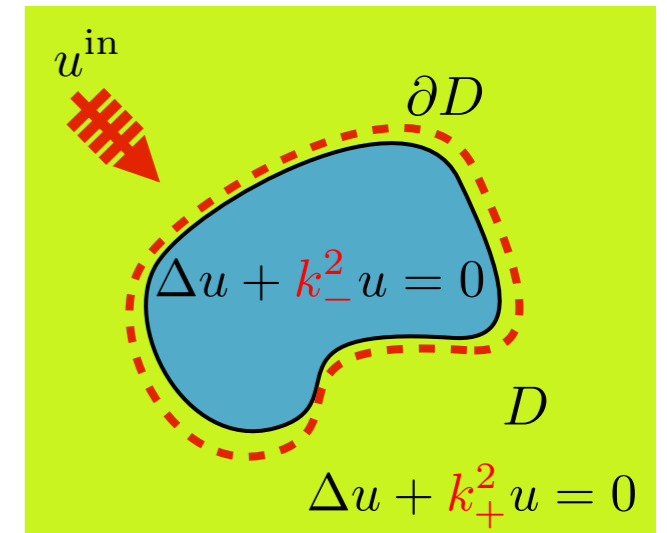
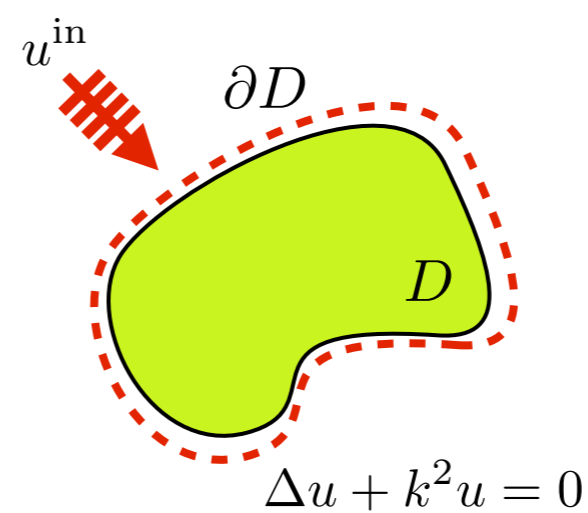
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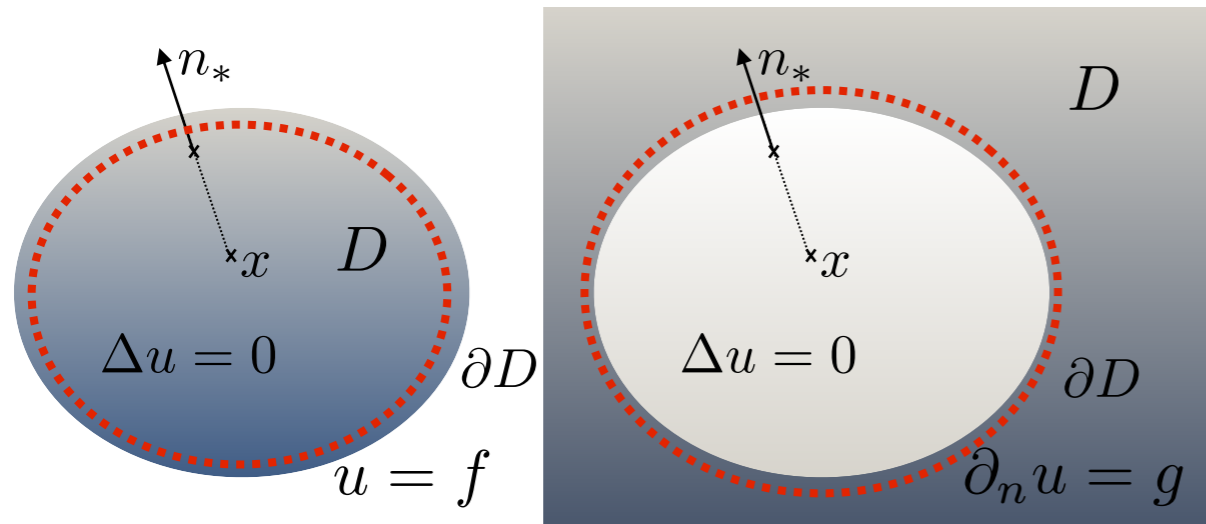
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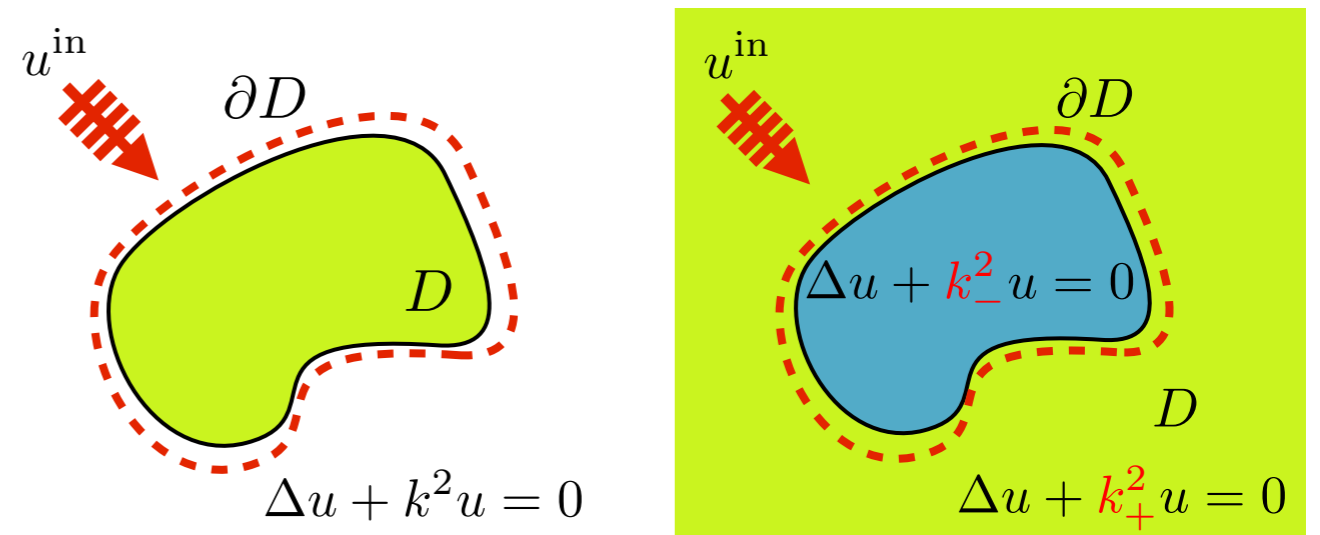
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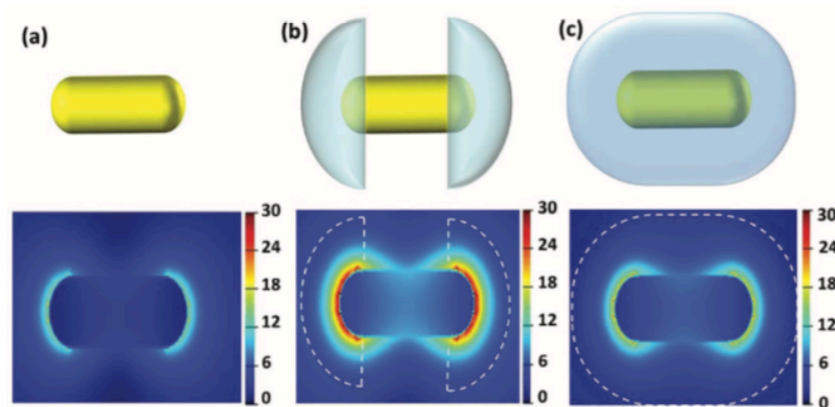
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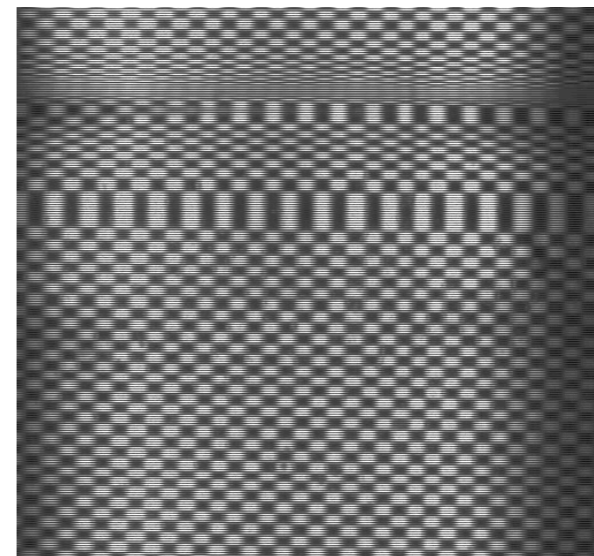
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Zhu et al., Nanoscale (2020)



JFL Lab.

An example in 2D

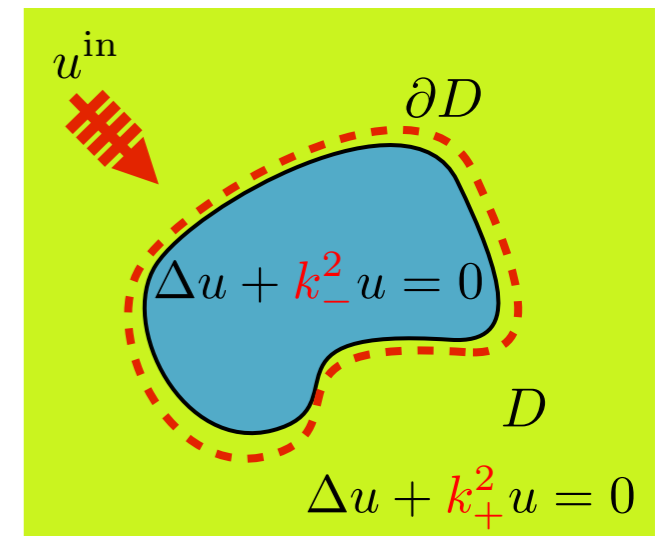
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+ Radiation condition at ∞



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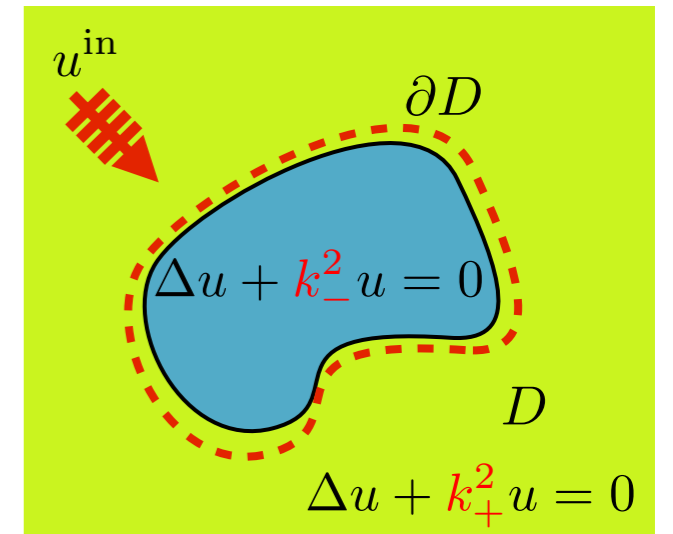
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Boundary integral methods represents the solution via layer potentials:

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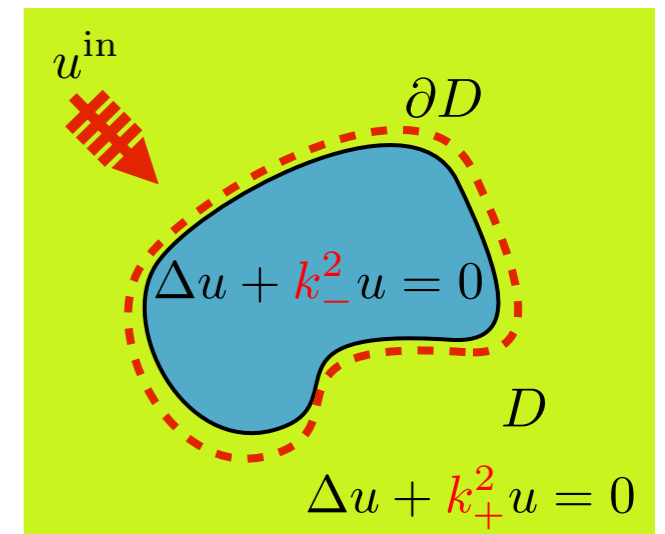
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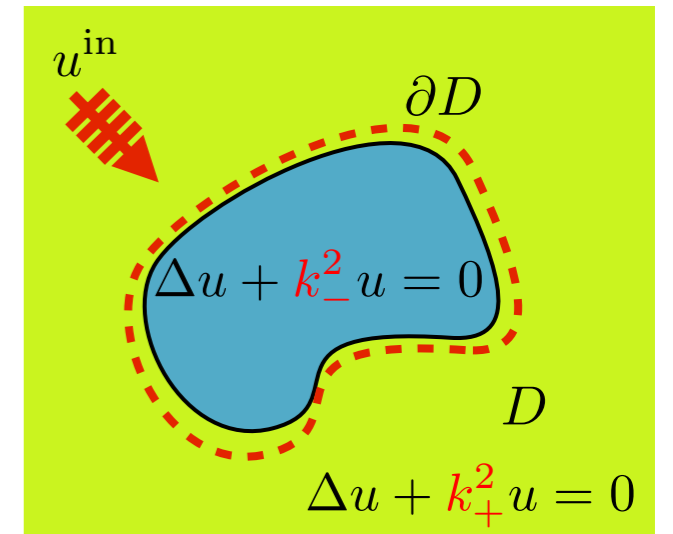
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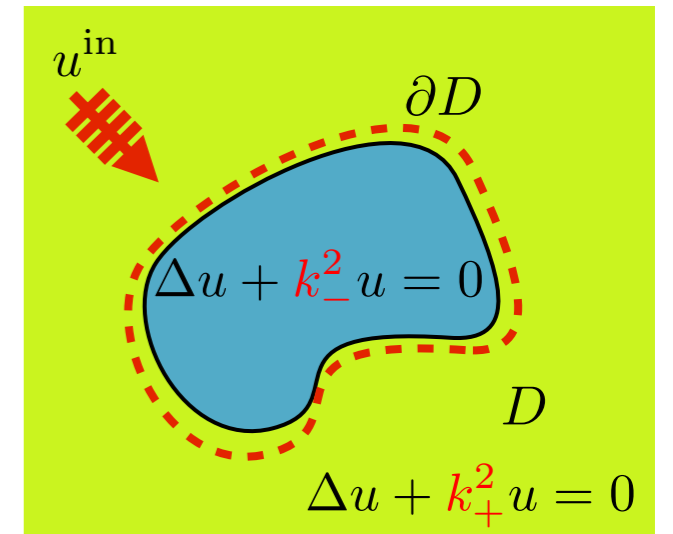
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$$G^-(x, y) := \frac{i}{4} H_0^{(1)}(k_- |x - y|)$$

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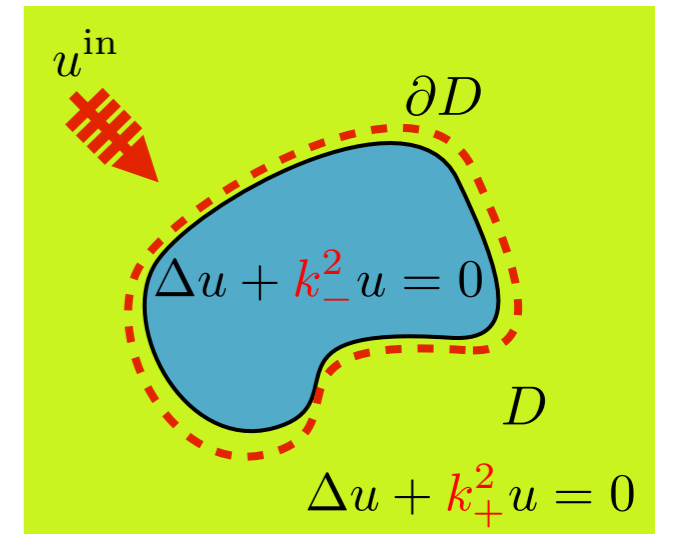
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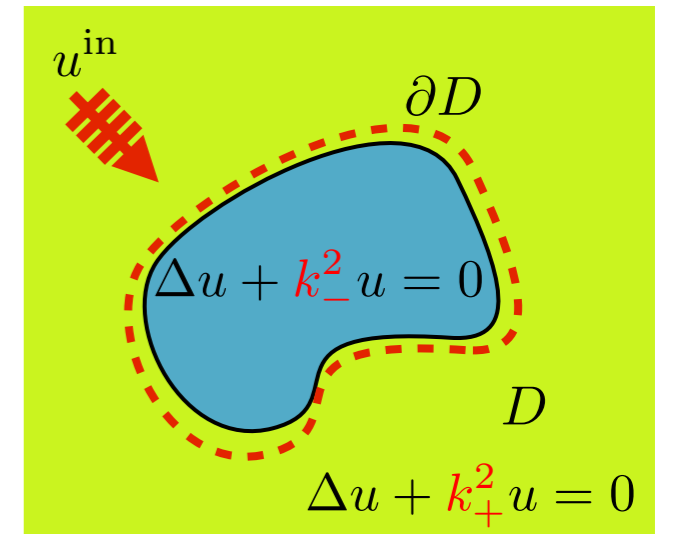
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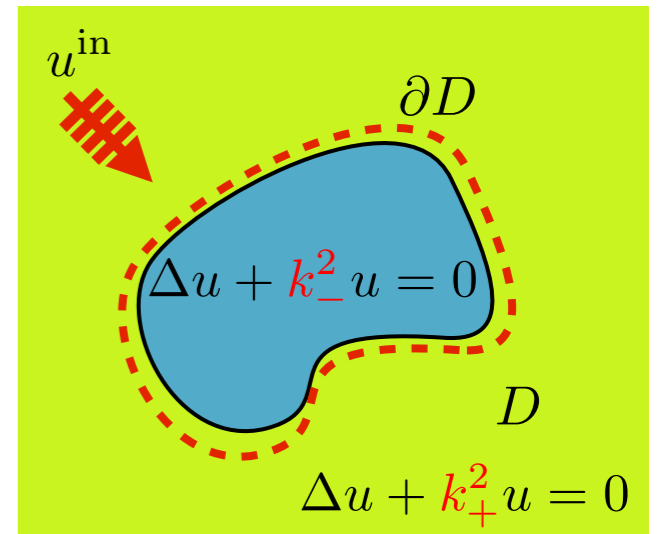
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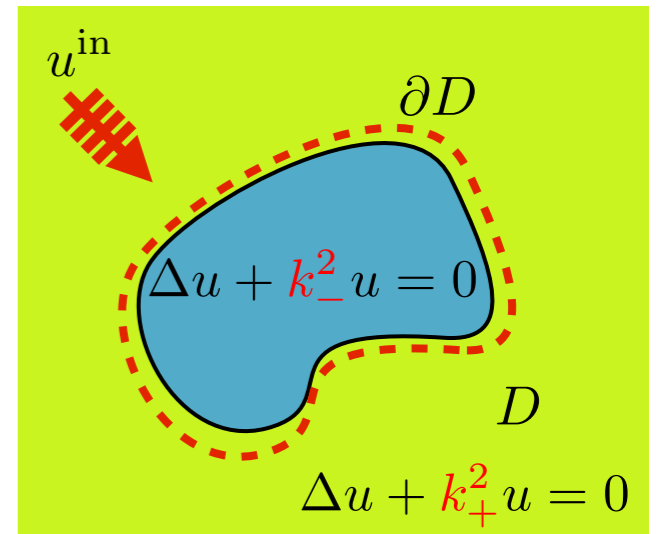
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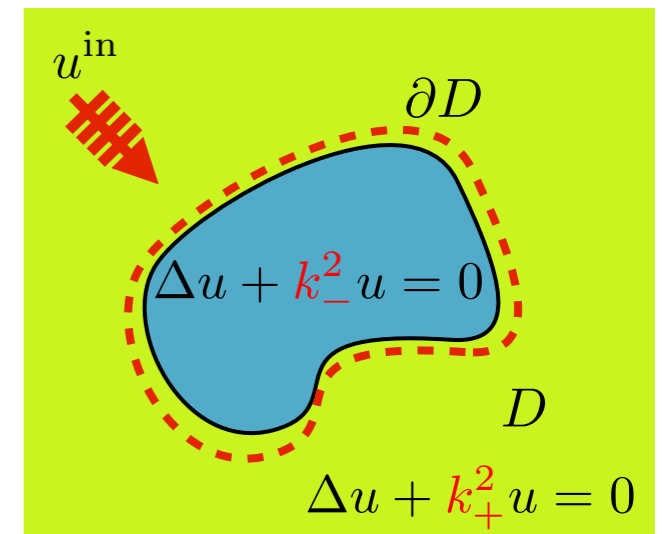
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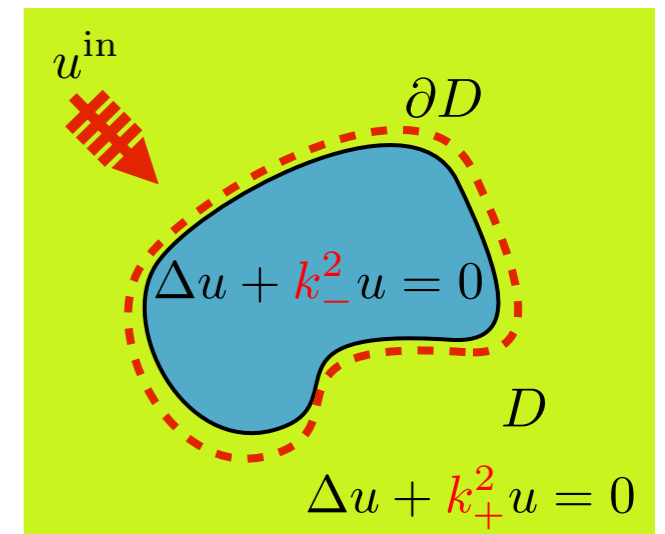
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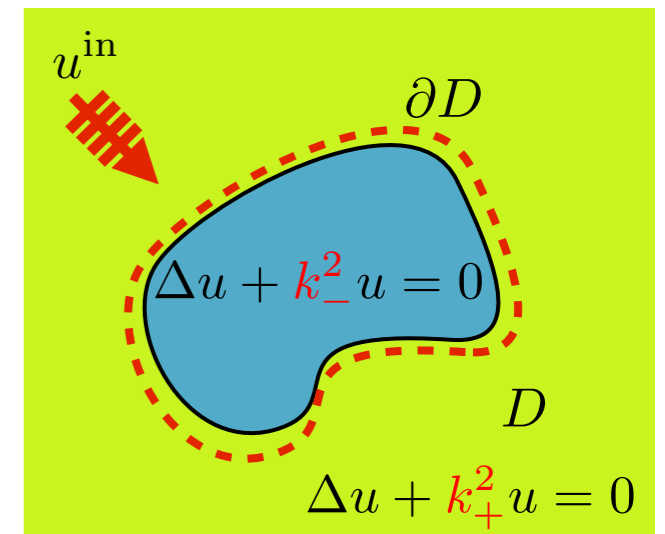
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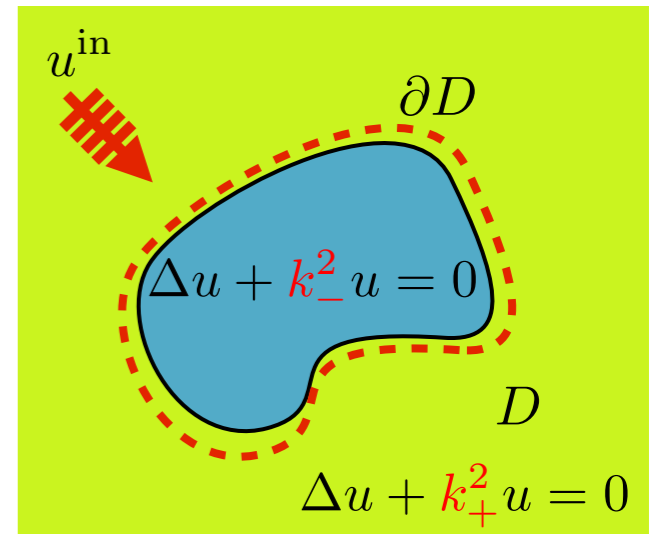
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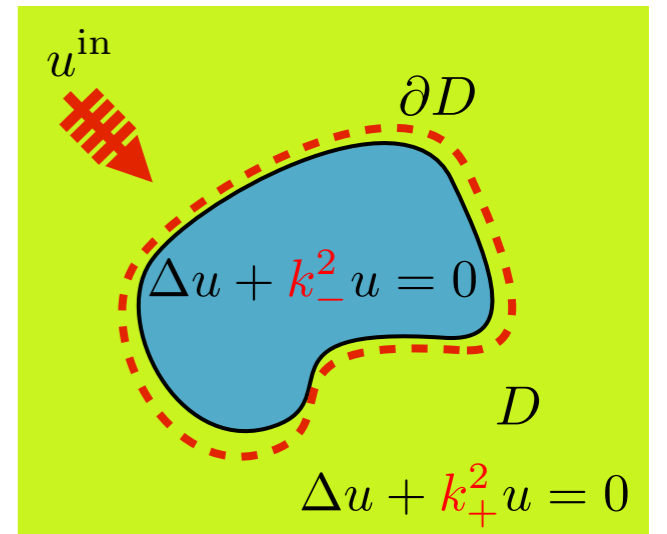
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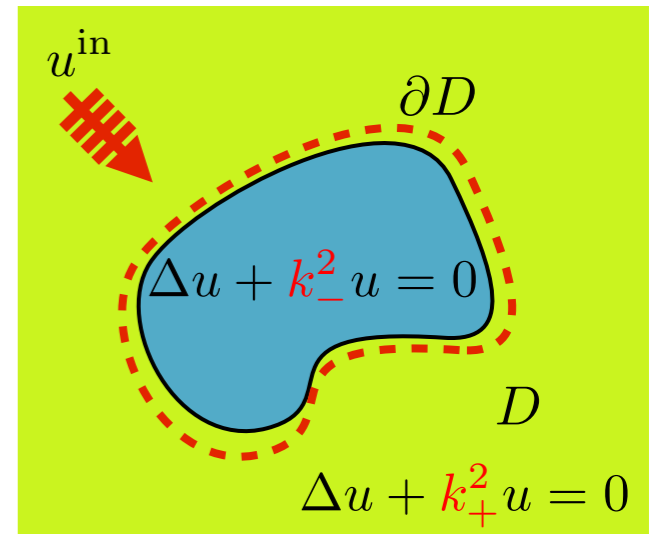
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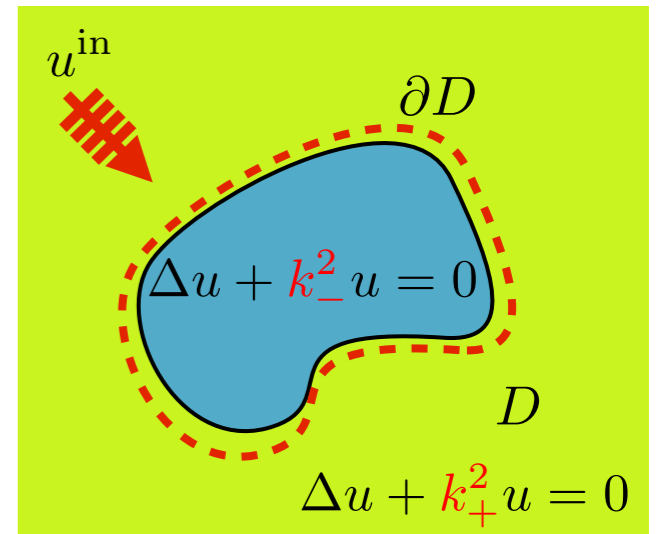
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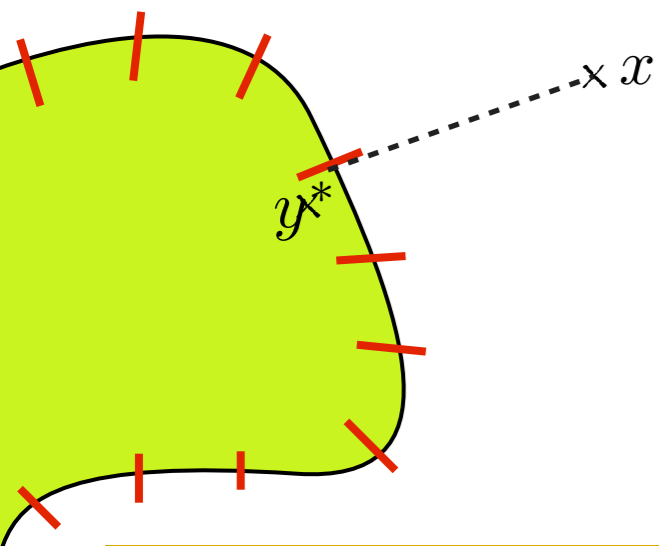
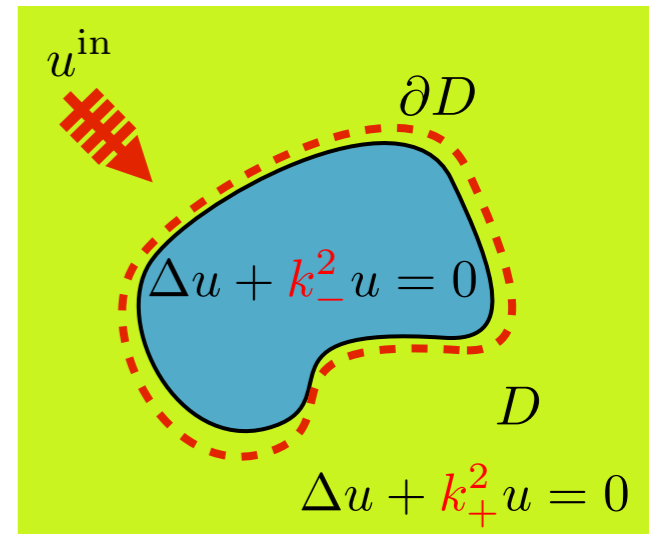
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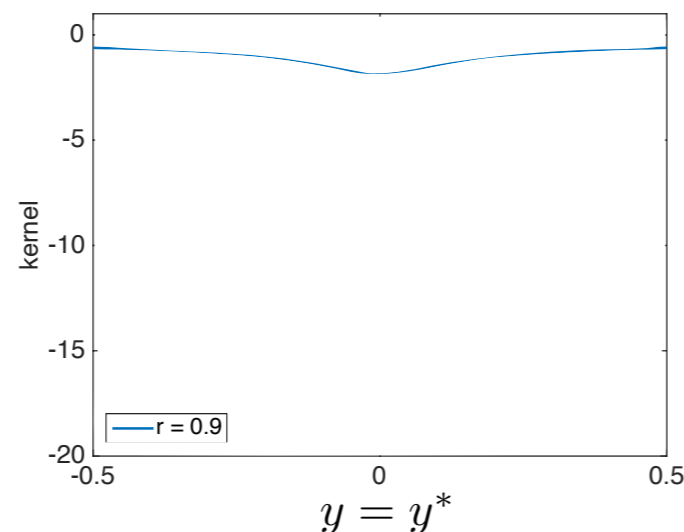
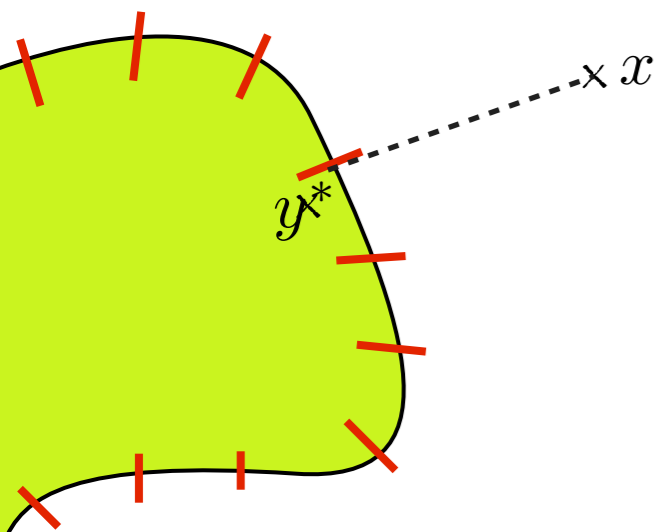
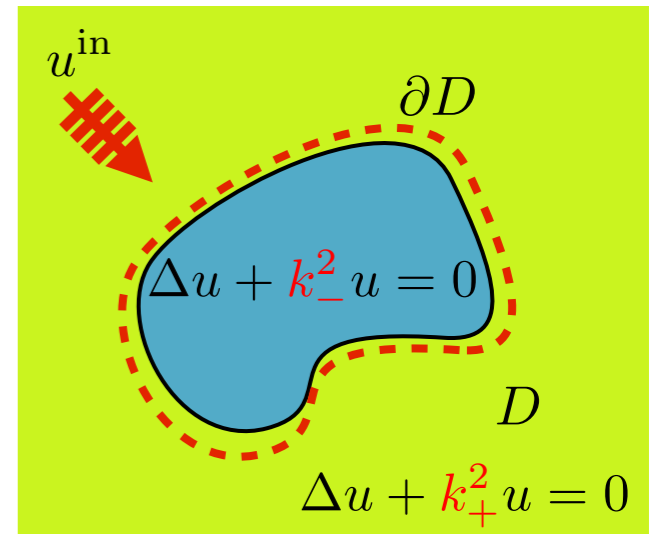
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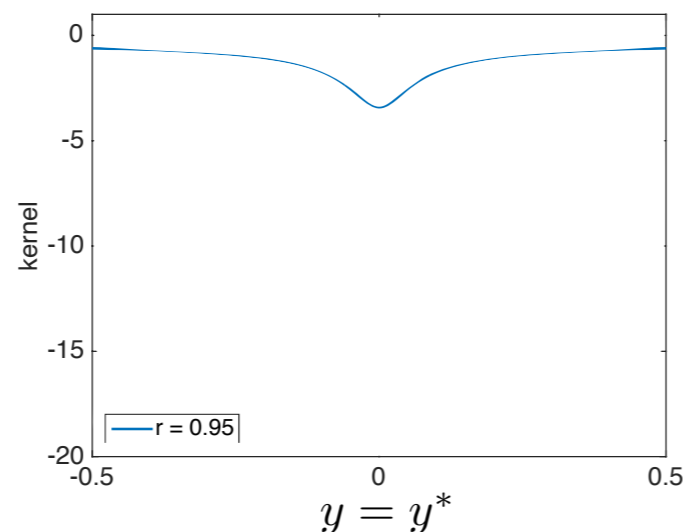
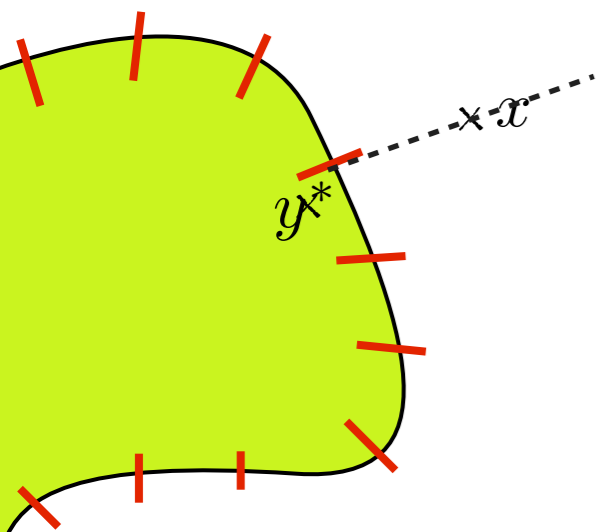
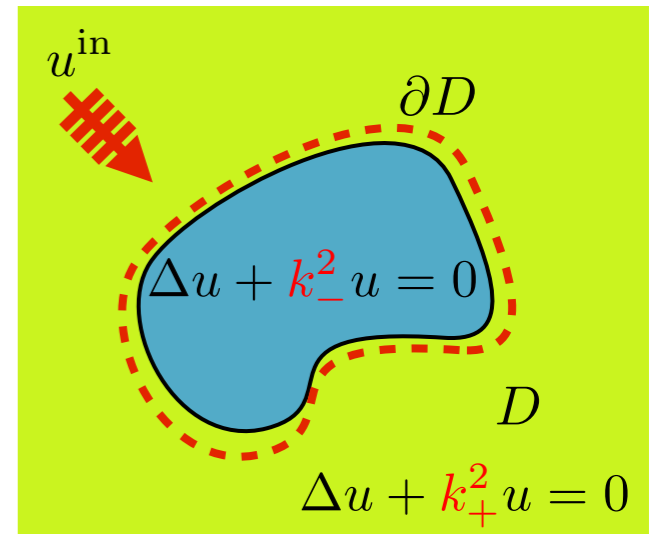
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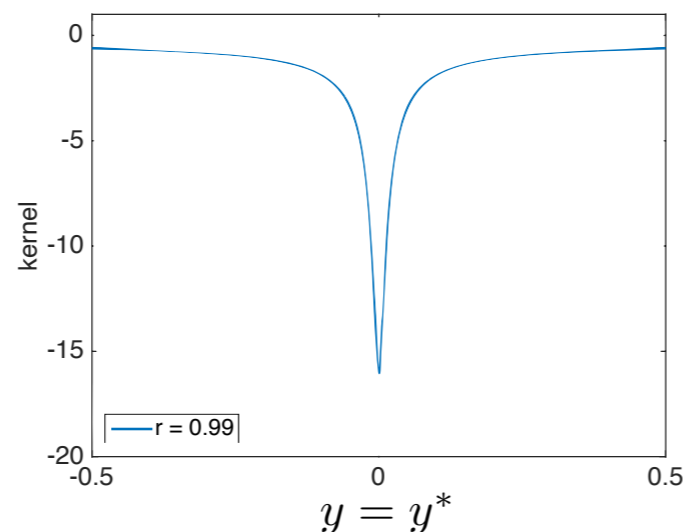
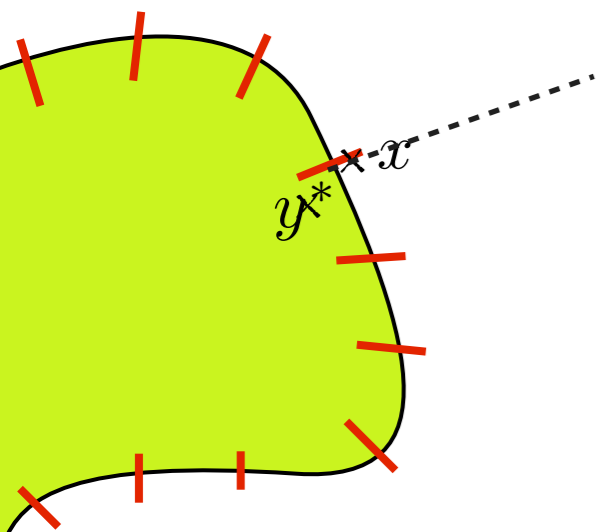
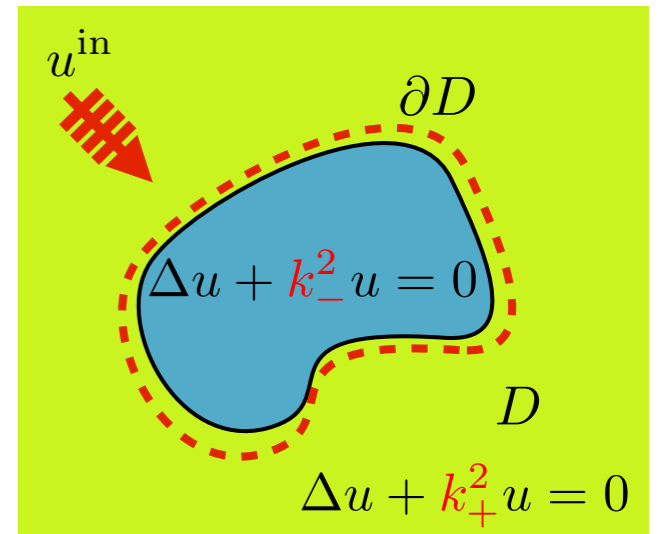
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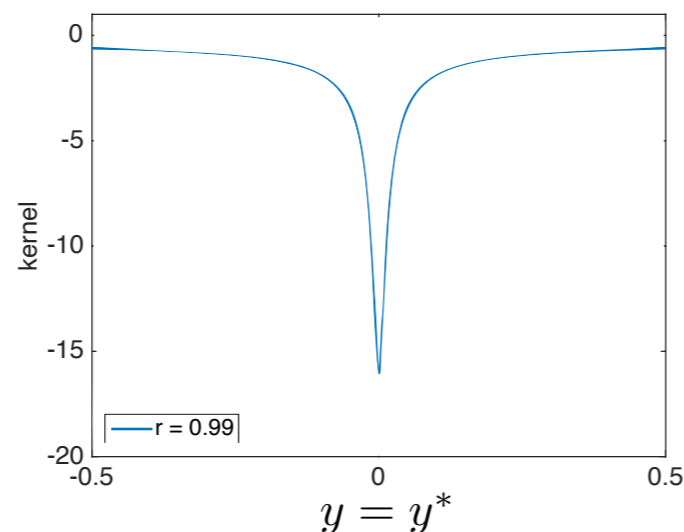
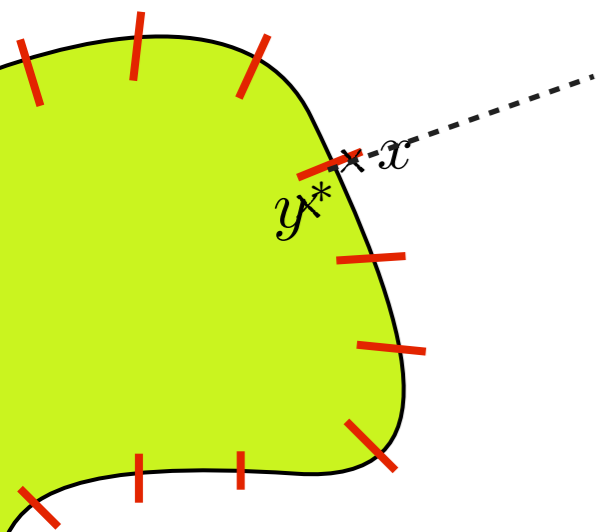
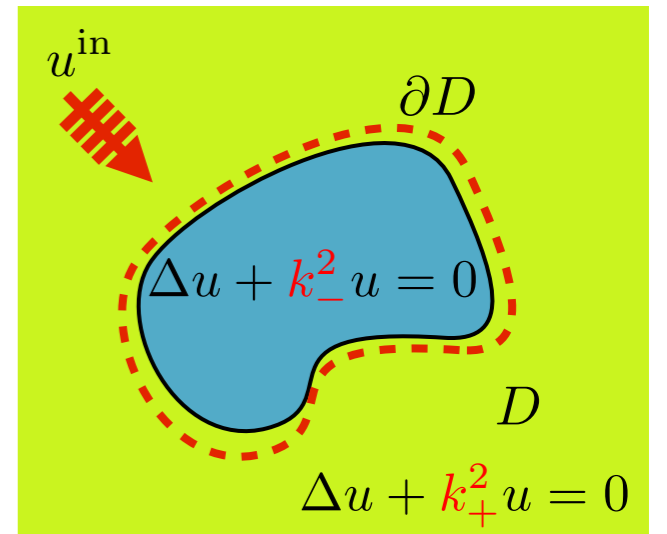
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Nearly singular integral

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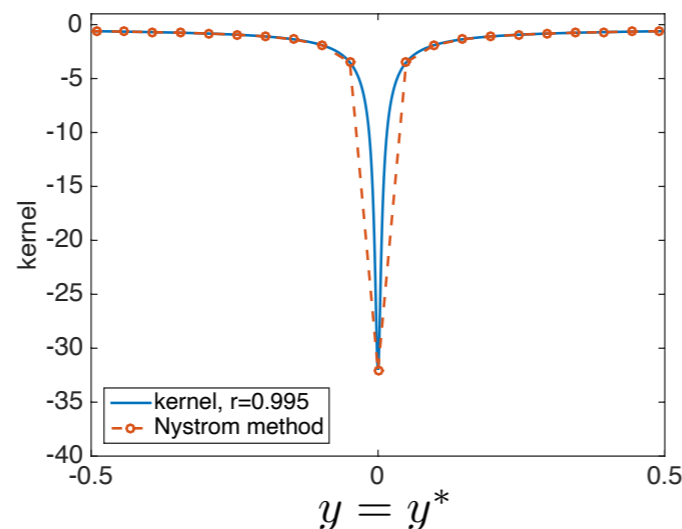
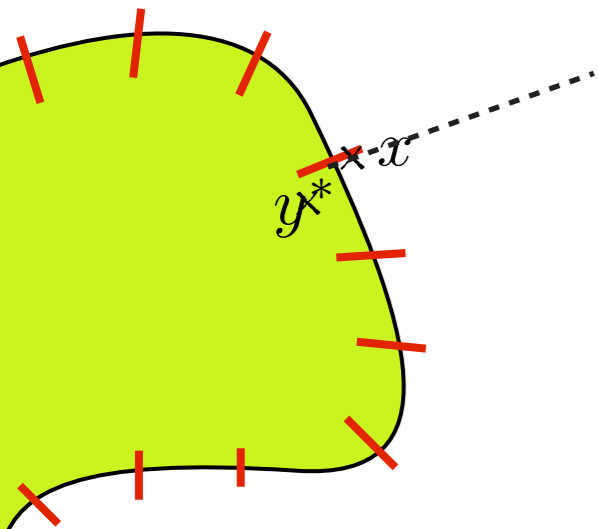
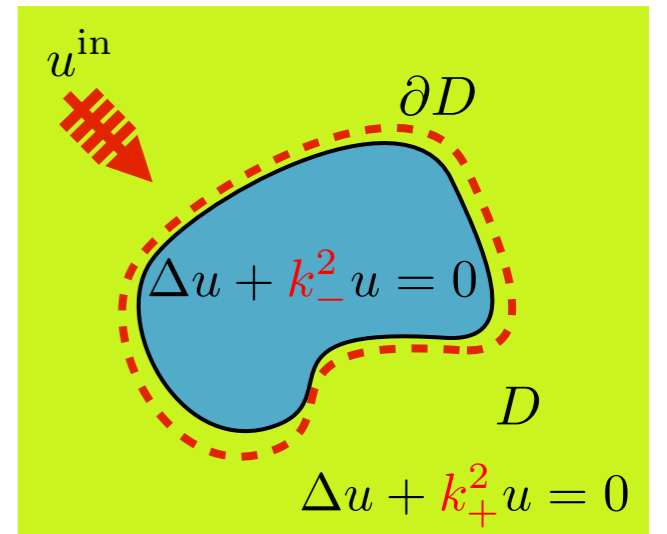
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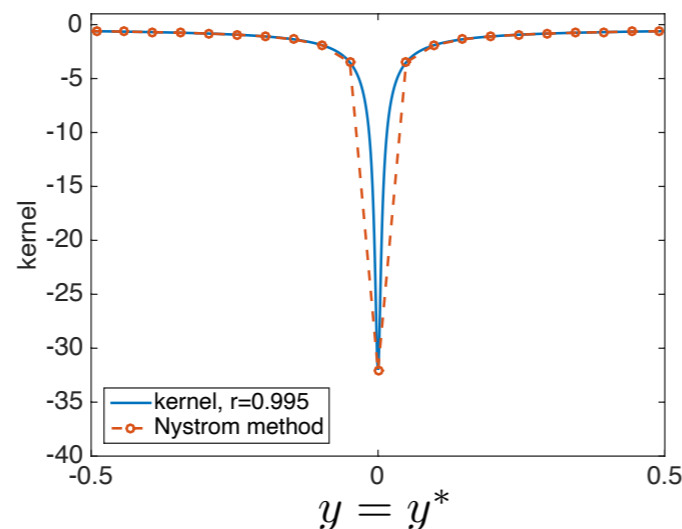
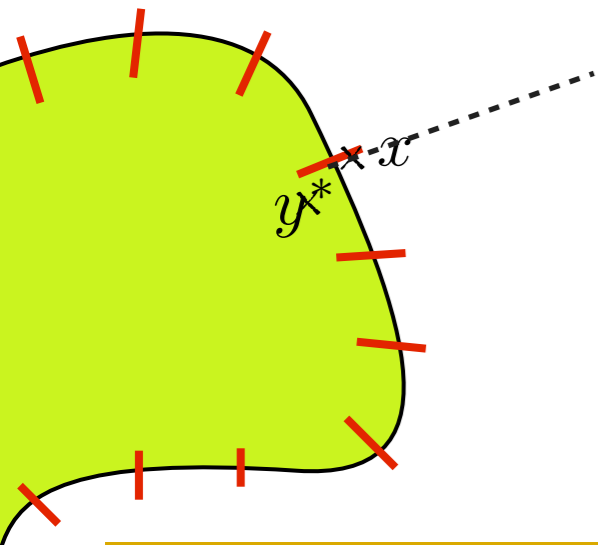
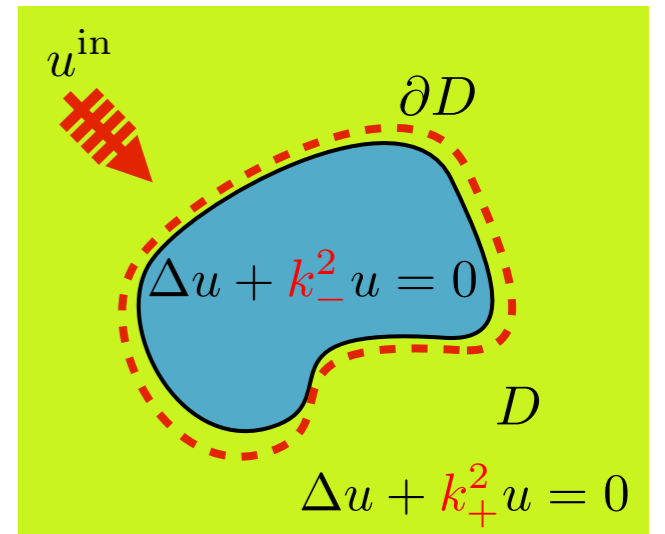
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Nearly singular integral
For a fixed number of
quadrature points, $O(1)$ error.



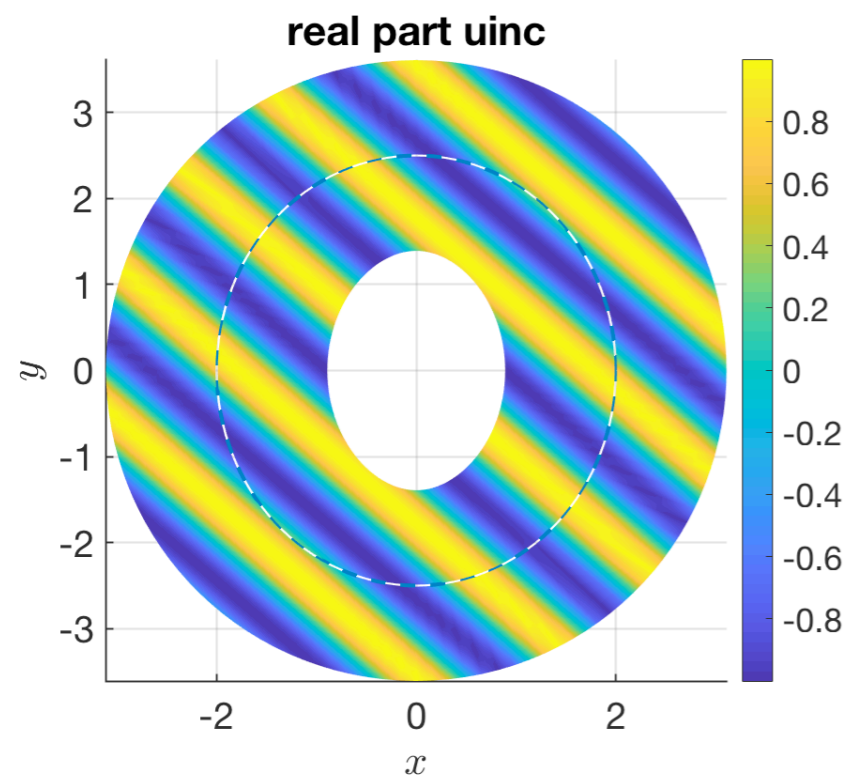
Barnett (2014).

Example

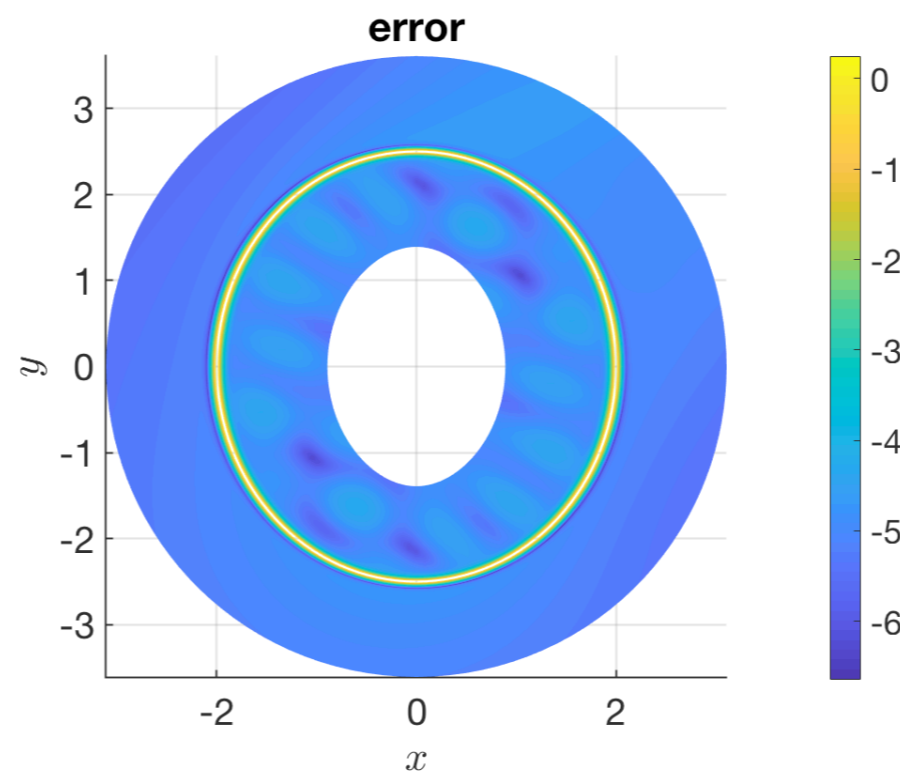
Example of a plane wave in an homogeneous domain (elliptic obstacle).

$$k_-^2 = k_+^2 = 5$$

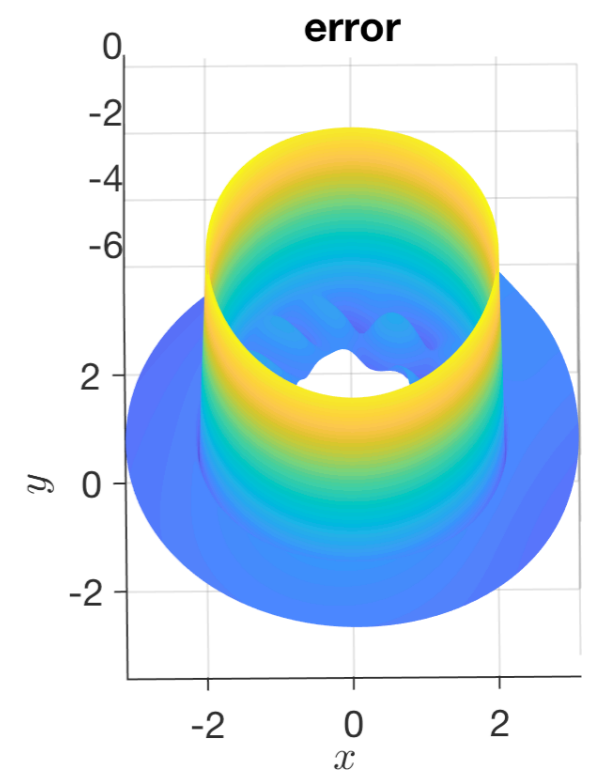
Use 128 points for the quadrature.



Real part solution



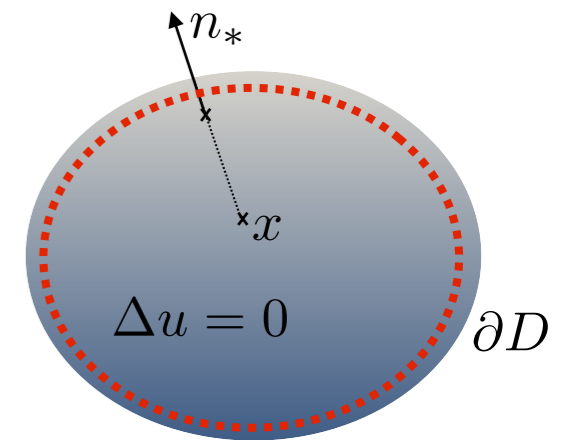
Log plot of the error (2D view, 3D view)



An example in 3D

Interior Dirichlet Laplace problem

$$\Delta u = 0 \quad \text{in } D, \quad u = f \quad \text{on } \partial D$$



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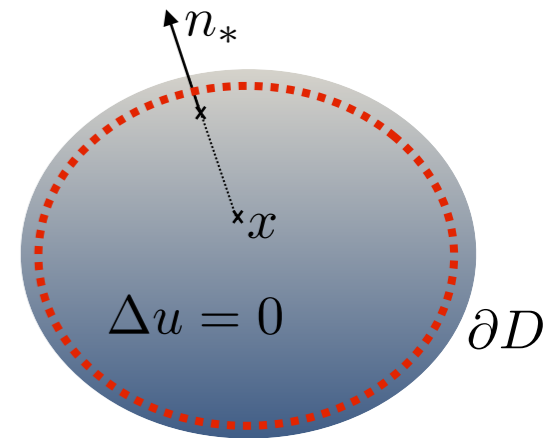
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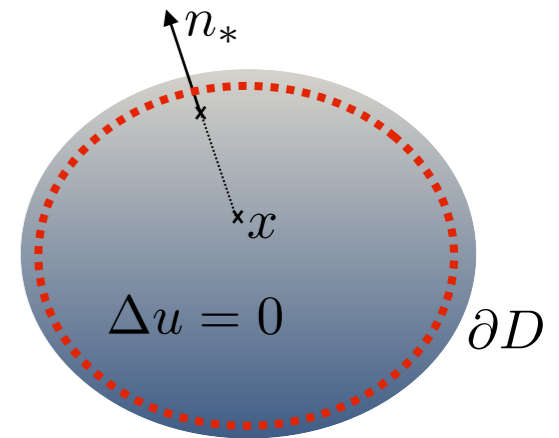
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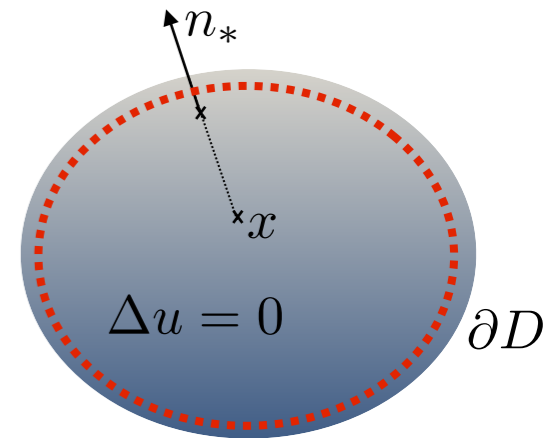
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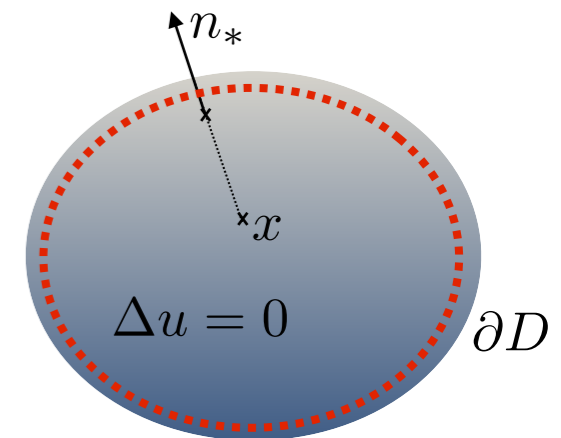
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Using parameterization

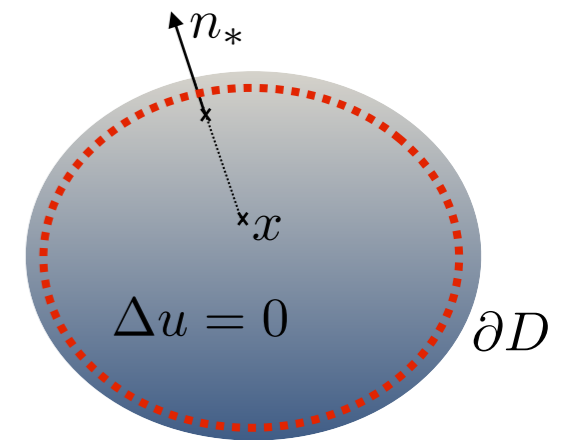
$$y = y(s, t), \quad s \in [0, \pi], \quad t \in [-\pi, \pi]$$

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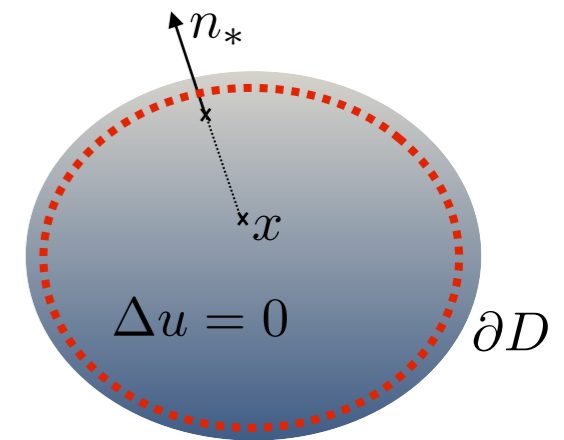


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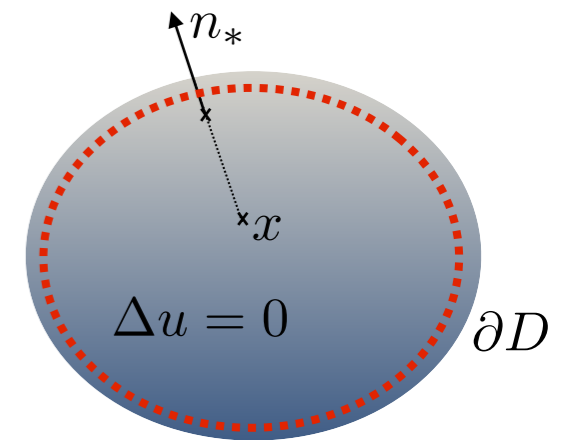
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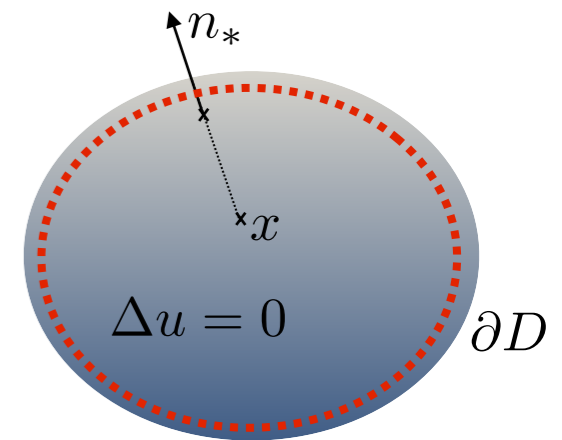
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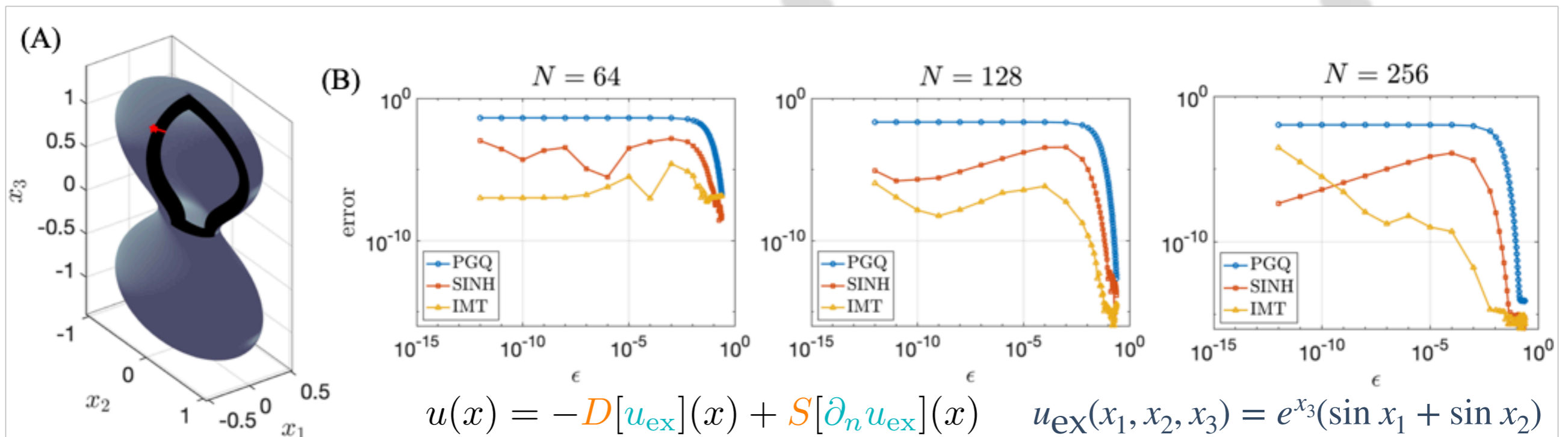
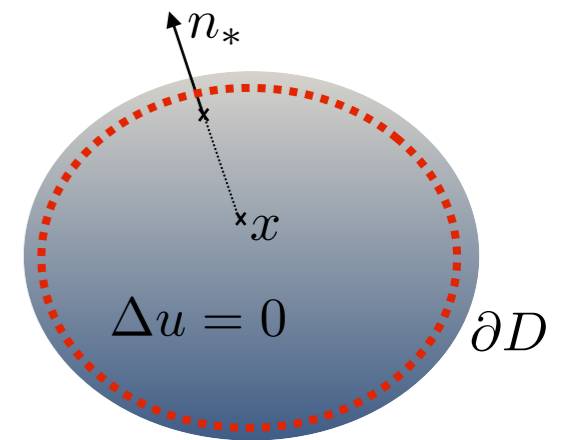
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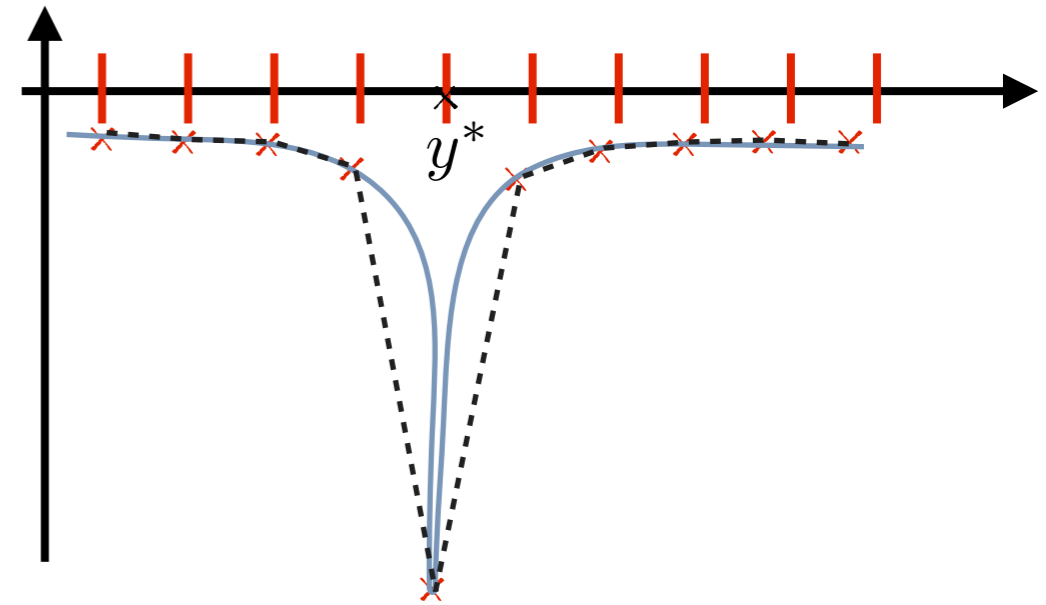
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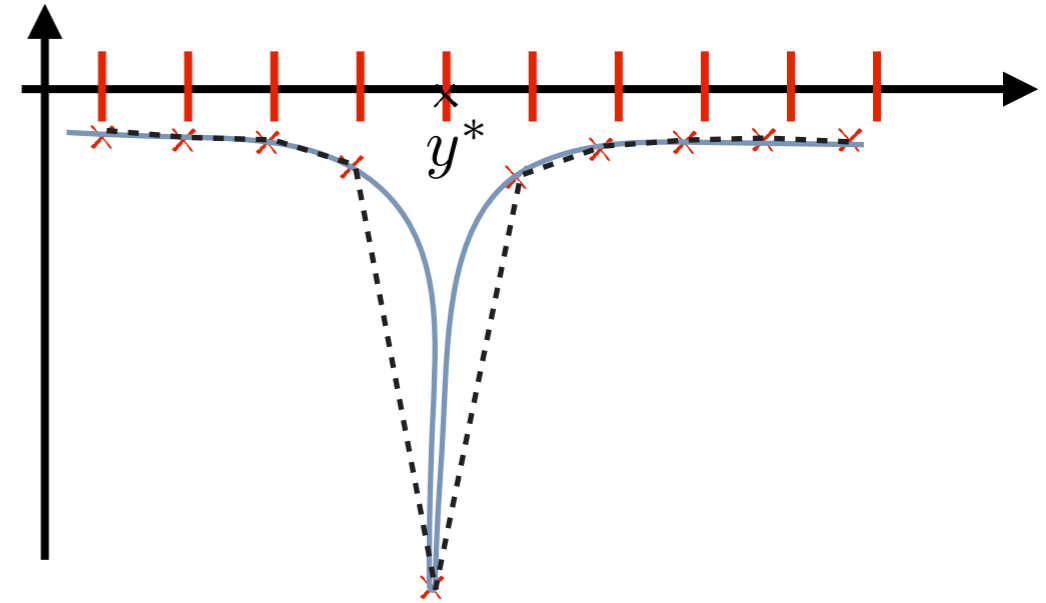


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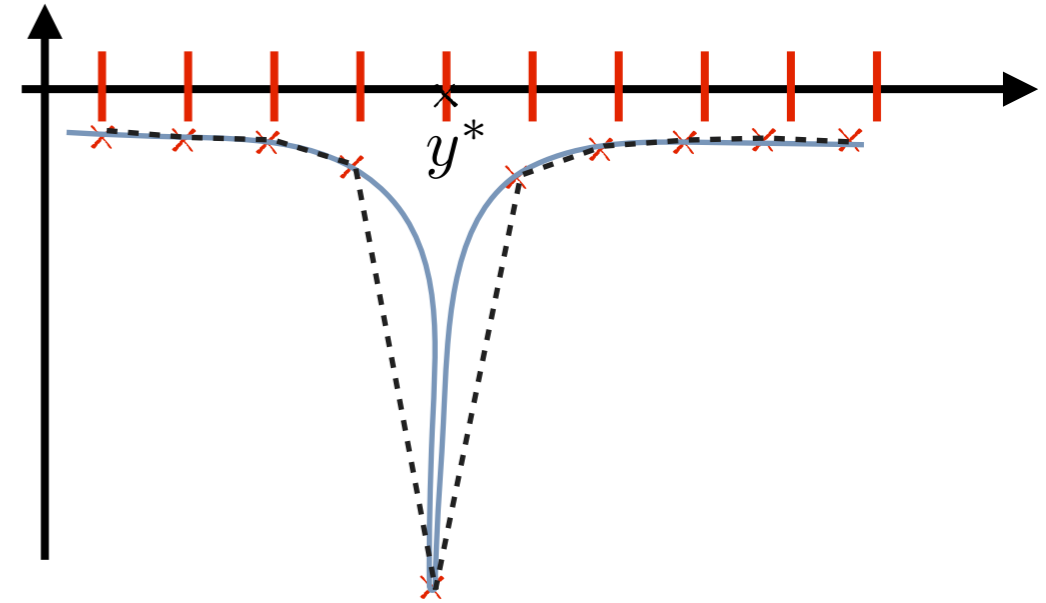
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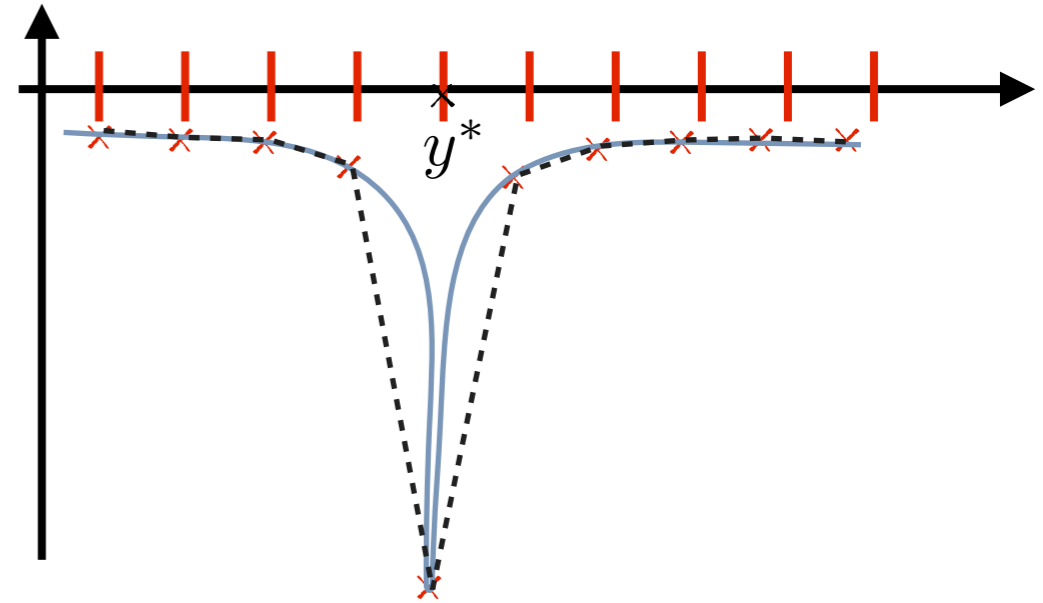


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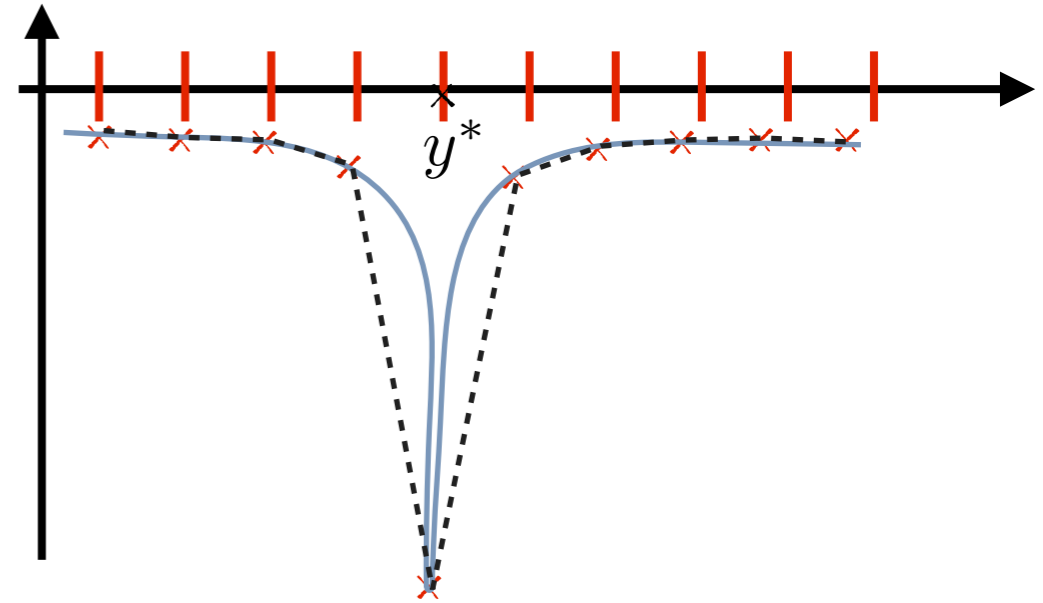
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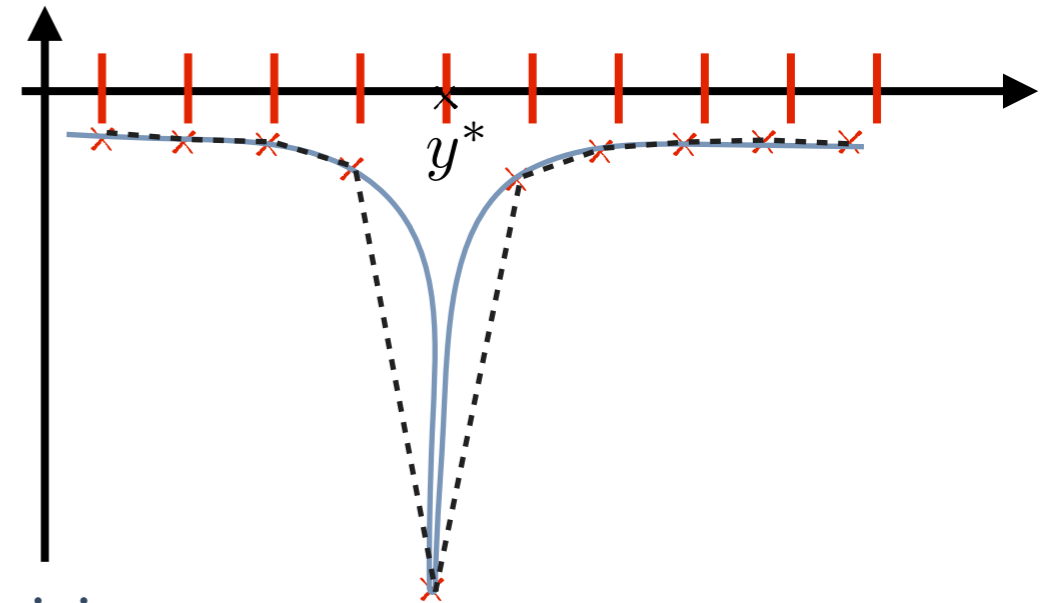
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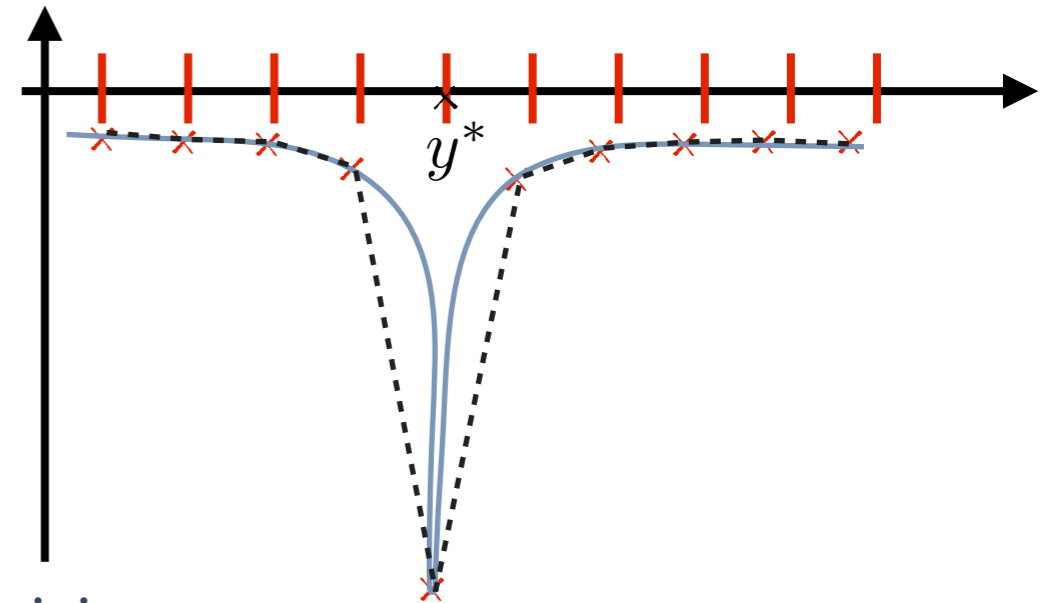
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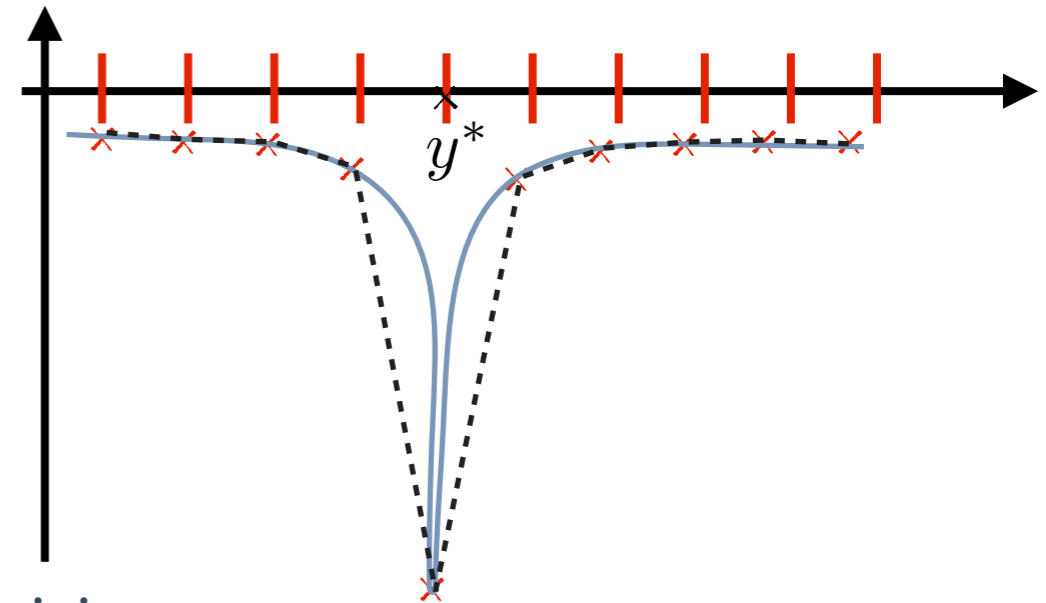
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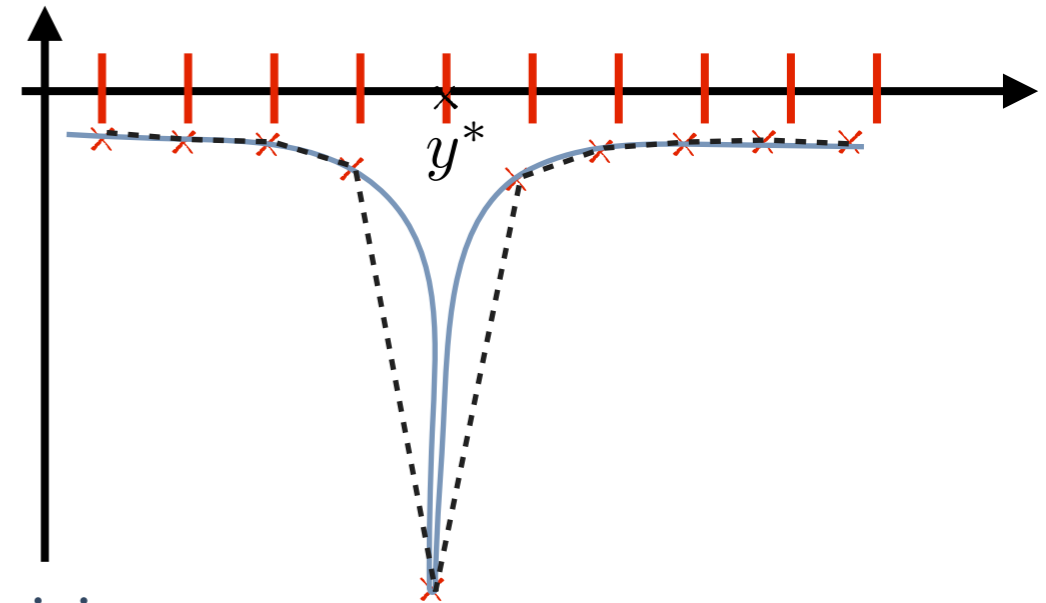
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Given a quadrature method, can we improve accuracy at close evaluation points ?

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C. et al. (2018, 2020, 2021), Khatri et al. (2020), C. (2021)

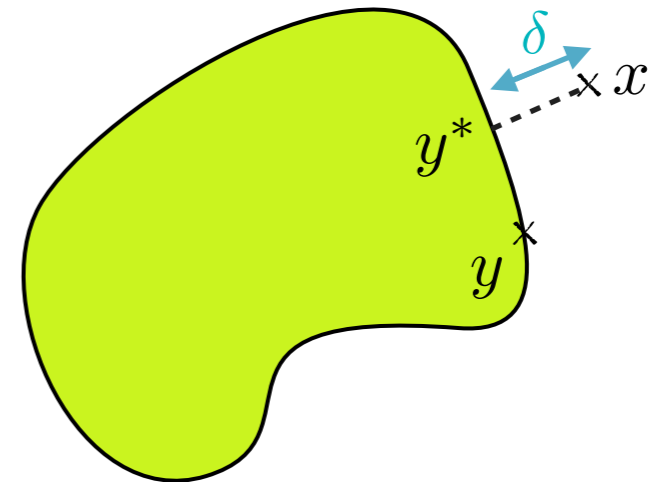
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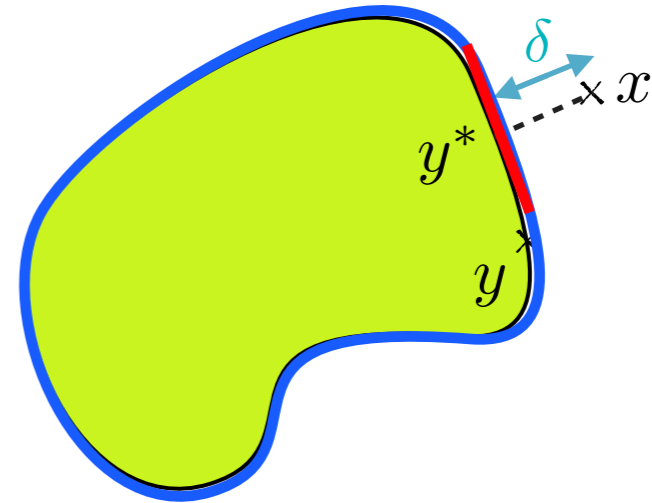
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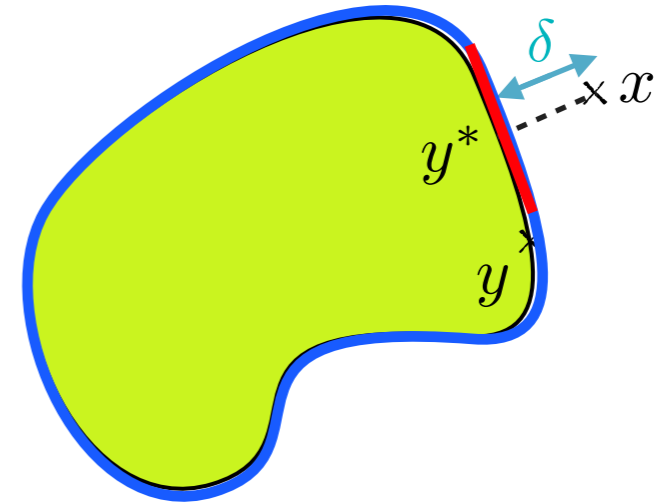
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- $K = K^{\text{in}} + K^{\text{out}} + O(\delta)$
- $\int_{\text{in}} K(x, y) \mu(y) d\sigma_y + \int_{\text{out}} K(x, y) \mu(y) d\sigma_y$

Base numerical method on asymptotic analysis



2) Modified representations

-
-



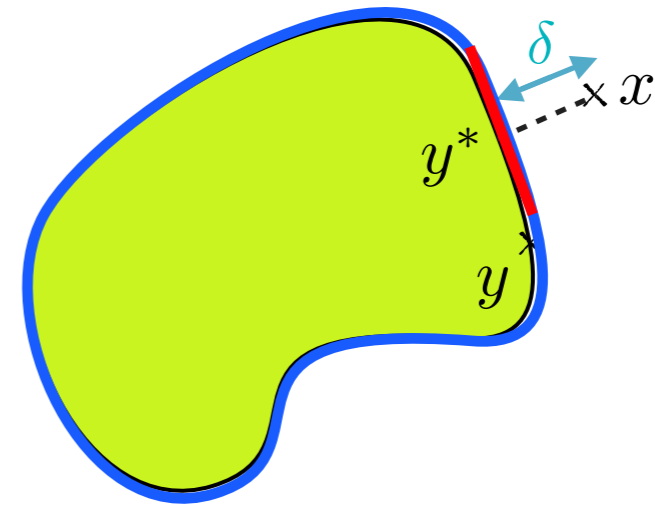
C. et al. (2018, 2020, 2021), Khatri et al. (2020), C. (2021)

Asymptotic methods for close evaluation

$$u(x) = \int_{\partial D} K(x, y) \mu(y) d\sigma_y \quad \text{peaked kernel} \quad \text{continuous function (spectral accuracy)}$$

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Base numerical method on asymptotic analysis

2) Modified representations

- $\int_{\partial D} K(x, y) [\mu(y) - \alpha(x, y)] d\sigma_y + \int_{\partial D} K(x, y) \alpha(x, y) d\sigma_y$
- $\int_{\partial D} [K(x, y) - \tilde{K}(x, y)] \mu(y) d\sigma_y + \int_{\partial D} \tilde{K}(x, y) \mu(y) d\sigma_y$



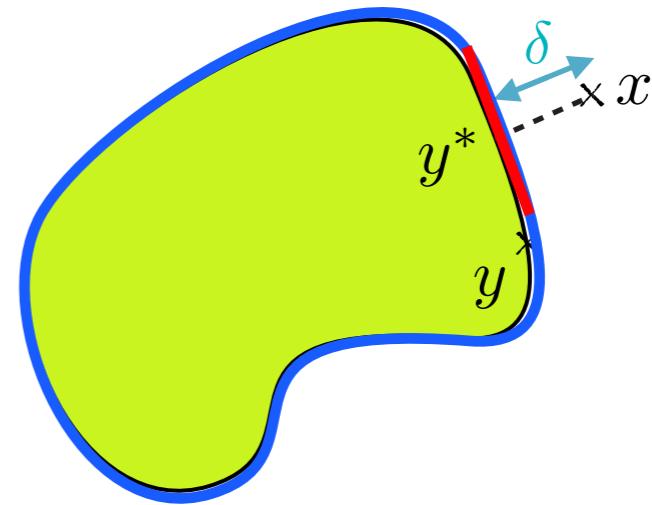
C. et al. (2018, 2020, 2021), Khatri et al. (2020), C. (2021)

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Base numerical method on asymptotic analysis

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- $\int_{\partial D} K(x, y) [\mu(y) - \alpha(x, y)] d\sigma_y + \int_{\partial D} K(x, y) \alpha(x, y) d\sigma_y$
 Vanishes at $x = y$ Spectral computation
- $\int_{\partial D} [K(x, y) - \tilde{K}(x, y)] \mu(y) d\sigma_y + \int_{\partial D} \tilde{K}(x, y) \mu(y) d\sigma_y$



C. et al. (2018, 2020, 2021), Khatri et al. (2020), C. (2021)

Asymptotic methods for close evaluation

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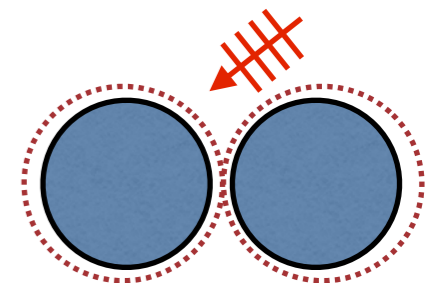
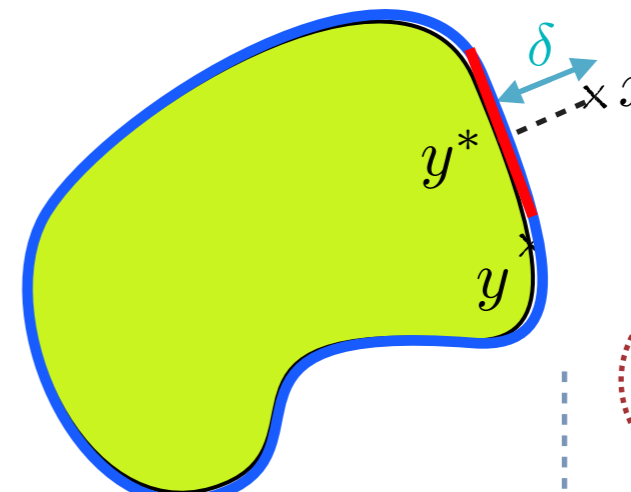
Base numerical method on asymptotic analysis

2) Modified representations

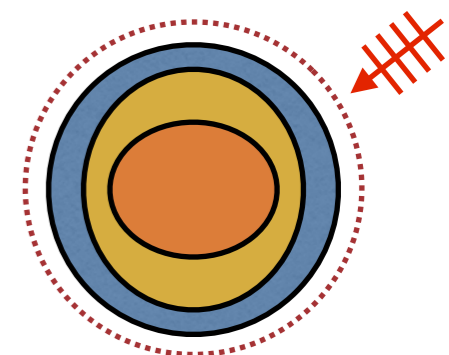
- $\int_{\partial D} K(x, y) [\mu(y) - \alpha(x, y)] d\sigma_y + \int_{\partial D} K(x, y) \alpha(x, y) d\sigma_y$
 Vanishes at $x = y$ Spectral computation
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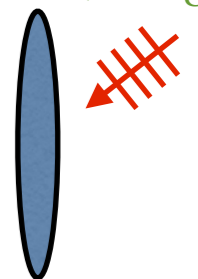
C. et al. (2018, 2020, 2021), Khatri et al. (2020), C. (2021)



Acoustic binding with Kim, McCullough



Optical cloaking with Chaillat, Cortes, Tsogka



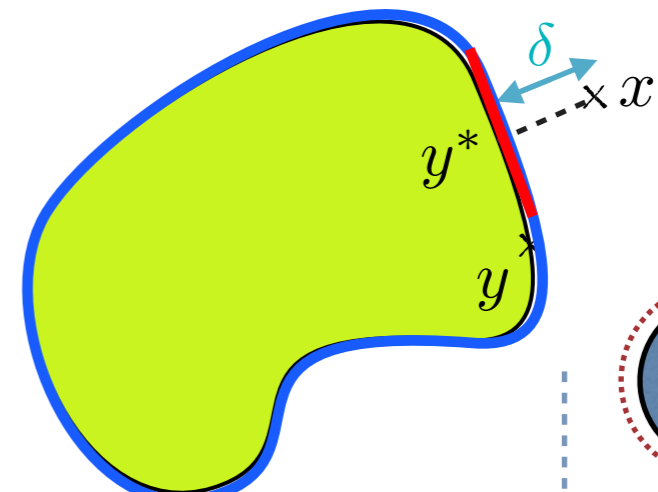
QPAX with Kim, Lewis, Moitier

Asymptotic methods for close evaluation

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- $K = K^{\text{in}} + K^{\text{out}} + O(\delta)$
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Base numerical method on asymptotic analysis in 3D

2) Modified representations

$$\int_{\partial D} K(x, y) [\mu(y) - \alpha(x, y)] d\sigma_y + \int_{\partial D} K(x, y) \alpha(x, y) d\sigma_y$$

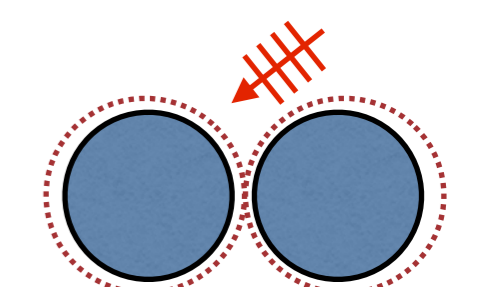
Vanishes at $x = y$

Spectral computation

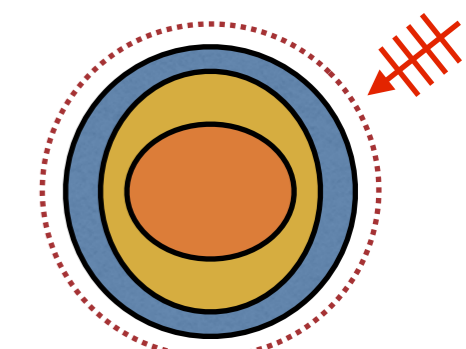
$$\int_{\partial D} [K(x, y) - \tilde{K}(x, y)] \mu(y) d\sigma_y + \int_{\partial D} \tilde{K}(x, y) \mu(y) d\sigma_y$$



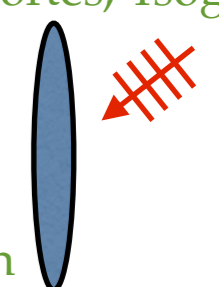
C. et al. (2018, 2020, 2021), Khatri et al. (2020), C. (2021)



Acoustic binding with Kim, McCullough



Optical cloaking with Chaillat, Cortes, Tsogka



QPAX with Kim, Lewis, Moitier

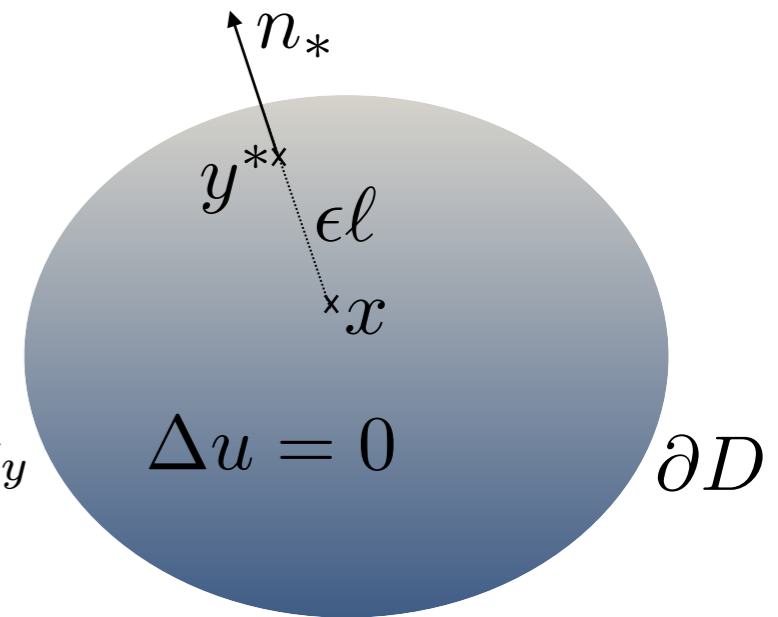
Outline

- ❖ The close evaluation problem
- ❖ Quadrature based on asymptotic methods
- ❖ Modified representations
- ❖ Conclusion

Nearly singular integrals

Interior Dirichlet Laplace problem

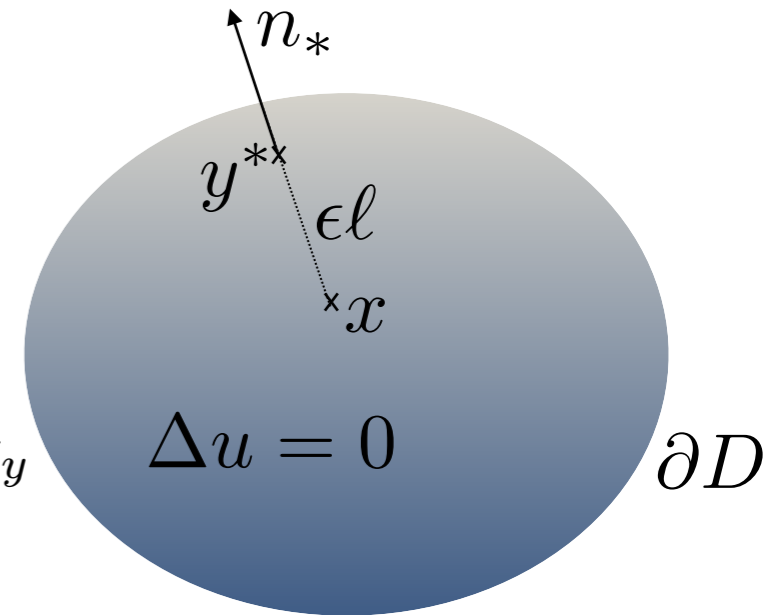
$$\begin{aligned} u(x) &= - \int_{\partial D} \partial_{n_y} G(x, y) u(y) d\sigma_y + \int_{\partial D} G(x, y) \partial_n u(y) d\sigma_y \\ &= - \frac{1}{4\pi} \int_{\partial D} \frac{n_y \cdot (x - y)}{|x - y|^3} u(y) d\sigma_y + \frac{1}{4\pi} \int_{\partial D} \frac{1}{|x - y|} \partial_n u(y) d\sigma_y \end{aligned}$$



Nearly singular integrals

Interior Dirichlet Laplace problem

$$\begin{aligned}
 u(x) &= - \int_{\partial D} \partial_{n_y} G(x, y) u(y) d\sigma_y + \int_{\partial D} G(x, y) \partial_n u(y) d\sigma_y \\
 &= - \frac{1}{4\pi} \int_{\partial D} \frac{n_y \cdot (x - y)}{|x - y|^3} u(y) d\sigma_y + \frac{1}{4\pi} \int_{\partial D} \frac{1}{|x - y|} \partial_n u(y) d\sigma_y
 \end{aligned}$$



Using $x = y^* - \epsilon\ell n_*$ and Gauss's law:

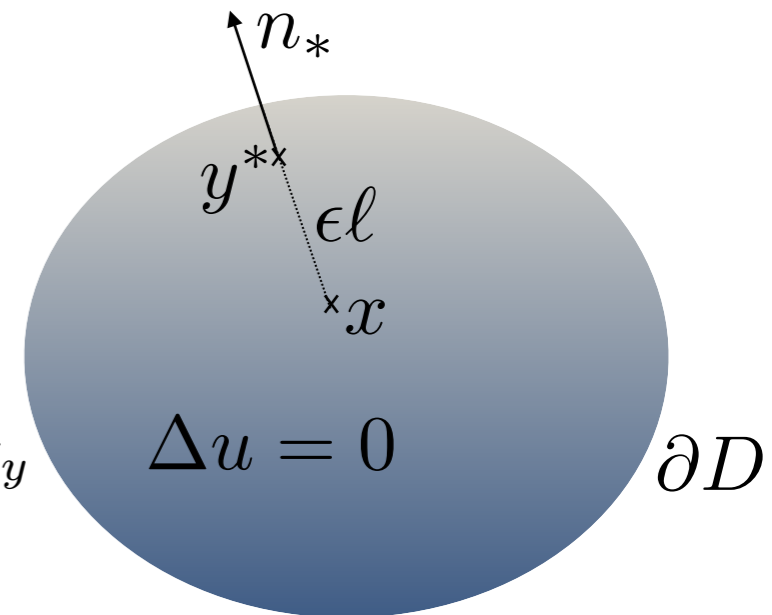
$$\int_{\partial D} \partial_{n_y} G(x, y) d\sigma_y = \begin{cases} 0 & x \in \mathbb{R}^3 \setminus \bar{D} \\ -\frac{1}{2} & x \in \partial D \\ -1 & x \in D \end{cases}$$

Nearly singular integrals

Interior Dirichlet Laplace problem

$$u(x) = - \int_{\partial D} \partial_{n_y} G(x, y) u(y) d\sigma_y + \int_{\partial D} G(x, y) \partial_n u(y) d\sigma_y$$

$$= - \frac{1}{4\pi} \int_{\partial D} \frac{n_y \cdot (x - y)}{|x - y|^3} u(y) d\sigma_y + \frac{1}{4\pi} \int_{\partial D} \frac{1}{|x - y|} \partial_n u(y) d\sigma_y$$



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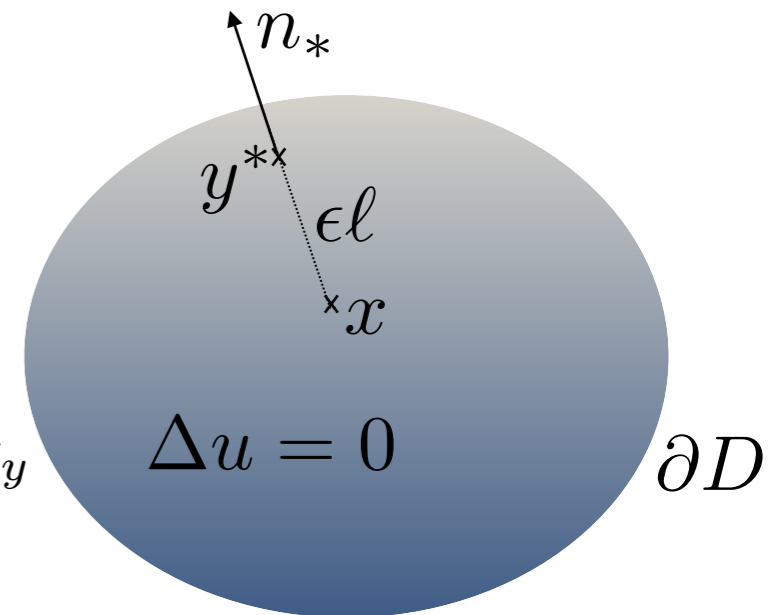
$$= u(y^*) - \frac{1}{4\pi} \int_{\partial D} \frac{n_y \cdot (y^* - y - \epsilon\ell n_*)}{|y^* - y - \epsilon\ell n_*|^3} [u(y) - u(y^*)] d\sigma_y + \frac{1}{4\pi} \int_{\partial D} \frac{1}{|y^* - y - \epsilon\ell n_*|} \partial_n u(y) d\sigma_y$$

Nearly singular integrals

Interior Dirichlet Laplace problem

$$u(x) = - \int_{\partial D} \partial_{n_y} G(x, y) u(y) d\sigma_y + \int_{\partial D} G(x, y) \partial_n u(y) d\sigma_y$$

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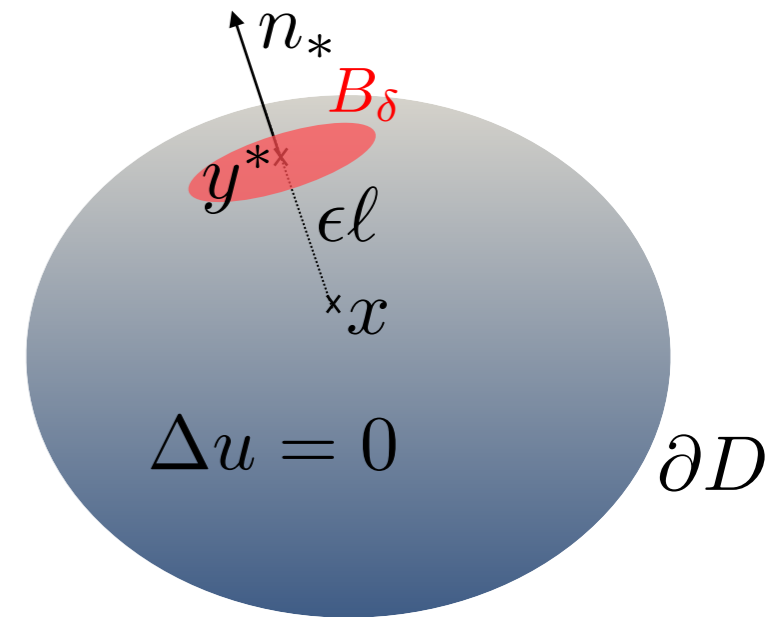
$$= u(y^*) - \frac{1}{4\pi} \int_{\partial D} \frac{n_y \cdot (y^* - y - \epsilon\ell n_*)}{|y^* - y - \epsilon\ell n_*|^3} [u(y) - u(y^*)] d\sigma_y + \frac{1}{4\pi} \int_{\partial D} \frac{1}{|y^* - y - \epsilon\ell n_*|} \partial_n u(y) d\sigma_y$$

local analysis of each layer potential about y^* when $\epsilon \rightarrow 0^+$

Close evaluation in 3D

Local Analysis of the Laplace double-layer potential:

$$-\frac{1}{4\pi} \int_{B_\delta} \frac{n_y \cdot (y^* - y - \epsilon \ell n_*)}{|y^* - y - \epsilon \ell n_*|^3} [u(y) - u(y^*)] d\sigma_y$$

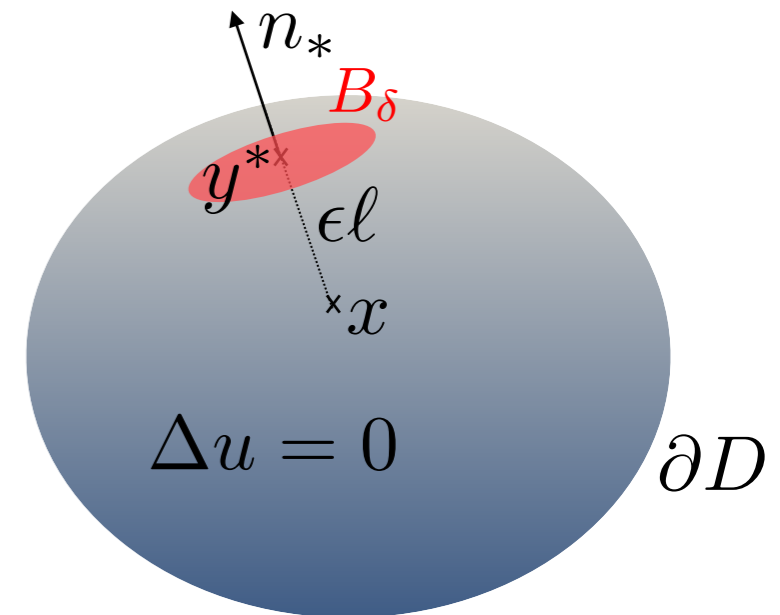


Close evaluation in 3D

Local Analysis of the Laplace double-layer potential:

$$-\frac{1}{4\pi} \int_{B_\delta} \frac{n_y \cdot (y^* - y - \epsilon \ell n_*)}{|y^* - y - \epsilon \ell n_*|^3} [u(y) - u(y^*)] d\sigma_y$$

1) Parameterization $y(s, t)$ with $(s, t) \in [0, \pi] \times [-\pi, \pi]$
with y^* corresponding to the north pole



Close evaluation in 3D

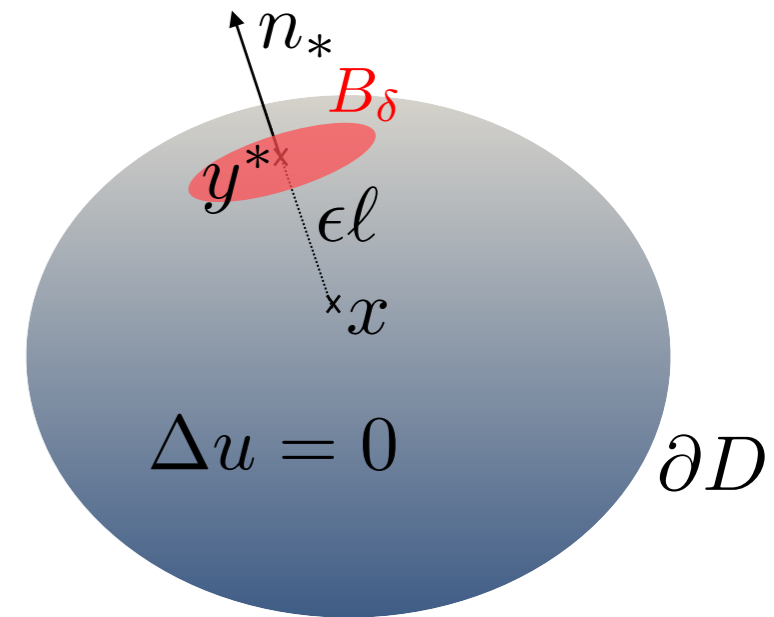
Local Analysis of the Laplace double-layer potential:

$$-\frac{1}{4\pi} \int_{B_\delta} \frac{n_y \cdot (y^* - y - \epsilon l n_*)}{|y^* - y - \epsilon l n_*|^3} [u(y) - u(y^*)] d\sigma_y$$

1) Parameterization $y(s, t)$ with $(s, t) \in [0, \pi] \times [-\pi, \pi]$
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$$-\frac{1}{4\pi} \int_{-\pi}^{\pi} \int_0^\delta \frac{\tilde{n}(s, t) \cdot (y(0, \cdot) - y(s, t) - \epsilon l n_*)}{|(y(0, \cdot) - y(s, t) - \epsilon l n_*)|^3} J(s, t) [\tilde{u}(s, t) - \tilde{u}(0, \cdot)] \sin(s) ds dt$$

$$J(s, t) = \frac{|y_s(s, t) \times y_t(s, t)|}{\sin(s)}$$



Close evaluation in 3D

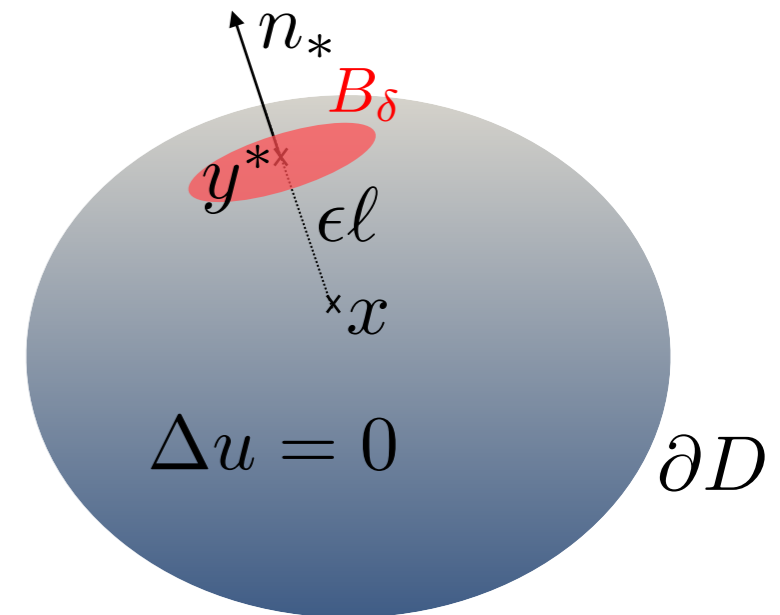
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2) Let $s = \epsilon S$ and expand about $\epsilon = 0$



$$J(s, t) = \frac{|y_s(s, t) \times y_t(s, t)|}{\sin(s)}$$

Close evaluation in 3D

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$$-\frac{1}{4\pi} \int_{B_\delta} \frac{n_y \cdot (y^* - y - \epsilon l n_*)}{|y^* - y - \epsilon l n_*|^3} [u(y) - u(y^*)] d\sigma_y$$

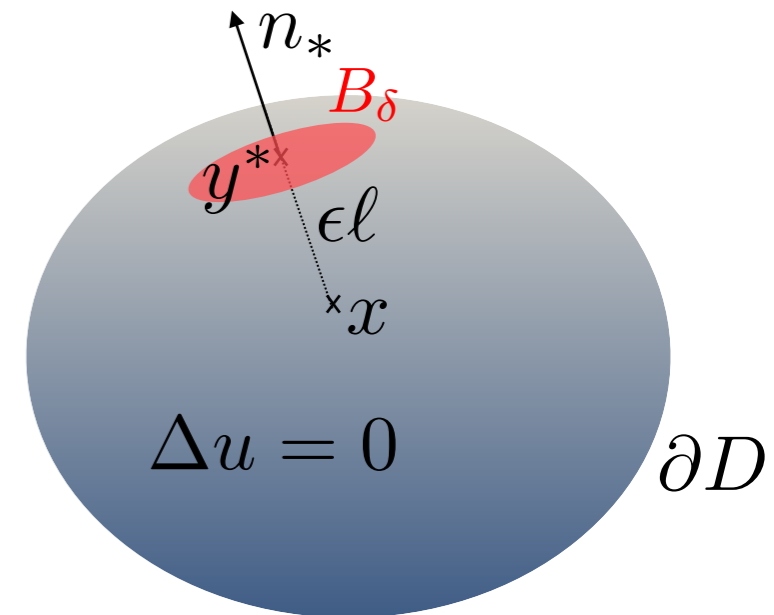
1) Parameterization $y(s, t)$ with $(s, t) \in [0, \pi] \times [-\pi, \pi]$ with y^* corresponding to the north pole

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$$J(s, t) = \frac{|y_s(s, t) \times y_t(s, t)|}{\sin(s)}$$

2) Let $s = \epsilon S$ and expand about $\epsilon = 0$

$$\frac{1}{2} \int_0^{\delta/\epsilon} \left(\frac{\ell J(0, \cdot) S}{(S^2 |y_s(0, \cdot)|^2 + \ell^2)^{3/2}} + O(\epsilon) \right) \left[\frac{1}{2\pi} \int_{-\pi}^{\pi} \tilde{u}(\epsilon S, t) - \tilde{u}(0, \cdot) dt \right] dS$$



Close evaluation in 3D

Local Analysis of the Laplace double-layer potential:

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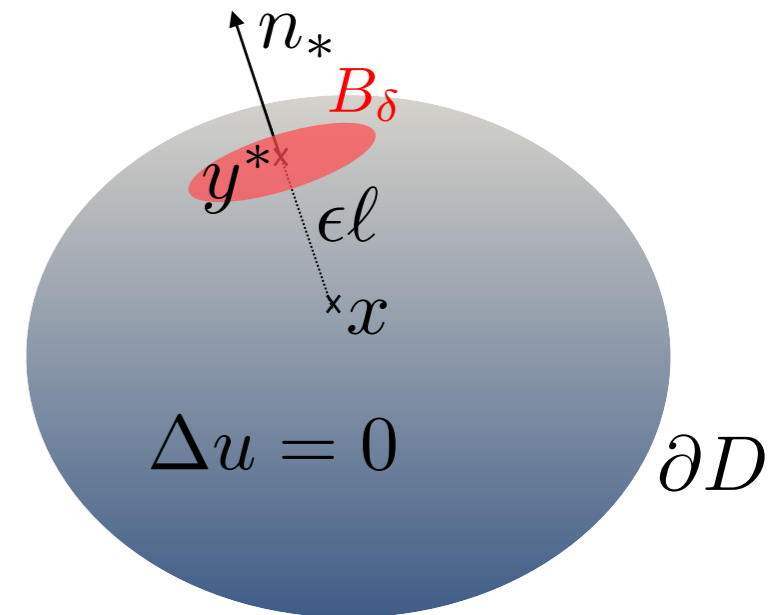
$$-\frac{1}{4\pi} \int_{-\pi}^{\pi} \int_0^{\delta} \frac{\tilde{n}(s, t) \cdot (y(0, \cdot) - y(s, t) - \epsilon \ell n_*)}{|(y(0, \cdot) - y(s, t) - \epsilon \ell n_*)|^3} J(s, t) [\tilde{u}(s, t) - \tilde{u}(0, \cdot)] \sin(s) ds dt$$

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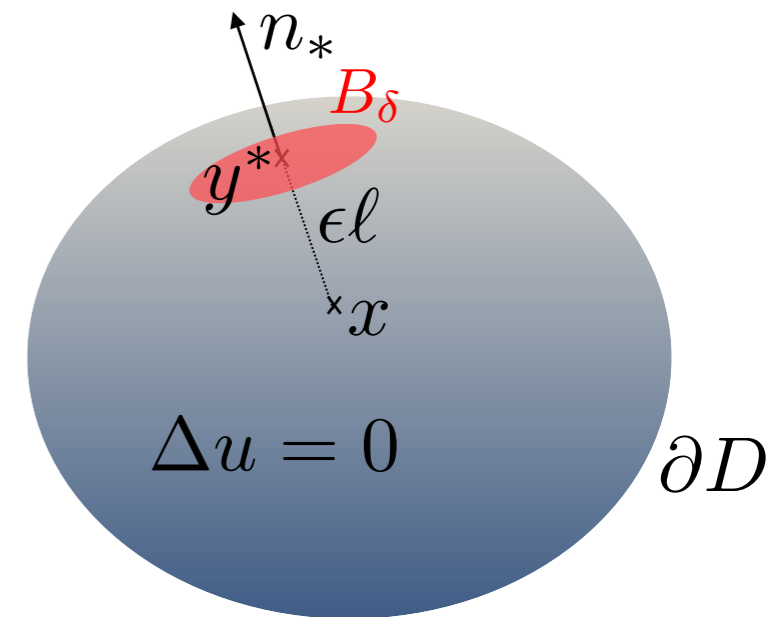
$$= \delta \frac{\epsilon \ell J(0, \cdot)}{8 |y_s(0, \cdot)|^3} \Delta_S \tilde{u}(0, \cdot) + O(\epsilon^2)$$



Close evaluation in 3D

Local Analysis of the single-layer potential:

$$\frac{1}{4\pi} \int_{B_\delta} \frac{1}{|y^* - y - \epsilon \ell n_*|} \partial_n u(y) d\sigma_y$$



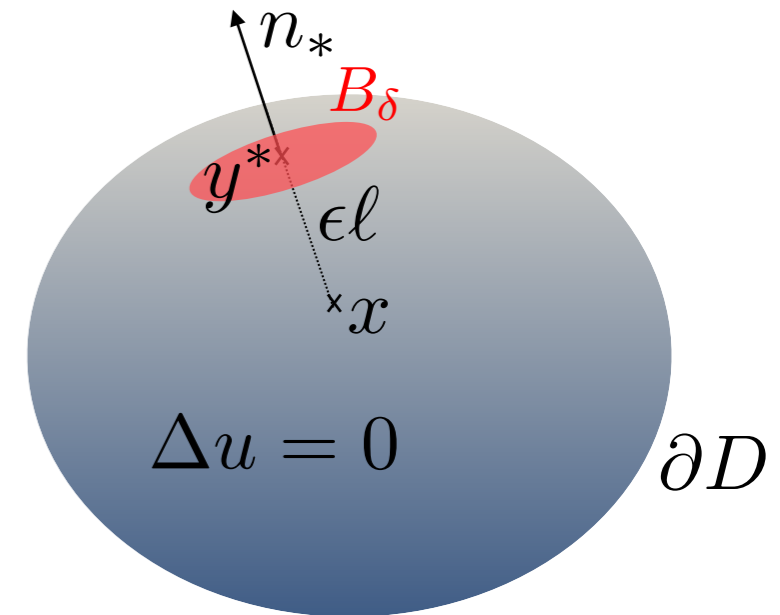
Close evaluation in 3D

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$$\frac{1}{4\pi} \int_{B_\delta} \frac{1}{|y^* - y - \epsilon \ell n_*|} \partial_n u(y) d\sigma_y$$

Steps 1) - 2)

$$\frac{1}{2} \int_0^{\delta/\epsilon} \left(\frac{J(0, \cdot) \epsilon S}{(S^2 |y_s(0, \cdot)|^2 + \ell^2)^{1/2}} + O(\epsilon^2) \right) \partial_n \tilde{u}(0, \cdot) dS = \delta \frac{J(0, \cdot)}{2|y_s(0, \cdot)|} \partial_n \tilde{u}(0, \cdot) + O(\epsilon)$$



Close evaluation in 3D

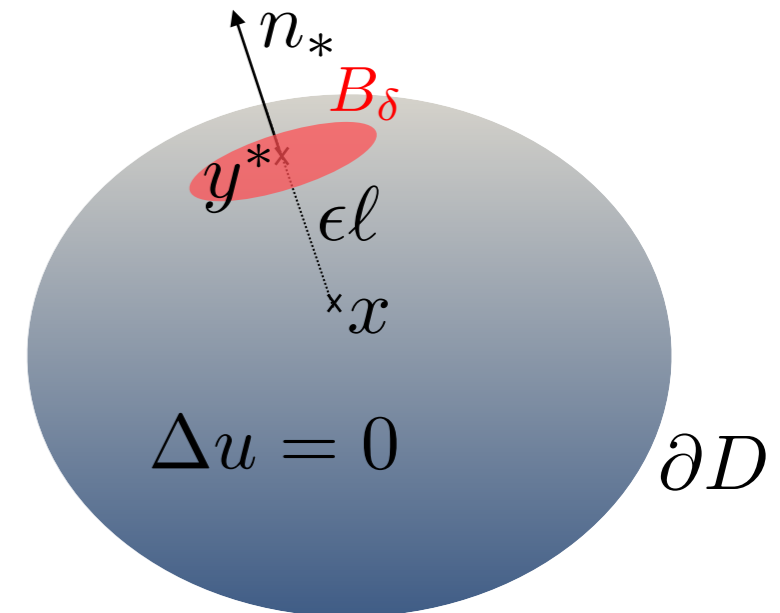
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From the local analysis:



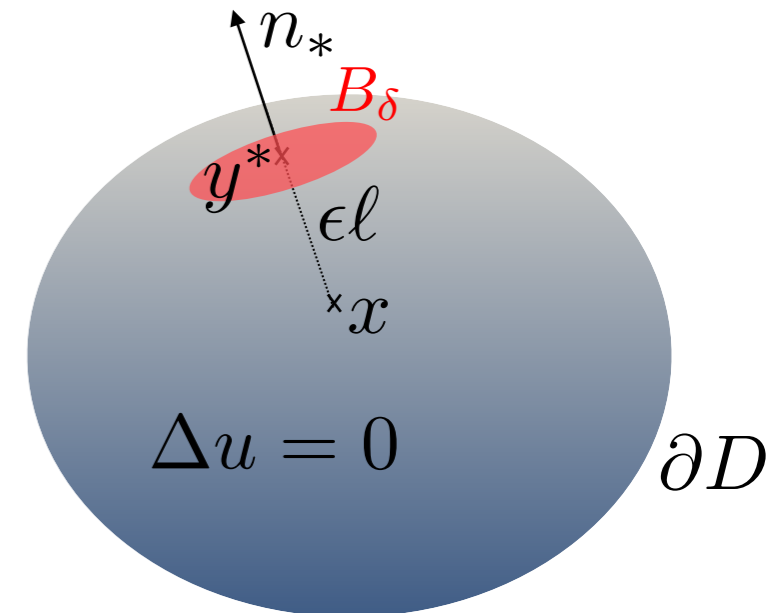
Close evaluation in 3D

Local Analysis of the single-layer potential:

$$\frac{1}{4\pi} \int_{B_\delta} \frac{1}{|y^* - y - \epsilon \ell n_*|} \partial_n u(y) d\sigma_y$$

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$$\frac{1}{2} \int_0^{\delta/\epsilon} \left(\frac{J(0, \cdot) \epsilon S}{(S^2 |y_s(0, \cdot)|^2 + \ell^2)^{1/2}} + O(\epsilon^2) \right) \partial_n \tilde{u}(0, \cdot) dS = \delta \frac{J(0, \cdot)}{2|y_s(0, \cdot)|} \partial_n \tilde{u}(0, \cdot) + O(\epsilon)$$



From the local analysis:

The **kernels** are azimuthally invariant about y^* as $\epsilon \rightarrow 0$

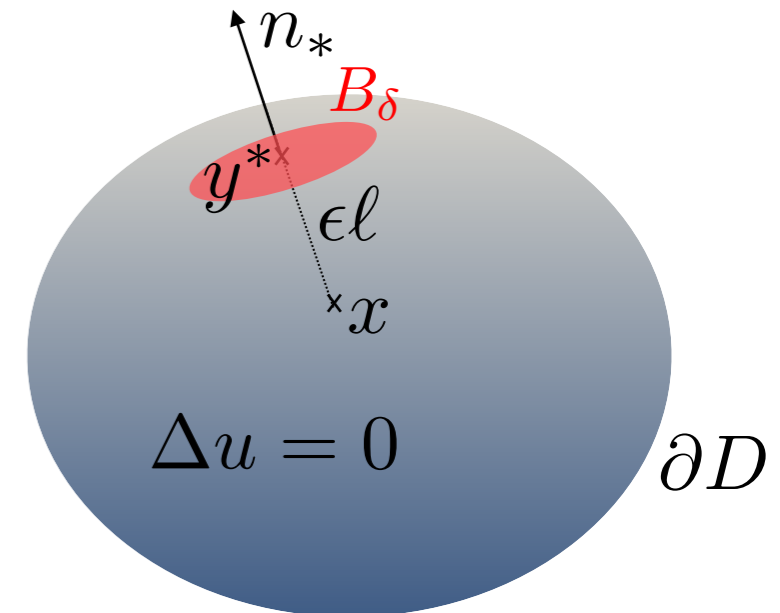
Close evaluation in 3D

Local Analysis of the single-layer potential:

$$\frac{1}{4\pi} \int_{B_\delta} \frac{1}{|y^* - y - \epsilon \ell n_*|} \partial_n u(y) d\sigma_y$$

Steps 1) - 2)

$$\frac{1}{2} \int_0^{\delta/\epsilon} \left(\frac{J(0, \cdot) \epsilon S}{(S^2 |y_s(0, \cdot)|^2 + \ell^2)^{1/2}} + O(\epsilon^2) \right) \partial_n \tilde{u}(0, \cdot) dS = \delta \frac{J(0, \cdot)}{2|y_s(0, \cdot)|} \partial_n \tilde{u}(0, \cdot) + O(\epsilon)$$



From the local analysis:

The **kernels** are azimuthally invariant about y^* as $\epsilon \rightarrow 0$

A **rotated spherical coordinate** system enhances the asymptotic behavior of the kernels

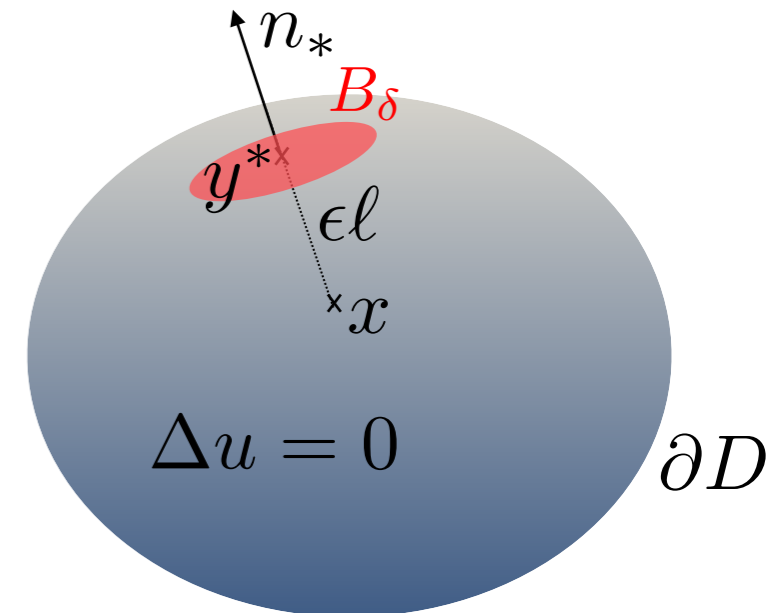
Close evaluation in 3D

Local Analysis of the single-layer potential:

$$\frac{1}{4\pi} \int_{B_\delta} \frac{1}{|y^* - y - \epsilon \ell n_*|} \partial_n u(y) d\sigma_y$$

Steps 1) - 2)

$$\frac{1}{2} \int_0^{\delta/\epsilon} \left(\frac{J(0, \cdot) \epsilon S}{(S^2 |y_s(0, \cdot)|^2 + \ell^2)^{1/2}} + O(\epsilon^2) \right) \partial_n \tilde{u}(0, \cdot) dS = \delta \frac{J(0, \cdot)}{2|y_s(0, \cdot)|} \partial_n \tilde{u}(0, \cdot) + O(\epsilon)$$



From the local analysis:

The **kernels** are azimuthally invariant about y^* as $\epsilon \rightarrow 0$

A **rotated spherical coordinate** system enhances the asymptotic behavior of the kernels

Explicit use of the spherical Jacobian **$\sin(s)$** results in smoother integrands

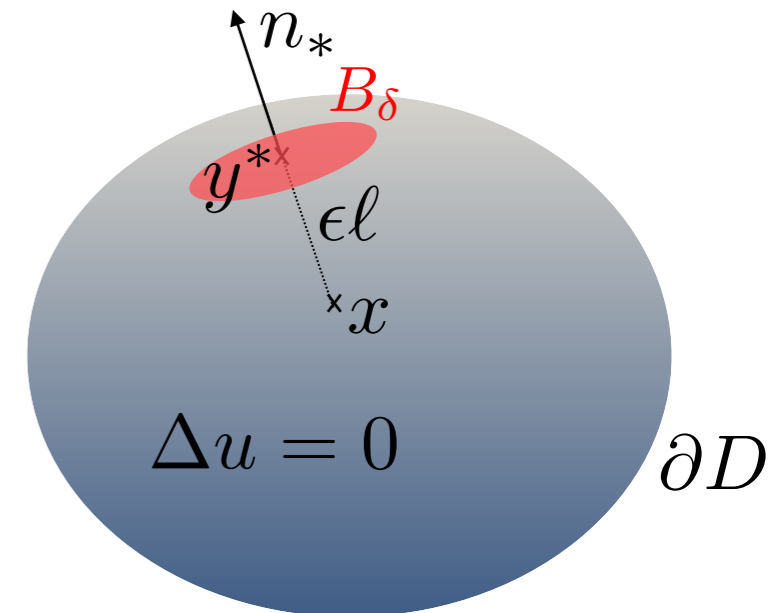
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Local Analysis of the single-layer potential:

$$\frac{1}{4\pi} \int_{B_\delta} \frac{1}{|y^* - y - \epsilon \ell n_*|} \partial_n u(y) d\sigma_y$$

Steps 1) - 2)

$$\frac{1}{2} \int_0^{\delta/\epsilon} \left(\frac{J(0, \cdot) \epsilon S}{(S^2 |y_s(0, \cdot)|^2 + \ell^2)^{1/2}} + O(\epsilon^2) \right) \partial_n \tilde{u}(0, \cdot) dS = \delta \frac{J(0, \cdot)}{2|y_s(0, \cdot)|} \partial_n \tilde{u}(0, \cdot) + O(\epsilon)$$



From the local analysis:

The **kernels** are azimuthally invariant about y^* as $\epsilon \rightarrow 0$

A **rotated spherical coordinate** system enhances the asymptotic behavior of the kernels

Explicit use of the spherical Jacobian $\sin(s)$ results in smoother integrands

Using the three above guidelines, how to proceed numerically ?

Numerical Method

Given a representation

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Atkinson (1982), Gimbutas, Veerapaneni (2013)

2) N-point Gauss Legendre $z_i \in (-1, 1)$, $i = 1, \dots, N$
with weights w_i

3) 2N Periodic Trapezoid rule

$$t_j = -\pi + \pi \frac{j-1}{N}, j = 1, \dots, 2N$$

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$$u^N(y^*; \epsilon) = \frac{\pi}{8N} \sum_{i=1}^N \sum_{j=1}^{2N} w_i \sin(s_i) F(s_i, t_j)$$

Other Quadrature Rules

$$u^N(y^*; \epsilon) = \frac{\pi}{N} \sum_{i=1}^N \sum_{j=1}^{2N} w_i F(s_i, t_j)$$

$$s_i = \cos^{-1}(z_i)$$

Product Gaussian Quadrature rule  Atkinson (1982)

SINH rule  Johnston, Elliott (2005)

IMT rule  Iri et al. (1987)

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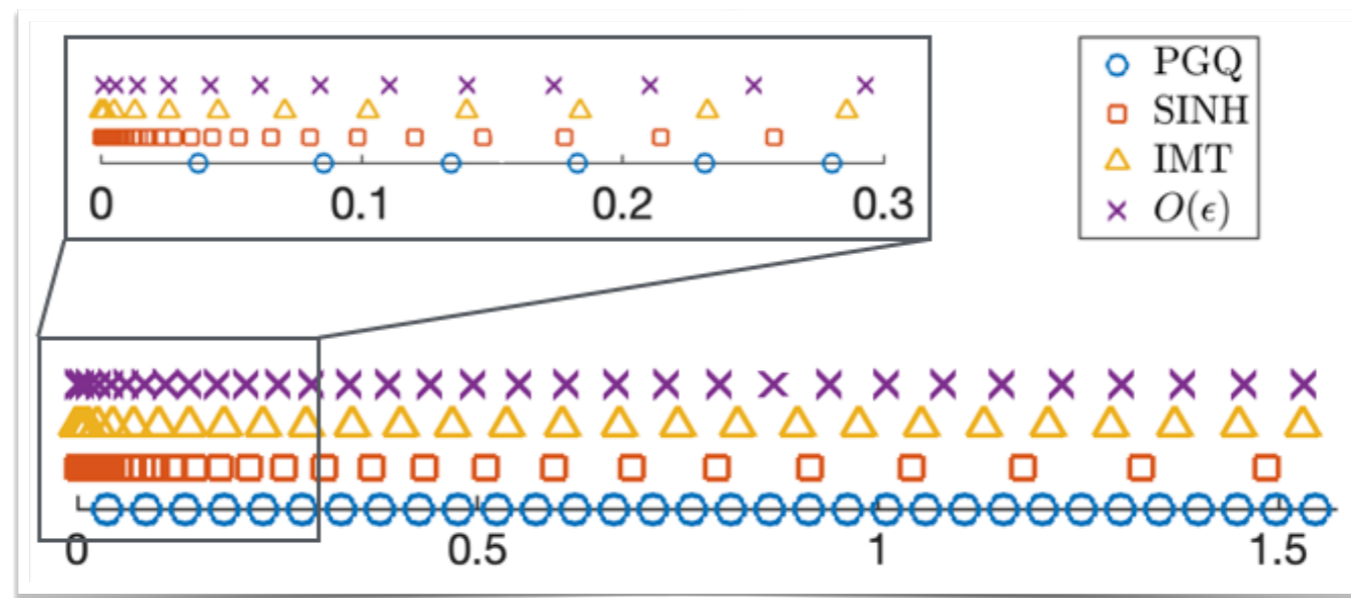
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✓ The quadrature clusters points near y^*

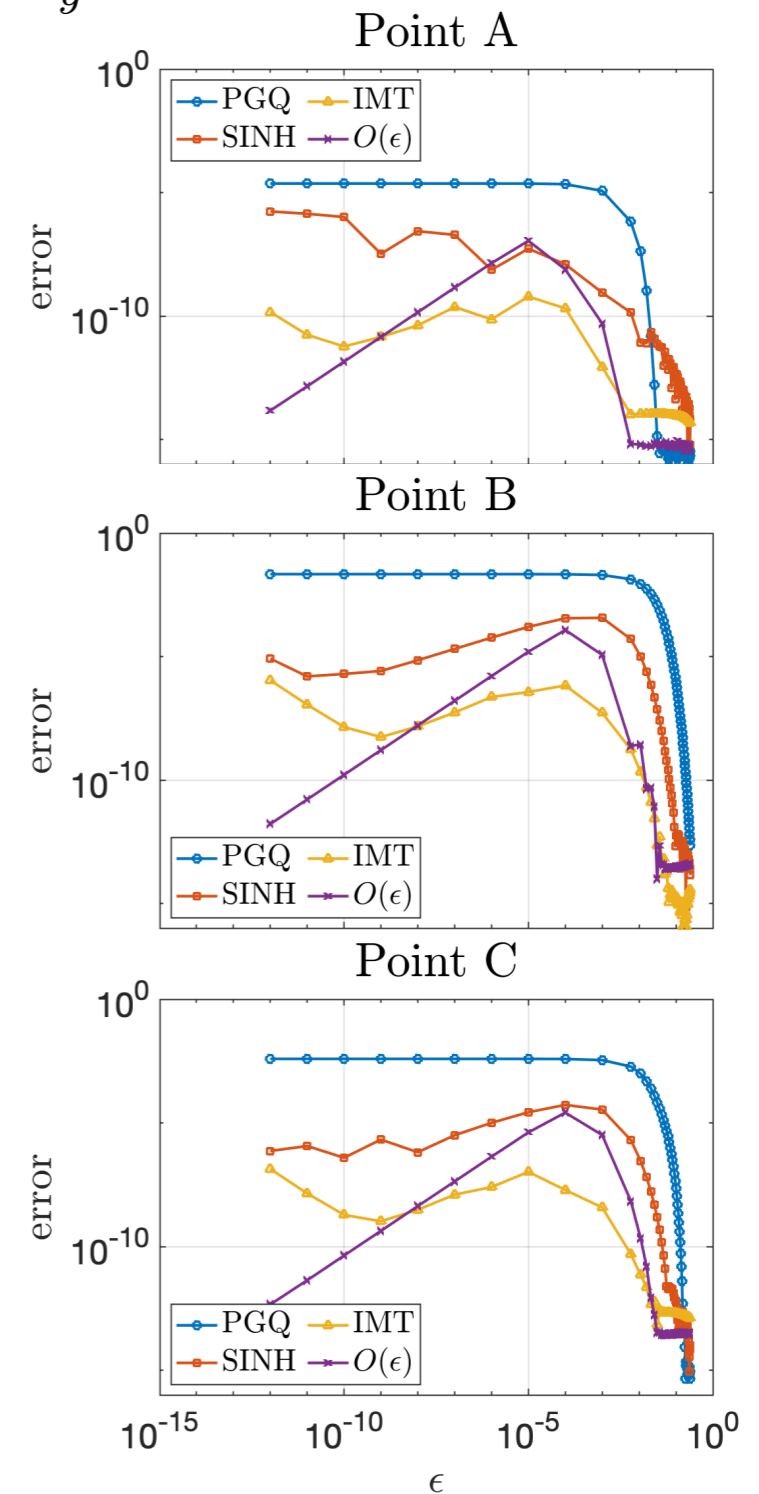
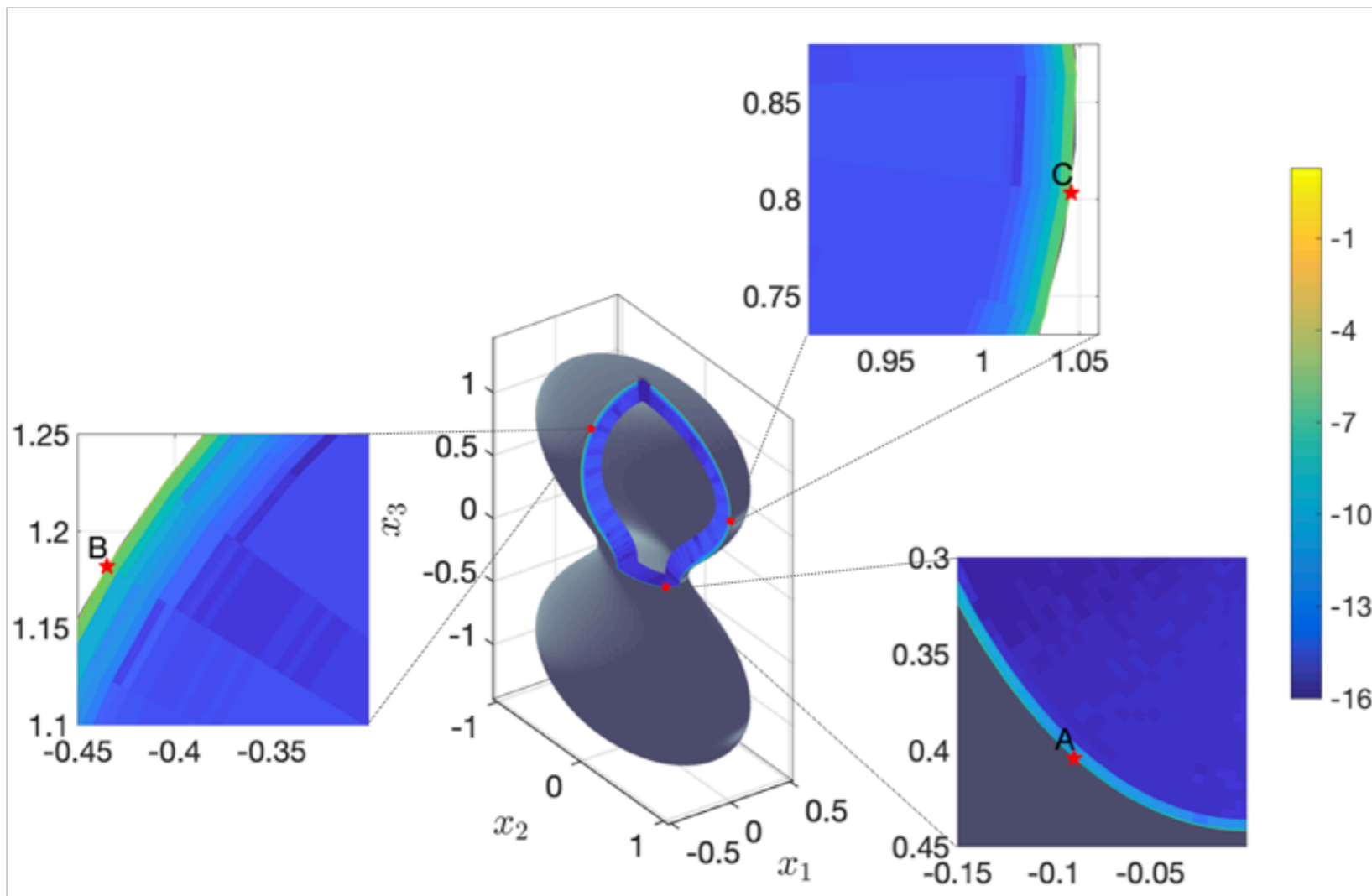
✓ The integrand is smoother at y^*

Numerical Results

$$u(x) = - \int_{\partial D} \partial_{n_y} G(x, y) u_{\text{ex}}(y) d\sigma_y + \int_{\partial D} G(x, y) \partial_n u_{\text{ex}}(y) d\sigma_y$$

Exact solution: $u_{\text{ex}}(x_1, x_2, x_3) = e^{x_3}(\sin x_1 + \sin x_2)$

Peanut shape with $N = 128$



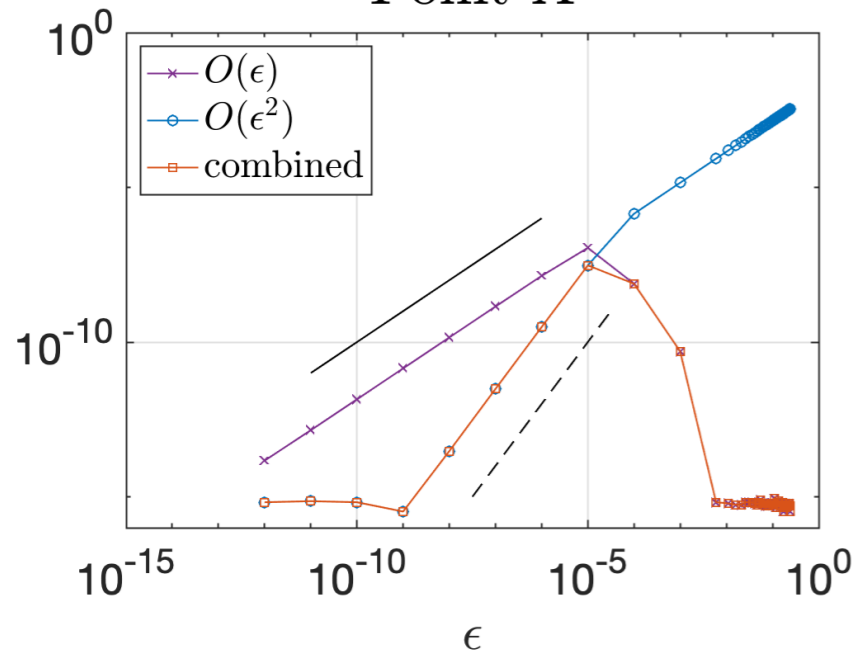
Extension to $O(\epsilon^2)$

Peanut shape with $N = 100$

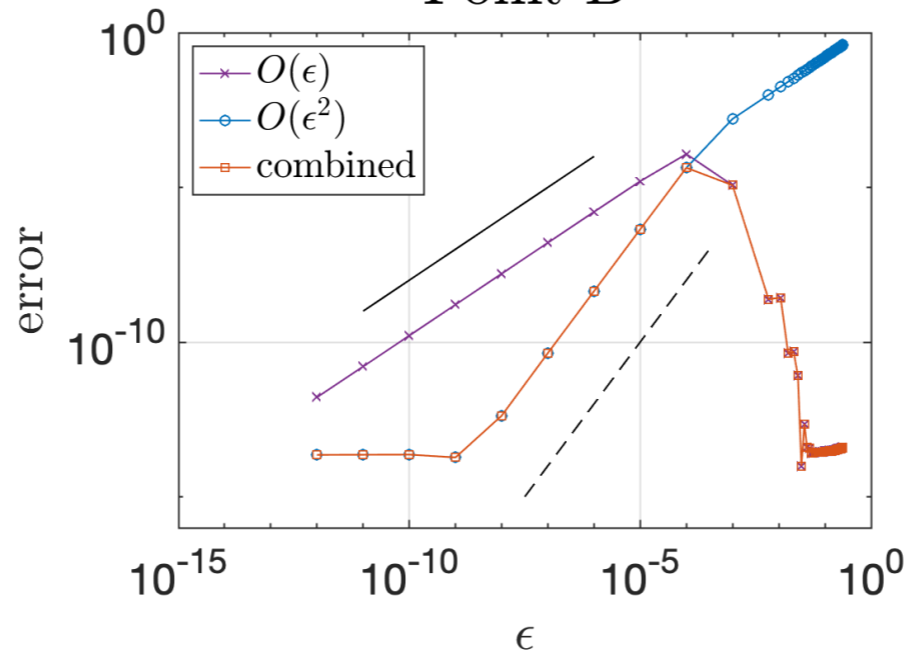
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Use asymptotic approximation of the single-layer potential

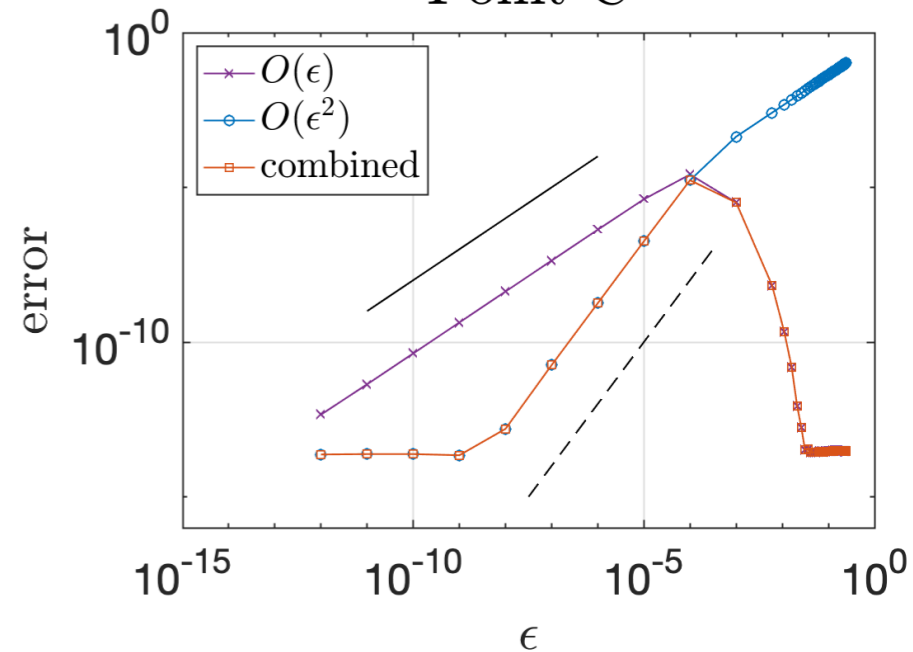
Point A



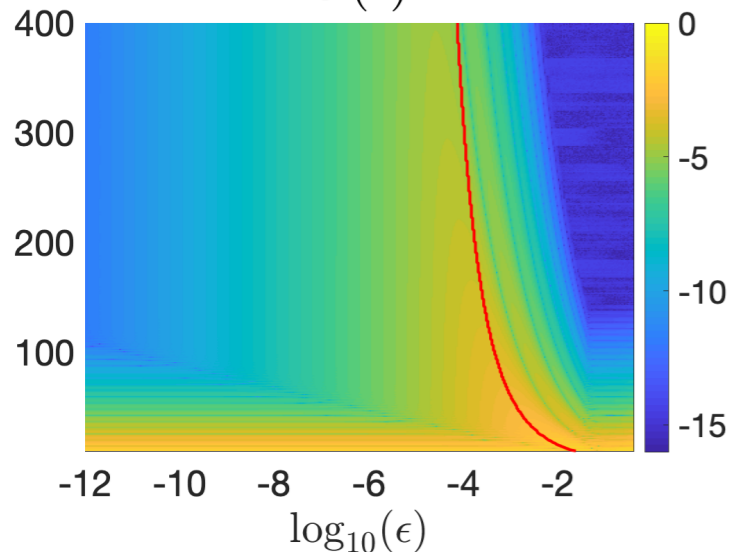
Point B



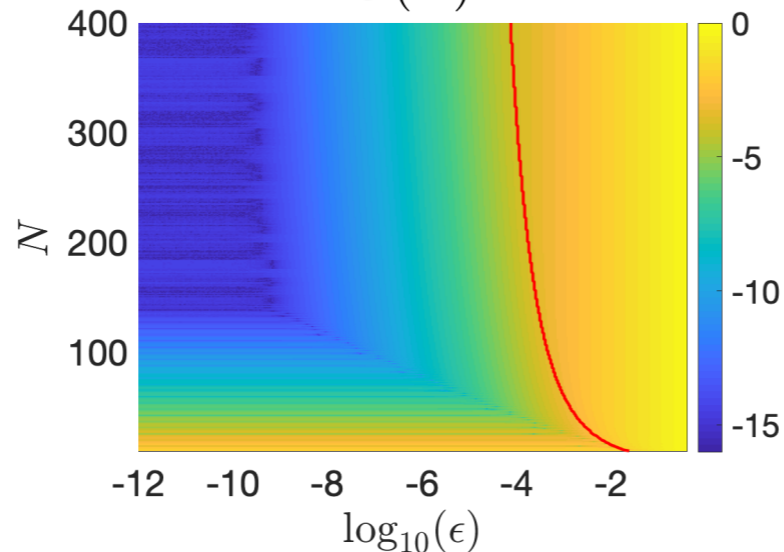
Point C



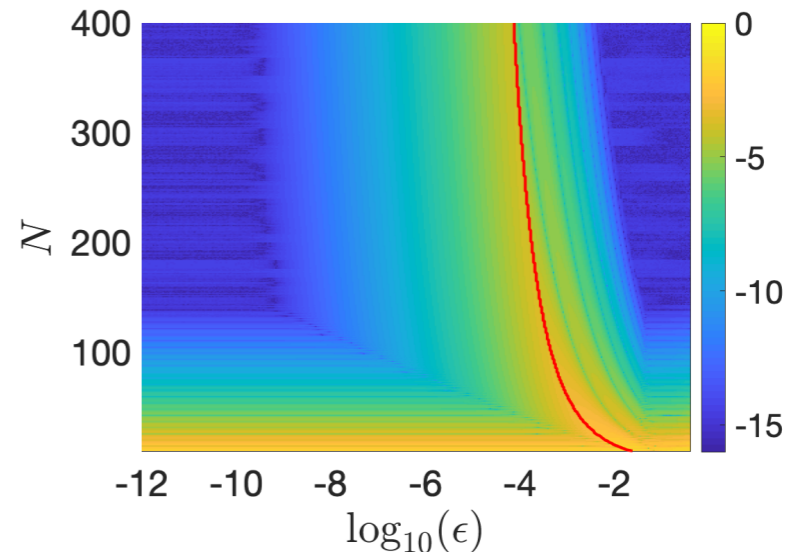
$O(\epsilon)$



$O(\epsilon^2)$



combined



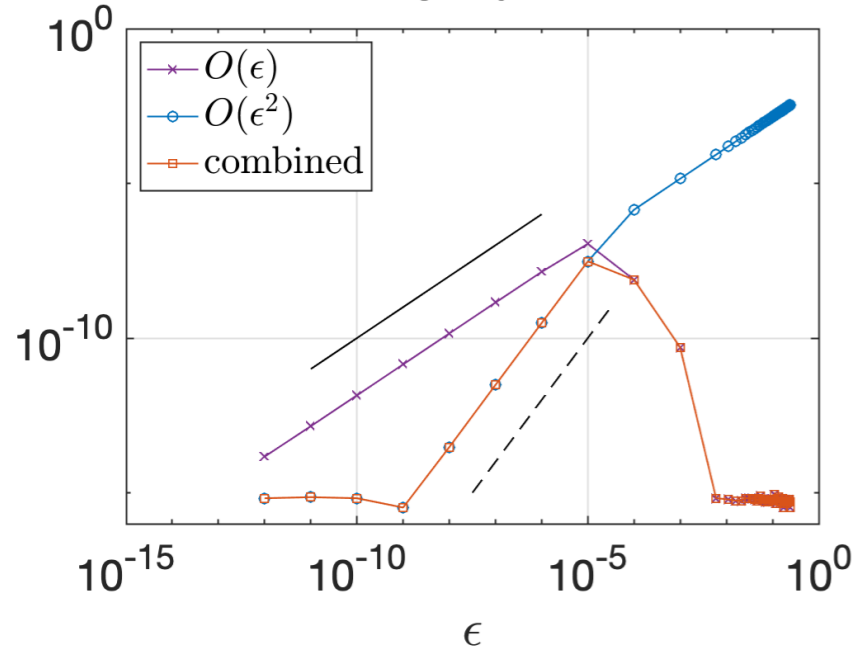
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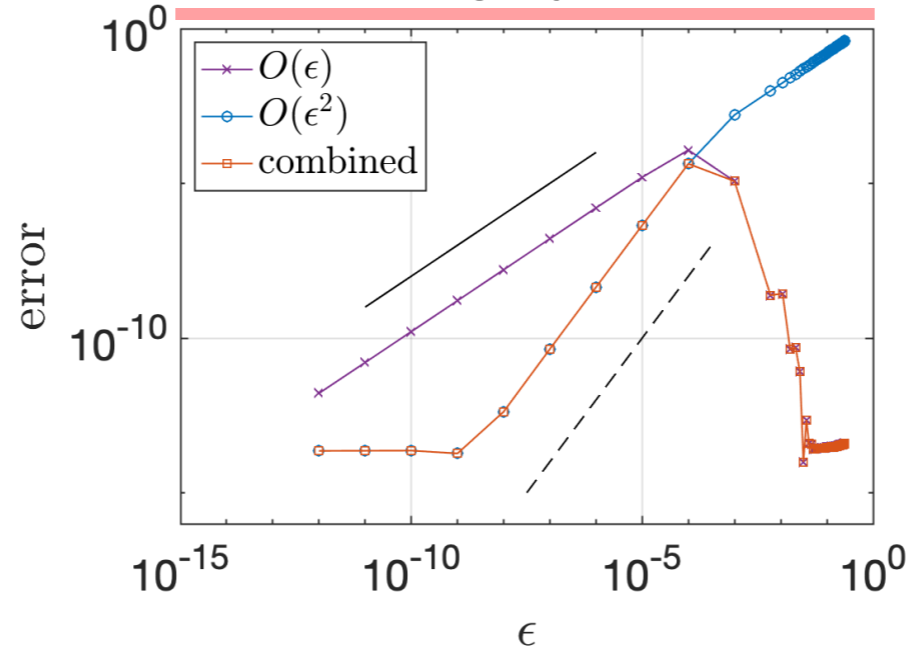
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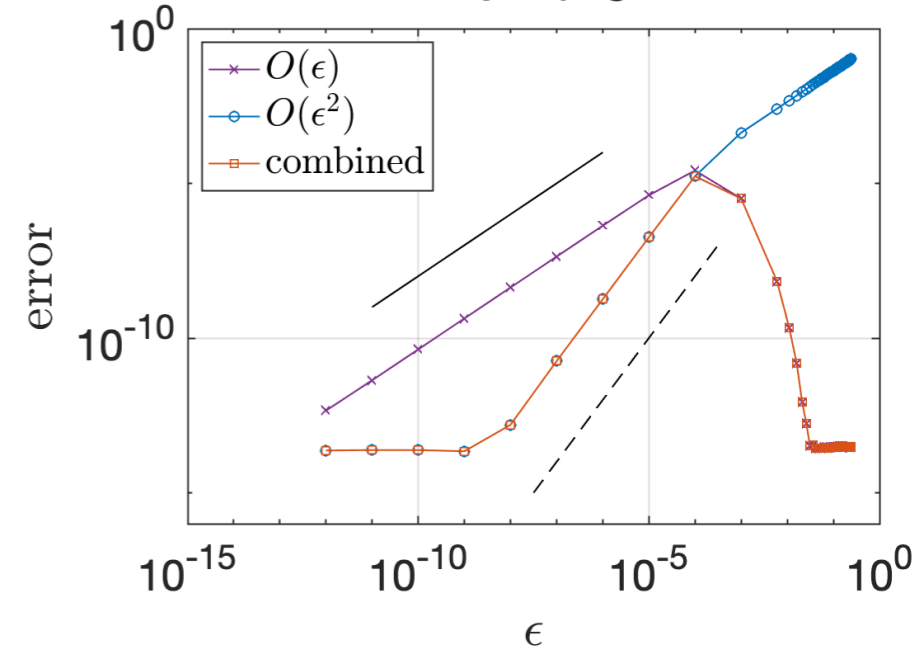
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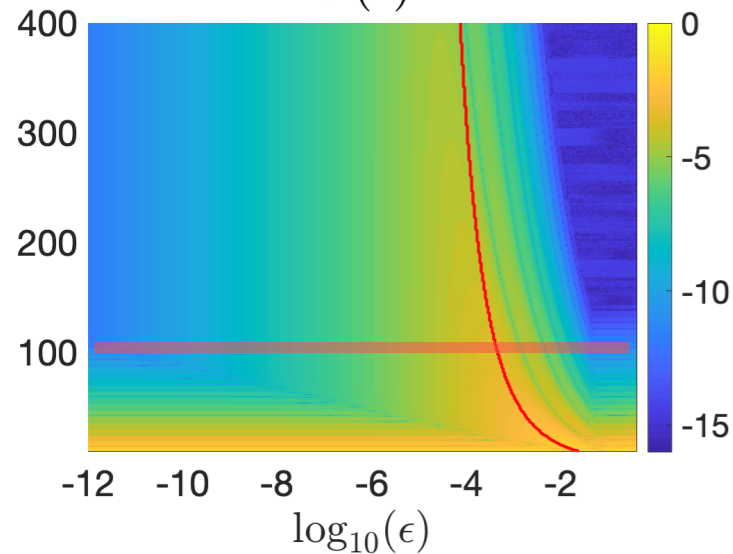
Point B



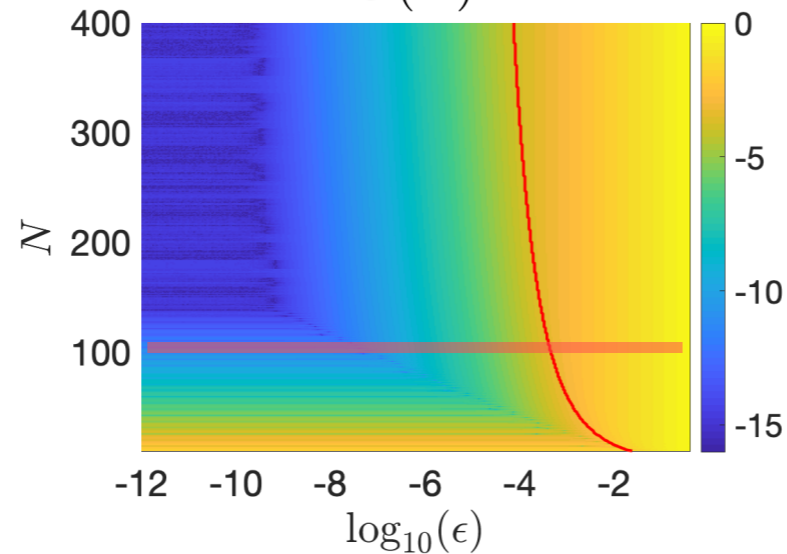
Point C



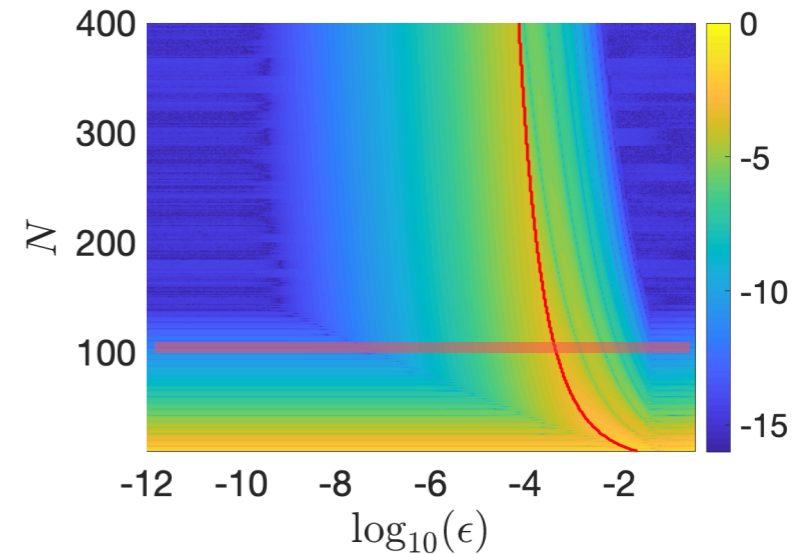
$O(\epsilon)$



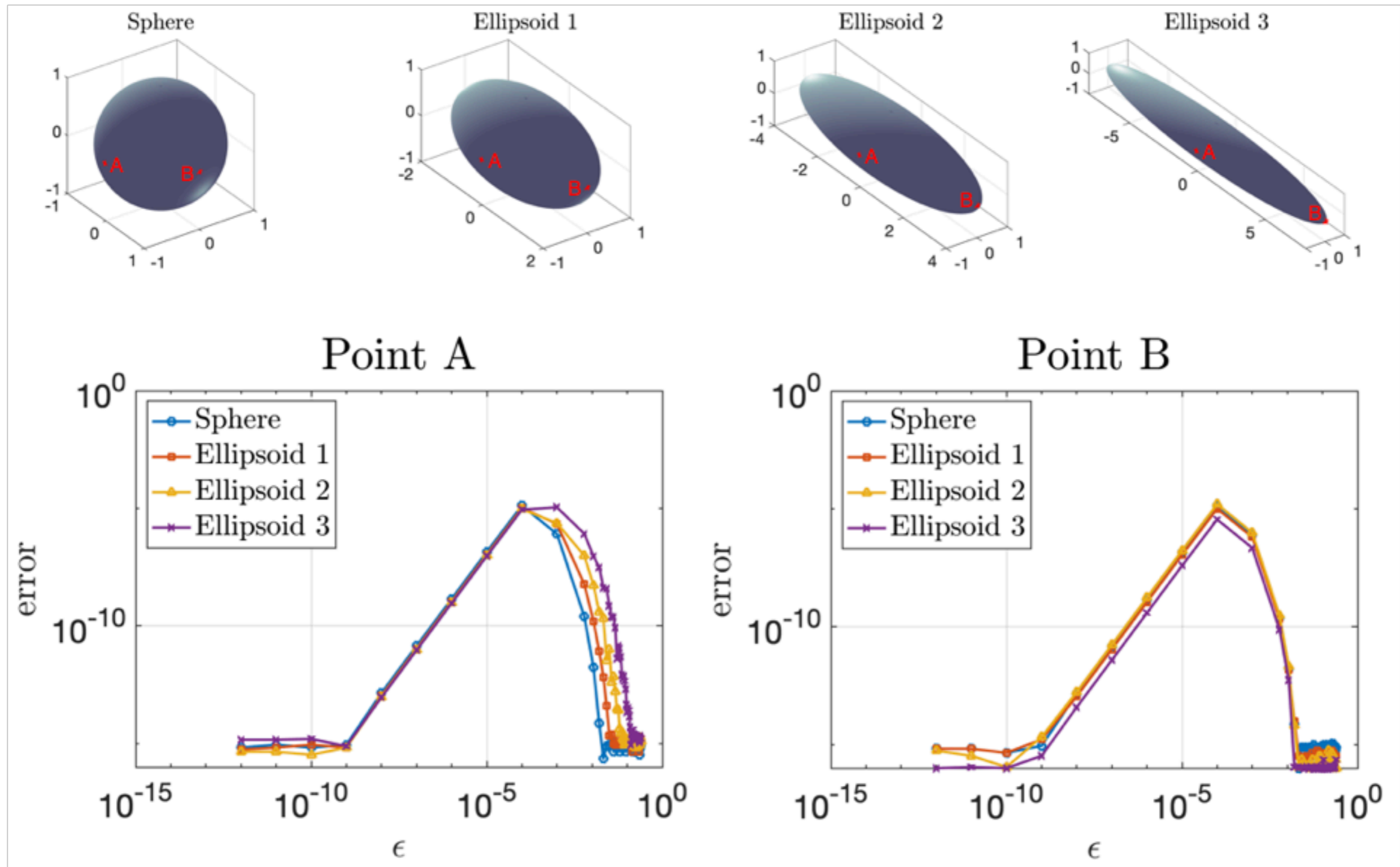
$O(\epsilon^2)$



combined



Effect of curvature



Outline

- ❖ The close evaluation problem
- ❖ Quadrature based on asymptotic methods
- ❖ Modified representations
- ❖ Conclusion

Modified representations

Previously we made use of Gauss' law to help attenuate the nearly singular behavior.

$$\int_{\partial D} \partial_n G(x, y) [u(y) - u(y^*)] d\sigma_y + u(y^*) \int_{\partial D} \partial_n G(x, y) d\sigma_y$$

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Representation formula gives us:

$$\int_{\partial D} \partial_{n_y} G(x, y) u(y) d\sigma_y - \int_{\partial D} G(x, y) \partial_n u(y) d\sigma_y = \begin{cases} -u(x) & x \in D \\ -\frac{1}{2}u(x) & x \in \partial D \\ 0 & x \in \mathbb{R}^3 \setminus \bar{D} \end{cases}$$

G : fundamental solution (Laplace or Helmholtz)

u : solution (of Laplace or Helmholtz) in D

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↓

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Modified representations

$$\text{Given } \int_{\partial D} \partial_{n_y} G(x, y) u^{\text{sol}}(y) d\sigma_y - \int_{\partial D} G(x, y) \partial_n u^{\text{sol}}(y) d\sigma_y = \begin{cases} -u^{\text{sol}}(x) & x \in D \\ -\frac{1}{2} u^{\text{sol}}(x) & x \in \partial D \\ 0 & x \in \mathbb{R}^3 \setminus \bar{D} \end{cases}$$

G : fundamental solution (Laplace or Helmholtz)

u^{sol} : solution (of Laplace or Helmholtz) in D

One can **modify the single-layer potential** $\int_{\partial D} G(x, y) \rho(y) d\sigma_y$

$$\int_{\partial D} G(x, y) \rho(y) [1 - \partial_n u^{\text{sol}}(y)] d\sigma_y + \int_{\partial D} G(x, y) \rho(y) \partial_n u^{\text{sol}}(y) d\sigma_y$$

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Exterior Neumann Laplace problem on a sphere ($N=16$)

Good resolution of ρ

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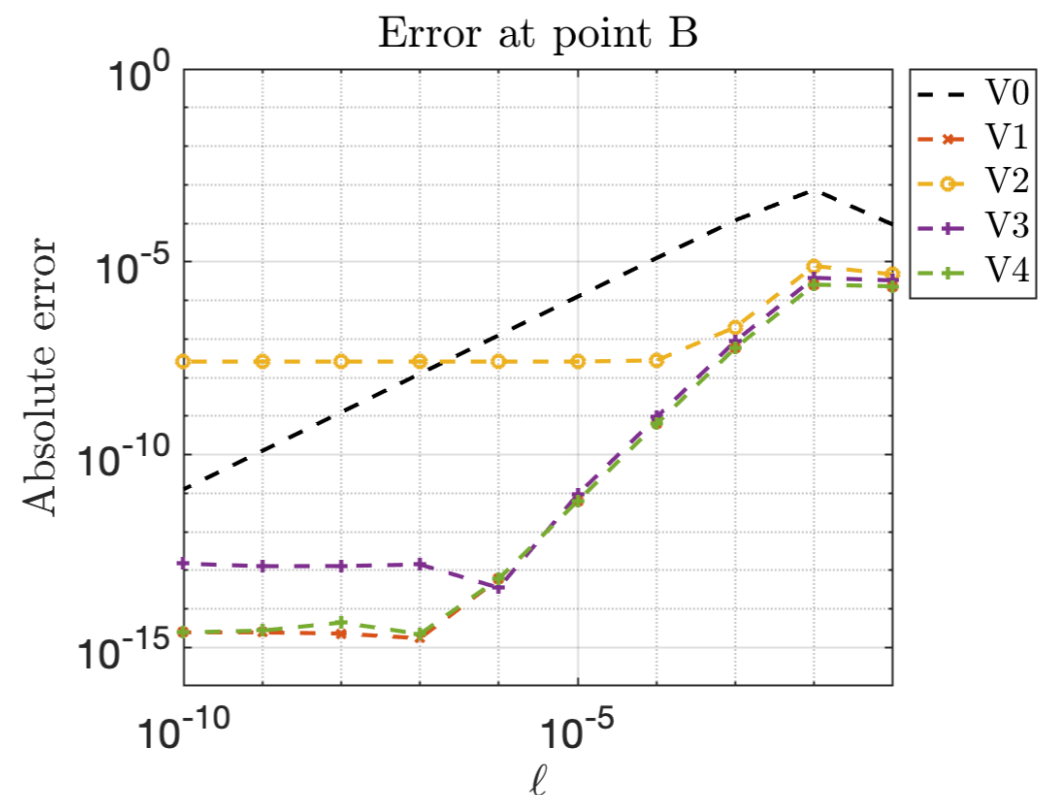
$$\partial_n u^{\text{sol}}(y^*) = 1$$

$$V1: u^{\text{sol}}(y) = n^* \cdot y$$

$$V2: u^{\text{sol}}(y) = 4\pi G(y, y^* + n^*)$$

$$V3: u^{\text{sol}}(y) = \frac{1}{2} \frac{y_1^2 - y_2^2}{n_{y^*,1} y_1^* - n_{y^*,2} y_2^*}$$

$$V4: u^{\text{sol}}(y) = \frac{(y_1-5)(y_2-5)}{n_{y^*,1}(y_2^*-5) + n_{y^*,2}(y_1^*-5)}$$



Results for Helmholtz layer potentials

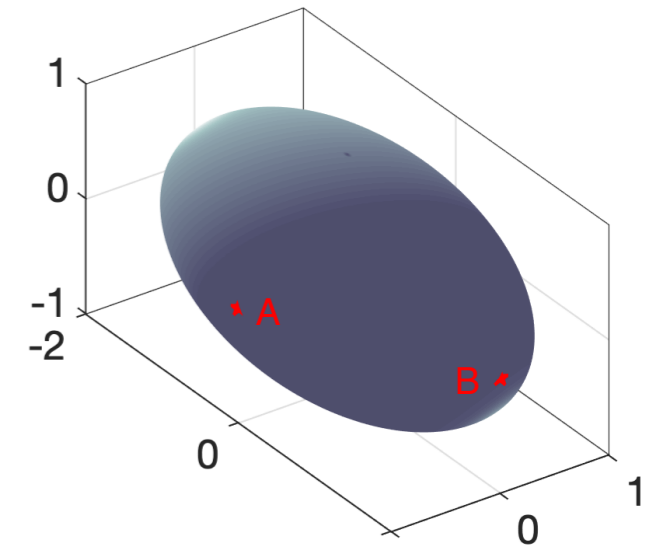
Sound-soft scattering problem on an ellipsoid ($k = 5$)

Limited resolution of μ (10^{-7})

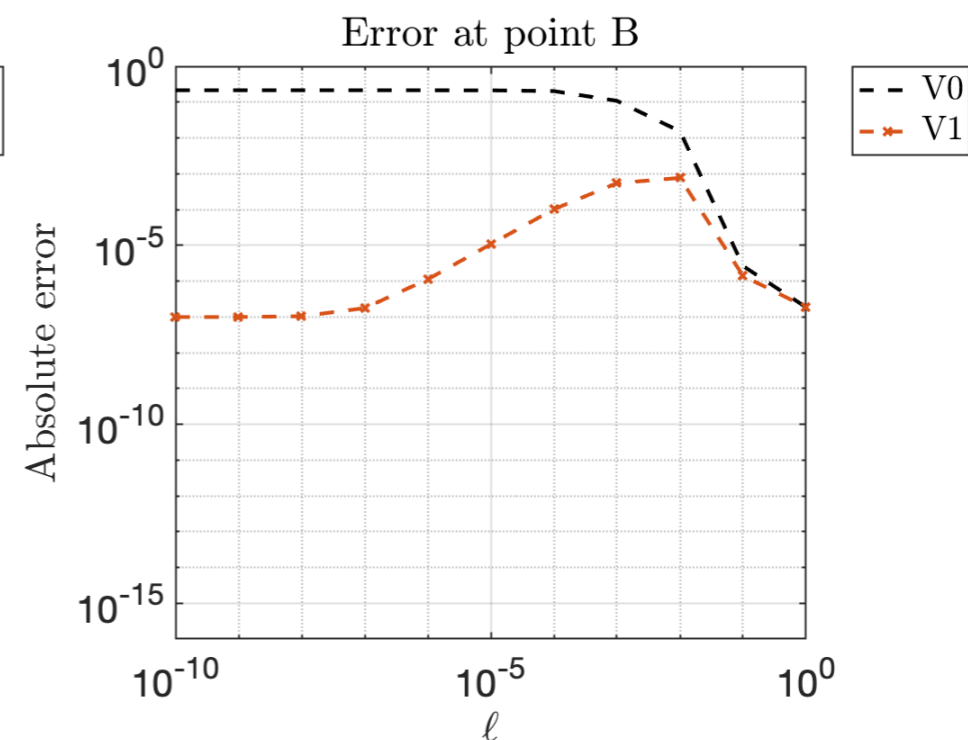
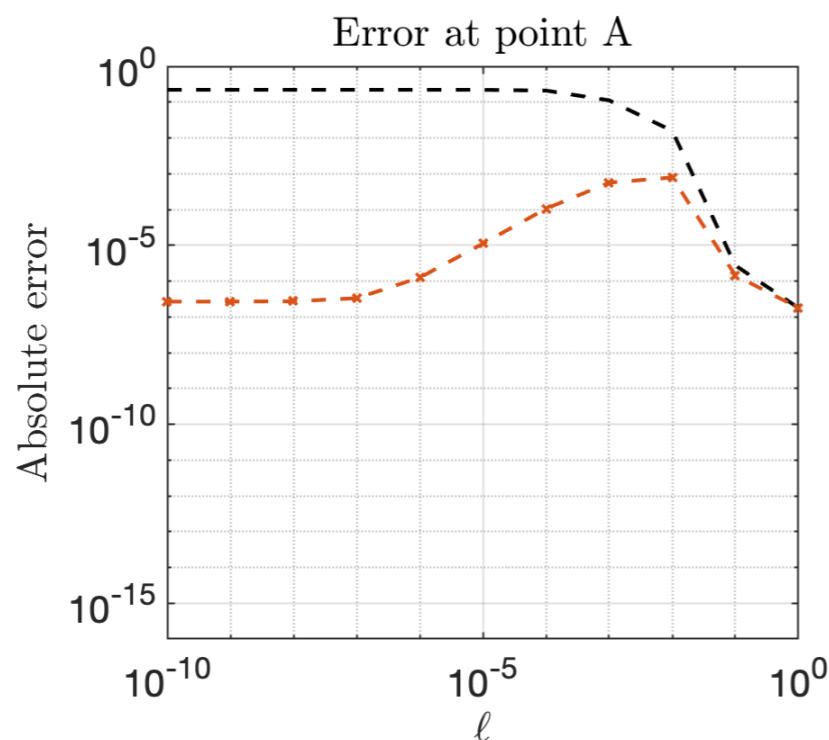
Representation V0:
$$u(x) = \int_{\partial D} [\partial_{n_y} G(x, y) - ikG(x, y)] \mu(y) d\sigma_y$$

Representation V1:
$$u^{\text{sol}}(y) = e^{ikn^* \cdot (y - y^*)}$$

$$u(x) = \int_{\partial D} [\partial_{n_y} G(x, y) - \partial_n u^{\text{sol}}(y) G(x, y)] [\mu(y) - \mu(y^*)] d\sigma_y + \int_{\partial D} G(x, y) [\partial_n u^{\text{sol}}(y) - ik] \mu(y) d\sigma_y + \mu(y^*) \int_{\partial D} \partial_{n_y} G(x, y) [1 - u^{\text{sol}}(y)] d\sigma_y$$



$N = 32$



Outline

- ❖ The close evaluation problem
- ❖ Quadrature based on asymptotic methods
- ❖ Modified representations
- ❖ Conclusion

Summary

Due to sharply peaked behavior of layer potentials' kernel, one makes an $O(1)$ error for close evaluation.

Local analysis provides valuable insights to design methods that naturally address the nearly singular behavior

When one has limited density resolution, modified representations help reduce the error (for free)

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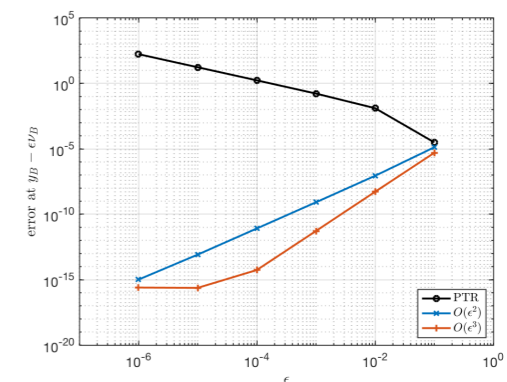
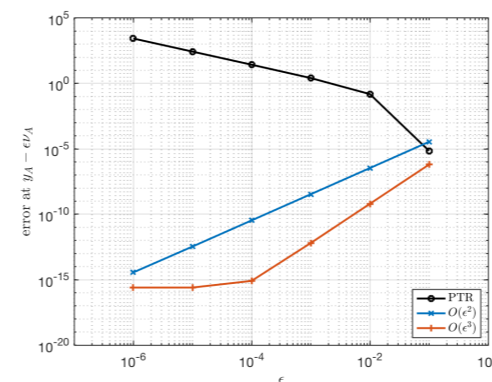
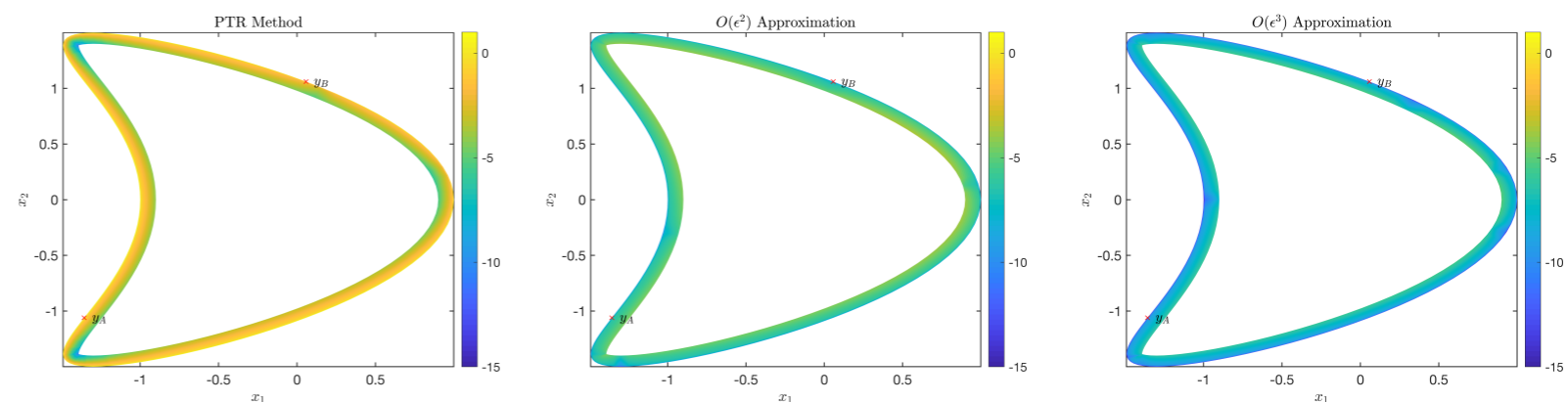
Kernel / singularity subtraction

Asymptotic approximations (2D)



Perez-Arancibia (2018)

Carvalho, Khatri, Kim (2020)



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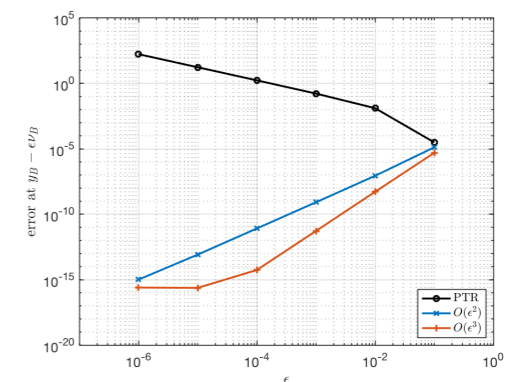
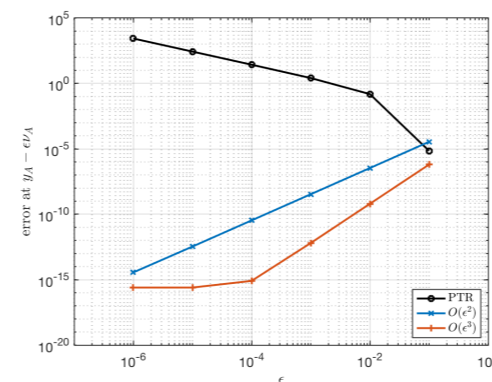
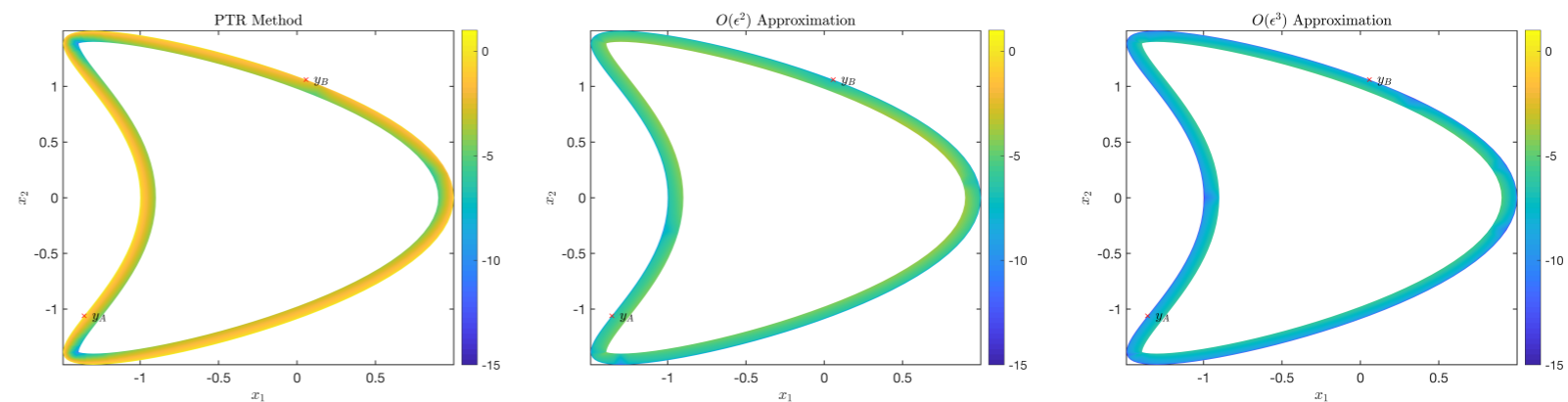
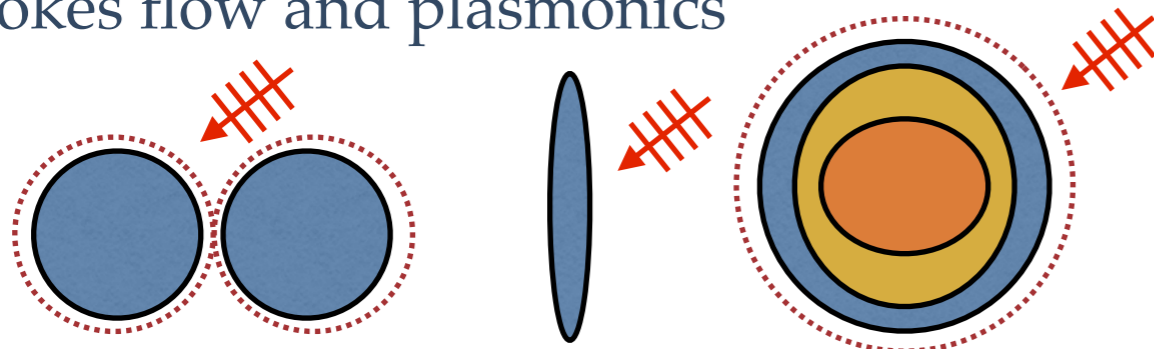


Perez-Arancibia (2018)

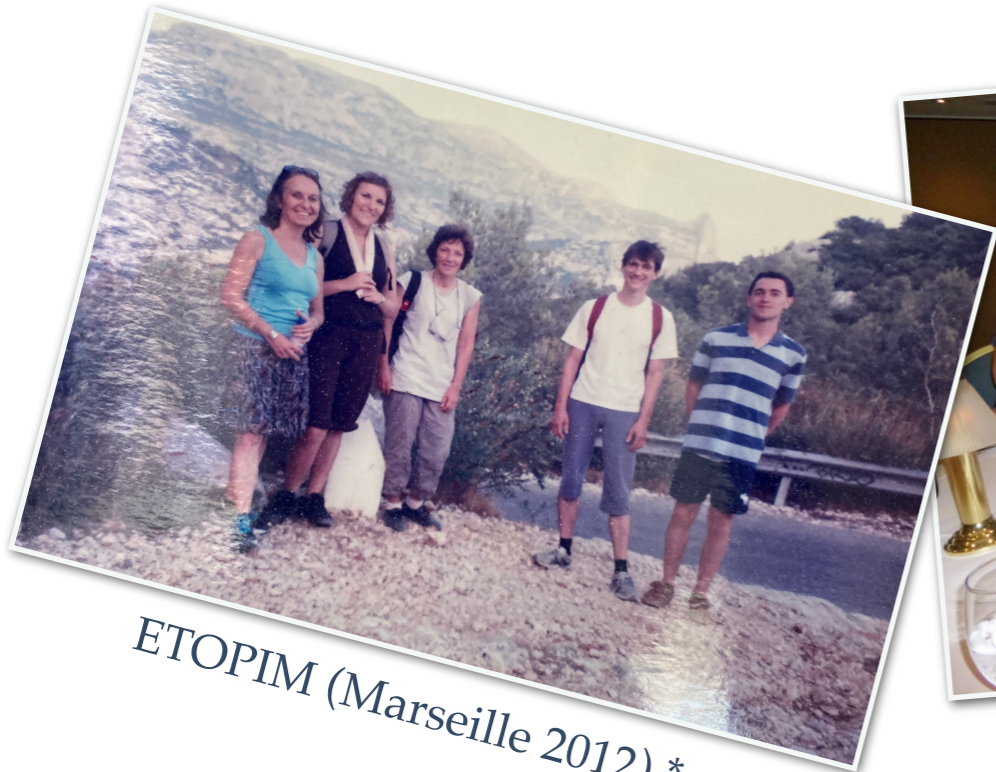
Carvalho, Khatri, Kim (2020)

Perspectives:

Stokes flow and plasmonics



Thank you for your attention.



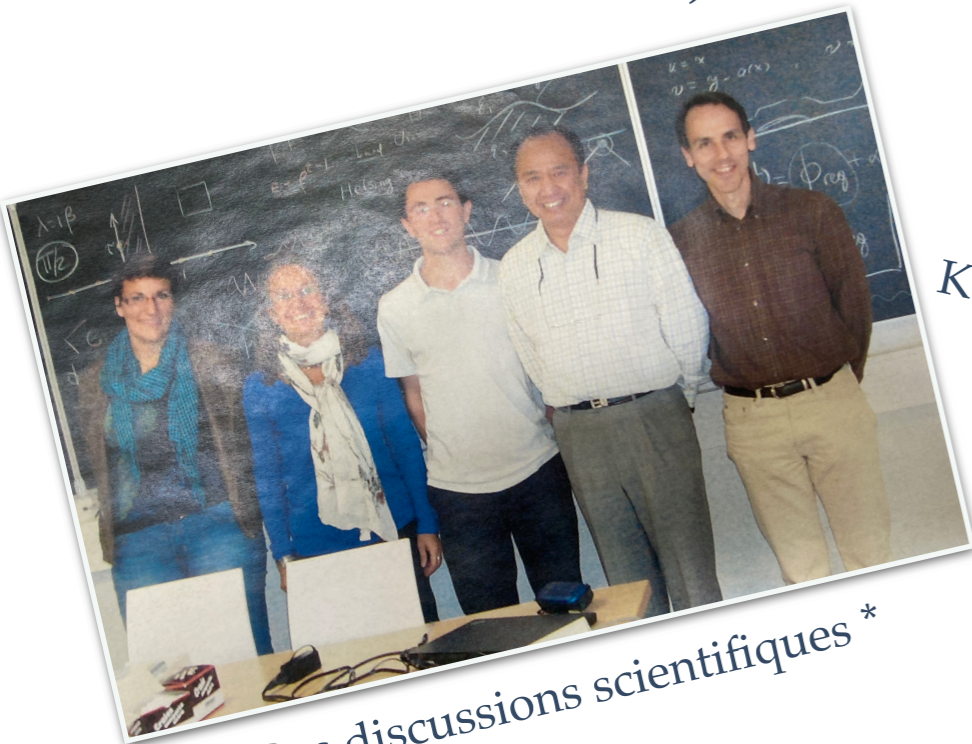
ETOPIM (Marseille 2012) *



CIRM (Marseille, 2022)



KOZ WAVES (New Castle/Sydney, 2014)



Des discussions scientifiques *



L'équipe de choc de l'UMA

* Photos de photos du bureau d'A.-S. Merci Lucas !