

# UNE MÉTHODE DE RECONSTRUCTION POUR UN PROBLÈME INVERSE EN GRAVIMÉTRIE

**Karim Ramdani**

avec **Anthony Gerber-Roth & Alexandre Munnier**



**JO des POÈTES - RELAIS 4x60**  
**17-19 avril 2024**

# DÉAMBULATION ONDULATOIRE

**Karim Ramdani**



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# Susceptibility: could age be the explanatory variable?

**Willard Oodendijk<sup>1\*</sup>, Michaël Rochoy<sup>2</sup>, Valentin Ruggeri<sup>3</sup>, Florian Cova<sup>4</sup>,  
Didier Lembrouille<sup>5</sup>, Sylvano Trottinetta<sup>6</sup>, Otter F. Hantome<sup>7</sup>,  
Nemo Macron<sup>8</sup> and Manis Javanica<sup>9</sup>**

<sup>1</sup>*Belgian Institute of Technology and Education (BITE), Couillet, Belgium.*

<sup>2</sup>*General Practitioner and Independent Seeker of Science, Ankh, Morpork, France.*

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<sup>7</sup>*University of Melon, Melon, France.*

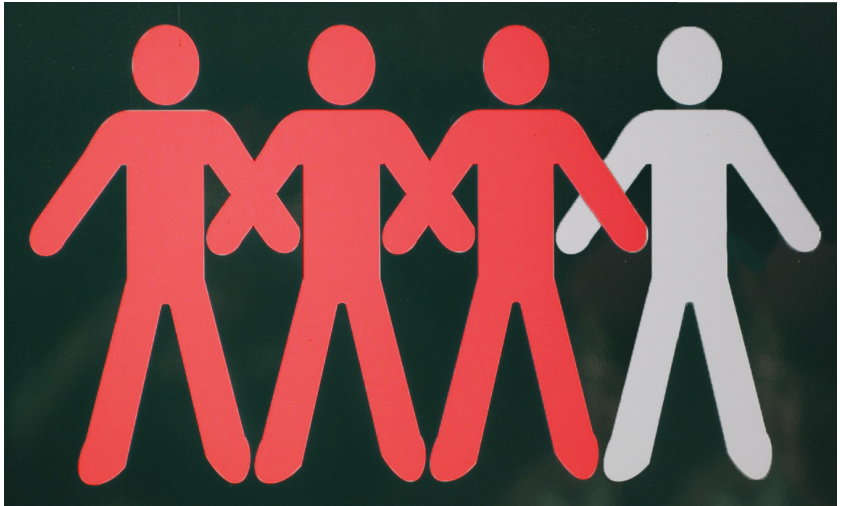
<sup>8</sup>*Palais de l'Élysée, Paris, France.*

<sup>9</sup>*Institute of Chiropteran Studies, East Timor.*

### ***Authors' contributions***

*This work was carried out in collaboration among all authors. Author WO launched idea on Twitter, added some sentences, submitted the paper, corresponded with the kind publisher. Author MR*





Après 60 ans, la susceptibilité augmente de 60% dans 3/4 des cas !

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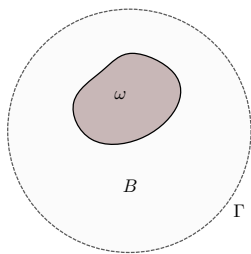
# OUTLINE

- 1 Problem formulation
- 2 Reconstruction method
  - Step 1 : Constructing quadrature formula
  - Step 2 : Constructing quadrature domains
- 3 Numerical results

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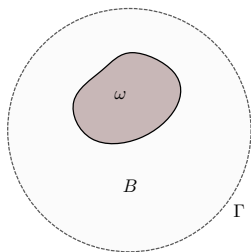
# The inverse problem



The gravitational potential generated by a uniform mass distribution in  $\omega$  satisfies

$$\begin{cases} -\Delta u_\omega &= \mathbb{1}_\omega & \text{on } \mathbb{R}^2, \\ u_\omega &= \mathcal{O}(\ln |x|) & |x| \rightarrow +\infty. \end{cases}$$

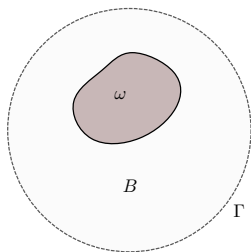
# The inverse problem



Obviously, this gravitational potential is given by

$$u_{\omega}(x) = \int_{\omega} G(x - y) \, dy = -(2\pi)^{-1} \int_{\omega} \ln |x - y| \, dy$$

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## Inverse problem

Reconstruct  $\omega$  given  $u_{\omega}$  (or  $\nabla u_{\omega}$ ) on  $\Gamma$ .

# Uniqueness

## Theorem [Isakov, 1990]

Let  $\omega_1, \omega_2$  be both

- star-shaped with respect to their centers of gravity

or

- convex in one direction

then  $\nabla u_{\omega_1} = \nabla u_{\omega_2}$  on  $\Gamma$  implies  $\omega_1 = \omega_2$ .

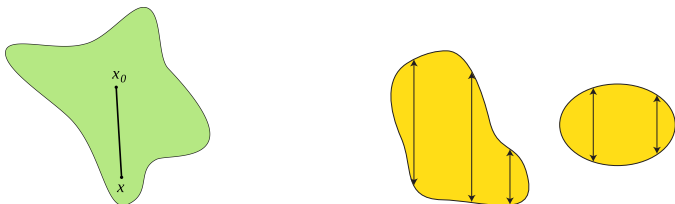


Figure – Star-shaped (left) and convex in one direction (right) domains.



# Stability

Given  $M, \rho_-, \rho_+, \alpha > 0$ , let  $\mathcal{U}$  be the set of star-shaped domains  $\omega$  described in polar coordinates by

$$\omega = \{x = (r, \theta) \in \mathbb{R}^2 \mid r < \rho(\theta), \quad \theta \in [0, 2\pi]\},$$

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**Theorem, [Isakov, 1990]**

There exists  $C > 0$  such that for all  $\omega_1, \omega_2 \in \mathcal{U}$ ,

$$\|\nabla u_{\omega_1} - \nabla u_{\omega_2}\|_{L^\infty(\Gamma)} < \varepsilon \Rightarrow \|\rho_1 - \rho_2\|_{L^\infty([0, 2\pi])} \leq C |\ln \varepsilon|^{-\frac{1}{C}}.$$

## Idea of the method

Our goal is to construct a sequence of domains  $(\omega_N)_N$  satisfying the following (asymptotical) gravi-equivalence property :

$$\|\nabla u_\omega - \nabla u_{\omega_N}\|_{L^\infty(\Gamma)} \xrightarrow{N \rightarrow +\infty} 0.$$

## Idea of the method

By Green's formula, knowing  $u_\omega$  and  $\partial_n u_\omega := \nabla u_\omega \cdot n$  over  $\Gamma$ , we can write for all function  $v$  :

$$\int_{\Gamma} (u_\omega \partial_n v - \partial_n u_\omega v) = \int_{\Omega_0} \Delta v u_\omega - \int_{\Omega_0} \Delta u_\omega v.$$

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If we choose  $v$  harmonic, we can compute from the measurements

$$\int_{\omega} v = \int_{\Gamma} u_\omega \partial_n v - \partial_n u_\omega v.$$

Using these harmonic moments as new measurements, we will construct the domains  $\omega_N$  such that

$$\forall 0 \leq m \leq N, \quad \int_{\omega} z^m = \int_{\omega_N} z^m,$$

and these equalities will ensure the gravi-equivalence property.

## Related problems

### Calderón's inverse problem for highly conducting inclusions

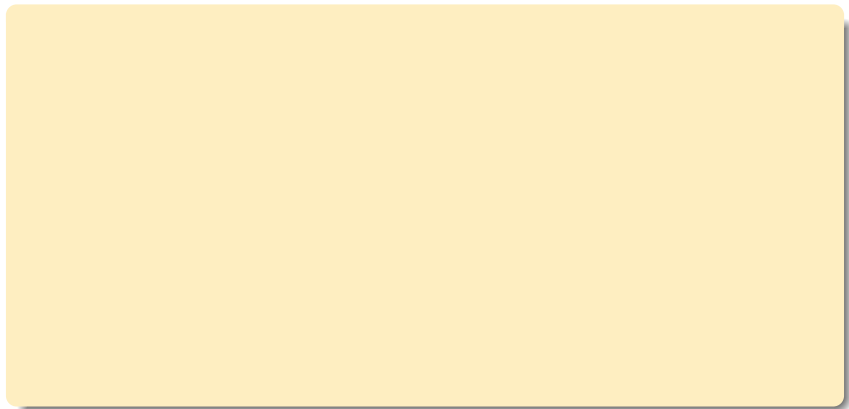
$$\begin{cases} -\Delta u = 0 & \text{in } \Omega_0 \setminus \bar{\omega}, \\ u = f & \text{on } \Gamma, \\ u = c & \text{on } \partial\omega, \end{cases}$$

where the constant  $c$  is such that :  $\int_{\partial\omega} \frac{\partial u}{\partial n} = 0$ .

Recovering the shape of a highly conducting inclusion  $\omega$  from the knowledge of the DtN operator can also be formulated as a shape from moments inverse problem. See for instance the contributions of AMMARI *et al.* on **Generalized Polarisation Tensors** and the two papers by MUNNIER & R. (2017, 2018).



## Related problems



# Métrologie

Multiples	Symboles	Rapport à l'unité de mesure	Sous-multiples	Symboles	Rapport à l'unité de mesure
quetta-	Q	$10^{30}$	déci-	d	$10^{-1}$
ronna-	R	$10^{27}$	centi-	c	$10^{-2}$
yotta-	Y	$10^{24}$	milli-	m	$10^{-3}$
zetta-	Z	$10^{21}$	micro-	$\mu$	$10^{-6}$
exa-	E	$10^{18}$	nano-	n	$10^{-9}$
péta-	P	$10^{15}$	<b>pico-</b>	p	$10^{-12}$
téra-	T	$10^{12}$	femto-	f	$10^{-15}$
giga-	G	$10^9$	atto-	a	$10^{-18}$
méga-	M	$10^6$	zepto-	z	$10^{-21}$
kilo-	k	$10^3$	yocto-	y	$10^{-24}$
hecto-	h	$10^2$	ronto-	r	$10^{-27}$
déca-	da	$10^1$	quecto-	q	$10^{-30}$

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## Reconstruction method

**Step 1** Given  $N \geq 1$ , find the (complex) weights  $c_1, \dots, c_N$  and the nodes  $z_1, \dots, z_N$  such that :

$$\forall 0 \leq m \leq 2N - 1, \quad \int_{\omega} z^m = \sum_{n=1}^N c_n z_n^m.$$

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**Step 2** Knowing  $c_1, \dots, c_N$  and  $z_1, \dots, z_N$ , construct (if possible)  $\omega_N$  such that

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$$\forall m \geq 0, \quad \int_{\omega_N} z^m = \sum_{n=1}^N c_n z_n^m.$$

In particular, we will have

$$\forall 0 \leq m \leq 2N - 1, \quad \int_{\omega} z^m = \int_{\omega_N} z^m.$$

## Step 1 : Quadrature formula

Knowing

$$\tau_m = \int_{\omega} z^m,$$

we want to solve the nonlinear system of  $2N$  equations and  $2N$  unknowns (with pairwise distinct nodes)

$$\forall 0 \leq m \leq 2N - 1, \quad \sum_{n=1}^N c_n z_n^m = \tau_m.$$

Such a system is known as a Prony's system, and appears for instance in the study of Padé's approximants, signal processing, error correction codes (see [BATENKOV and YOMDIN, 2013]).

# Prony's method

Set

$$\mathbb{H}_0 = \begin{pmatrix} \tau_0 & \tau_1 & \cdots & \tau_{N-1} \\ \tau_1 & \tau_2 & \cdots & \tau_N \\ \vdots & \vdots & \ddots & \vdots \\ \tau_{N-1} & \tau_N & \cdots & \tau_{2N-2} \end{pmatrix}$$

and

$$P(z) = \left| \begin{array}{cccc|c} & & & & \tau_N \\ & & & & \tau_{N+1} \\ & & & & \vdots \\ & & & & \tau_{2N-1} \\ \hline 1 & z & \cdots & z^{N-1} & z^N \end{array} \right|.$$



# Prony's method

## Theorem

*The Prony's system*

$$\forall 0 \leq m \leq 2N - 1, \quad \sum_{n=1}^N c_n z_n^m = \tau_m.$$

*admits a solution if and only if the polynomial  $P$  admits  $N$  simple roots. In this case, this solution is unique and the nodes  $z_1, \dots, z_N$  are the roots of  $P$ .*

# Prony's method

Knowing  $(z_n)$ , we can compute  $(c_n)$  by solving the Vandermonde system

$$\mathbb{V}(z)\mathbf{c} = \boldsymbol{\tau}.$$

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$$\mathbb{V}(z)\mathbf{c} = \boldsymbol{\tau}.$$

Following [GOLUB, MILANFAR AND VARAH, 1999],  $(z_n)$  are the eigenvalues of the generalized eigenvalue problem

$$\mathbb{H}_1\mathbf{x} = z\mathbb{H}_0\mathbf{x},$$

with

$$\mathbb{H}_0 = \begin{pmatrix} \tau_0 & \tau_1 & \cdots & \tau_{N-1} \\ \tau_1 & \tau_2 & \cdots & \tau_N \\ \vdots & \vdots & \ddots & \vdots \\ \tau_{N-1} & \tau_N & \cdots & \tau_{2N-2} \end{pmatrix} \quad \mathbb{H}_1 = \begin{pmatrix} \tau_1 & \tau_2 & \cdots & \tau_N \\ \tau_2 & \tau_3 & \cdots & \tau_{N+1} \\ \vdots & \vdots & \ddots & \vdots \\ \tau_N & \tau_{N+1} & \cdots & \tau_{2N-1} \end{pmatrix}.$$

## Prony's method

### Remark

*Solving Prony's system involves 2 ill-conditioned problems : a generalized eigenvalue problem with Hankel matrices and a Vandermonde linear system. To improve the conditioning, we solve a modified generalized eigenvalue problem*

$$\mathbf{H}_1 \mathbf{x} = z \mathbf{H}_0 \mathbf{x},$$

where

$$\mathbf{H}_\ell = \Phi \mathbf{H}_\ell \Phi^T, \quad \ell = 0, 1$$

and

$$\Phi = \text{Diag} \left[ \frac{1}{\hat{\rho}^n} \right]_{n=0}^{N-1}, \quad \hat{\rho} \sim |\tau_N|^{1/N}.$$

The unknown domain  $\omega$  can be covered by plotting the disks centered at the nodes  $z_k$  with radii  $\sqrt{|\operatorname{Re}(c_k)|}/\pi$ .

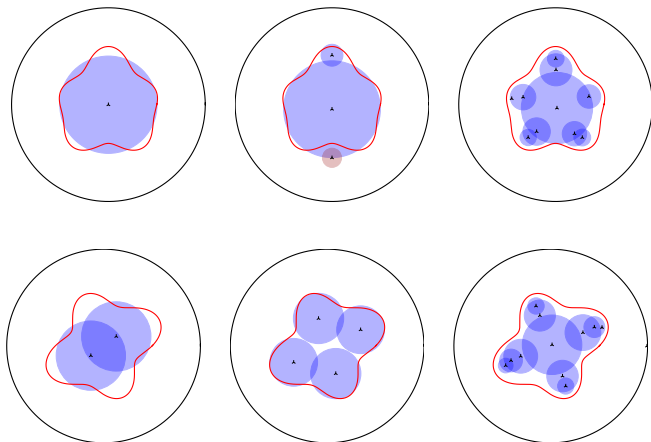


Figure – Covering the unknown domain by balls : two examples.

## Step 2 : Quadrature domains

Given the weights  $c_1, \dots, c_N$  and the nodes  $z_1, \dots, z_N$ , construct  $\omega_N$  such that

$$\forall m \geq 0, \quad \int_{\omega_N} z^m = \sum_{n=1}^N c_n z_n^m.$$

This leads to the notion of **quadrature domains**.

## Quadrature domains

### Definition

$\Omega$  is a **Harmonic Quadrature Domain (HQD)**<sup>1</sup> if there exists nodes  $(z_n)_{1 \leq n \leq N}$  in  $\Omega$  and (real) weights  $(c_n)_{1 \leq n \leq N}$  such that for all harmonic function  $v$  :

$$\int_{\Omega} v = \sum_{n=1}^N c_n v(z_n).$$

- 
1. "What is a quadrature domain?", B. GUSTAFSSON and H. S. SHAPIRO, 2005.

The **disk** is the simplest HQD, since for every harmonic function  $v$  in the disk  $B(a, r)$  (**mean value property**)

$$\int_{B(a,r)} v = \pi r^2 v(a).$$

## Quadrature domains

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- **Existence/Uniqueness of HQD ?** Existence is an open question. Uniqueness does not hold (AMEUR, HELMER AND TEL-LANDER, 2021).
- **Generalization** : The above definition of HQD can be extended to arbitrary measures (and not only atomic).

## Quadrature domains

### Definition

$\Omega$  is a **Sub-Harmonic Quadrature Domain (SHQD)** if there exists nodes  $(z_n)_{1 \leq n \leq N}$  in  $\Omega$  and (positive) weights  $(c_n)_{1 \leq n \leq N}$  such that for every *subharmonic* function  $v$  (i.e.  $-\Delta v \leq 0$ ) :

$$\int_{\Omega} v \geq \sum_{n=1}^N c_n v(z_n).$$

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In the class of **SHQD**, **existence and uniqueness** are ensured ([GUSTAFSSON, 1990]).



## Numerical construction of quadrature domains

The method proposed here is strongly related to **partial balayage of measures** [GUSTAFSSON, SAKAI, 1994] and **obstacle problem** [GUSTAFSSON, SHAHGHOLIAN, 1995].

We assume that the weights  $c_1, \dots, c_N$  are **positive** and that the nodes  $z_1, \dots, z_N$  are contained in the unit disk  $B$ .

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The (unknown) function

$$u_{\Omega}(x) = -\frac{1}{2\pi} \int_{\Omega} \ln|x-y| \, dy,$$

gives access to the characteristic function of  $\Omega$ , since :

$$\mathbb{1}_{\Omega} = -\Delta u_{\Omega}.$$

# Numerical construction of quadrature domains

It can be proved that  $u_\Omega$  is the unique minimizer of the convex functional

$$\mathcal{E}(v) := \frac{1}{2} \int_B |\nabla v|^2 - \int_B v,$$

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$$K := \{v \in H^1(B) \mid v \leq G_N \text{ dans } B; v = G_N \text{ sur } \partial B\},$$

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## Numerical construction of quadrature domains

To obtain the **SHQD**  $\Omega$ , it suffices to compute  $u_\Omega$  by solving the above (standard) minimization problem \*, and use the fact that

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\*. Computations are made with FreeFem++ and we use  $\mathbb{P}_2$  finite elements.

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### Remark

Since  $G_N \notin H^1(B)$ , we will use a slightly modified (regularized) version of this result, that shows that  $\mathbb{1}_\Omega = -\Delta \tilde{u}_\Omega$  where  $\tilde{u}_\Omega$  is the unique minimizer of the convex functional  $\mathcal{E}(v)$  over the closed convex set

$$\tilde{K} := \{v \in H^1(B) \mid v \leq \tilde{G}_N \text{ dans } B; v = \tilde{G}_N \text{ sur } \partial B\},$$

with  $\tilde{G}_N(x) := -\frac{1}{2\pi} \sum_{n=1}^N \int_{B(z_n, r_n)} \ln|x-y| \, dy, r_n := \sqrt{c_n/\pi}.$

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## Algorithm (for positive weights)

- 1 Choose  $N \geq 1$ .
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- 1 Choose  $N \geq 1$ .
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- 3 **Step 1** : Using Prony's method, compute  $(z_n)$  and  $(c_n)$  s.t.

$$\sum_{n=1}^N c_n z_n^m = \int_{\omega} z^m, \quad \forall 0 \leq m \leq 2N - 1.$$

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- 1 Choose  $N \geq 1$ .
- 2 For  $0 \leq m \leq 2N - 1$ , compute  $\int_{\omega} z^m$  from the measurements.
- 3 **Step 1** : Using Prony's method, compute  $(z_n)$  and  $(c_n)$  s.t.

$$\sum_{n=1}^N c_n z_n^m = \int_{\omega} z^m, \quad \forall 0 \leq m \leq 2N - 1.$$

- 4 **Step 2** : Determine the unique **SHQD**  $\omega_N$  associated to  $(z_n)$  and  $(c_n)$ .

# Convergence

## Theorem

Assume that all the weights computed by the algorithm are positive.

- ① If  $\omega$  is a **SHQD** associated to a finite number of points, then there exists  $N$  such that  $\omega_N = \omega$ .
- ② If there exists a compact set  $K \subset B$  and a constant  $C > 0$  such that  $\omega_N \subset K$  for all  $N \geq 1$  :, then

$$\|\nabla u_{\omega_N} - \nabla u_{\omega}\|_{L^{\infty}(\Gamma)} \xrightarrow{N \rightarrow \infty} 0.$$

If, in addition,  $\omega_N$  and  $\omega$  are star-shaped with respect to their centers of gravity and belong to  $\mathcal{U}$ , then  $\omega_N$  converges to  $\omega$ , in the sense that

$$\|\rho_{\omega_N} - \rho_{\omega}\|_{L^{\infty}(0,2\pi)} \longrightarrow 0.$$

# Convergence

The equality determining the quadrature formula

$$\int_{\omega} z^m = \sum_{n=1}^N c_n z_n^m, \quad \forall 0 \leq m \leq 2N - 1$$

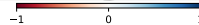
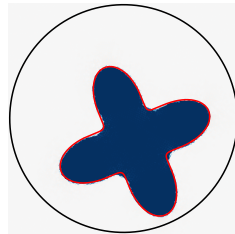
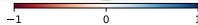
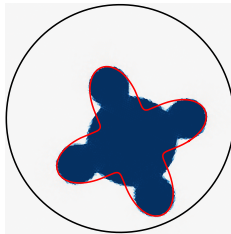
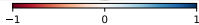
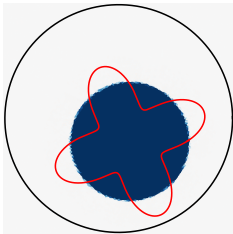
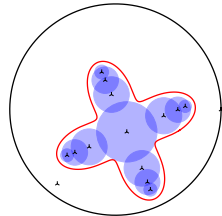
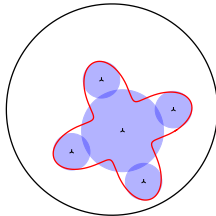
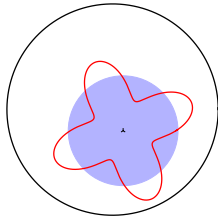
can also be written

$$\int z^m \mathbb{1}_{\omega} = \int z^m d\mu_N, \quad \mu_N := \sum_{n=1}^N c_n \delta_{z_n}.$$

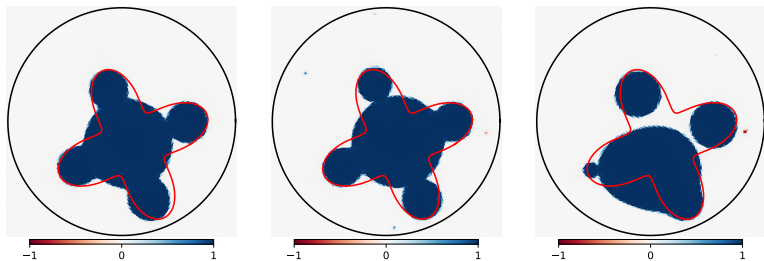
It is thus natural to expect the weights to be positive. In most of our numerical examples, the weights were real and positive. When negative weights appear, one needs to use another algorithm to construct **HQD**; see [AMEUR, HELMER AND TELLANDER, 2021] and [GERBER-ROTH, MUNNIER, RAMDANI, 2023].

# OUTLINE

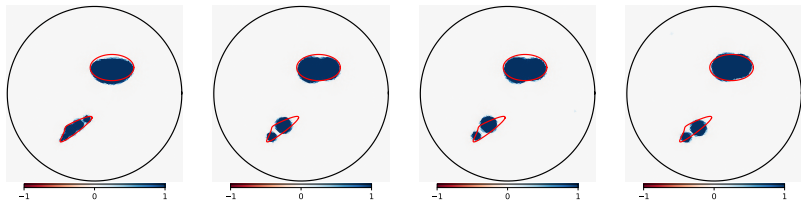
- 1 Problem formulation
- 2 Reconstruction method
  - Step 1 : Constructing quadrature formula
  - Step 2 : Constructing quadrature domains
- 3 Numerical results



Reconstructing a star-shaped domain without noise :  $N = 1, 5, 15$ .

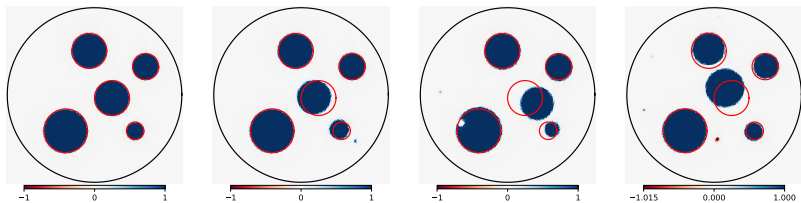


Reconstructing a star-shaped domain with noisy data ( $N = 15$ ) :  
1%, 3%, 5%.

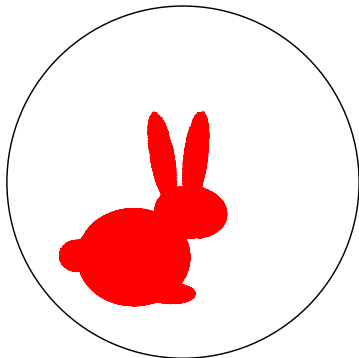


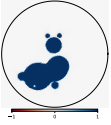
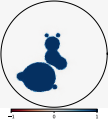
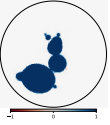
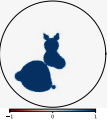
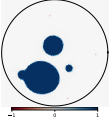
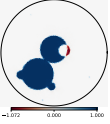
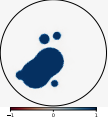
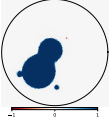
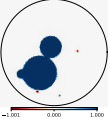
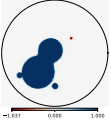
Reconstructing a domain convex in one direction with noisy data  
( $N = 12$ ) : 0%, 1%, 2%, 3%.

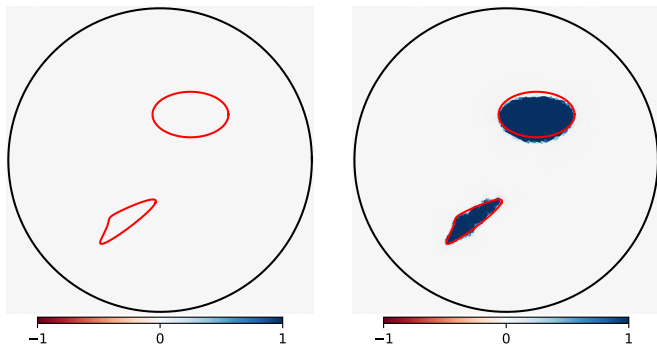




Reconstruction of disks with noisy data ( $N = 10$ ) : 0%, 5%, 8%,  
10%.



	$N = 10$	$N = 15$	$N = 20$	$N = 25$
0 %				
1 %				
3 %				
5 %				



Instabilities may occur : false reconstruction with  $N = 31$  (all the weights are almost zero) and correct one for  $N = 32$ .

# Conclusion

- We presented a reconstruction method based on an original combination of Prony's method and quadrature domains :  
*A reconstruction method for the inverse gravimetric problem*  
Anthony Gerber-Roth, Alexandre Munnier and Karim Ramdani  
SMAI Journal of Computational Mathematics (2023)  
<https://smaj-jcm.centre-mersenne.org/articles/10.5802/smai-jcm.99/>
- What's next ?
  - Surface uniform mass distribution.
  - Non uniform mass distribution.
  - 3D
  - Stokes problem ([Alexandre Munnier](#)).

## Selected bibliography

- *"Quadrature domains"* Makoto Sakai (1982)
- *"Inverse source problems"* Victor Isakov (1990)
- *"Properties of some balayage operators, with applications to quadrature domains and moving boundary problems"* Björn Gustafsson and Makoto Sakai (1994)
- *"A stable numerical method for inverting shape from moments"* Gene H. Golub, Peyman Milanfar and James Varah (1999)
- *"Lectures on balayage"* Björn Gustafsson (2004)
- *"What is a quadrature domain?"* Björn Gustafsson and Harold S. Shapiro (2005)
- *"On the accuracy of solving confluent Prony systems"* Dmitry Batenkov and Yosef Yomdin (2013)

Interpolation géométrique et lien avec le cone huygens (au pres <sup>1</sup> ~~pres~~)  
 on a toujours  

$$T_{\mathbb{Z}}K \subset \text{Ker}(f'(\bar{z}))$$
  
 En effet,  $d \in T_{\mathbb{Z}}K \Rightarrow \exists x_k \in K, h_k \in \mathbb{R}^n, \frac{x_k - \bar{z}}{h_k} \rightarrow d$   

$$0 = f'(x_k) = f'(\bar{z})(x_k - \bar{z}) + o(\|x_k - \bar{z}\|)$$
  

$$0 = f'(\bar{z})\left(\frac{x_k - \bar{z}}{h_k}\right) + \frac{1}{h_k} o(\|x_k - \bar{z}\|) \xrightarrow{h_k \rightarrow 0} f'(\bar{z})d$$

4.1. Approche naïve: les idées, sans trop de technique...

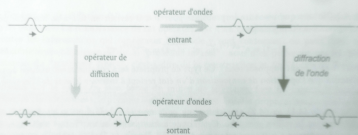


Figure 4.2: Opérateurs d'ondes et opérateurs de diffusion

(où la seconde égalité découle du caractère unitaire<sup>8</sup> de  $\exp(-iA^{1/2}t)$ ). Ainsi les états initiaux  $u(t)$  et  $\hat{u}(t)$  correspondent à une même onde incidente (ou entrante) si cette quantité tend vers  $-\infty$ . De même, ils correspondent à une même onde sortante si elle tend vers  $+\infty$ . C'est pourquoi la diffraction entrante

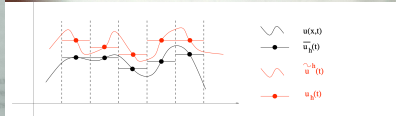
Soit  $v \in D(A^*)$ . Alors il existe  $w \in L^2$   

$$\forall u \in D(A) \int_{\mathbb{R}^2} d\omega (E\nabla u) v = \int_{\mathbb{R}^2} u w$$
  
 Soit  $\chi$  une fonction de troncature de la forme  

$$\chi(x) = \Theta(|x - x_0|)$$
  
 où  $\Theta$  est une fonction de troncature et  $x_0$  un point.  
 On a  $\forall u \in D(A) \chi u \in D(A)$  donc  

$$\int_{\mathbb{R}^2} d\omega (E\nabla(\chi u)) v = \int_{\mathbb{R}^2} \chi u w$$
  
 d'où  

$$\int_{\mathbb{R}^2} d\omega (E\nabla u)(\chi v) = \int_{\mathbb{R}^2} u(\chi w - \nabla u \cdot \nabla \chi v)$$



...MERCI....

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## Further comments

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# Solving Prony's system

$$P(Z) := \prod_{n=1}^N (Z - z_n) = Z^N + \sum_{n=0}^{N-1} \alpha_n Z^n.$$

$$\left\{ \begin{array}{l} c_1 + c_2 + \dots + c_N = \tau_0 \\ c_1 z_1 + c_2 z_2 + \dots + c_N z_N = \tau_1 \\ c_1 z_1^2 + c_2 z_2^2 + \dots + c_N z_N^2 = \tau_2 \\ \vdots \\ c_1 z_1^{N-1} + c_2 z_2^{N-1} + \dots + c_N z_N^{N-1} = \tau_{N-1} \\ c_1 z_1^N + c_2 z_2^N + \dots + c_N z_N^N = \tau_N \\ \vdots \\ c_1 z_1^{2N-1} + c_2 z_2^{2N-1} + \dots + c_N z_N^{2N-1} = \tau_{2N-1} \end{array} \right.$$

# Solving Prony's system

$$P(Z) := \prod_{n=1}^N (Z - z_n) = Z^N + \sum_{n=0}^{N-1} \alpha_n Z^n.$$

$$\left\{ \begin{array}{l} c_1 + c_2 + \dots + c_N = \tau_0 \quad \times \alpha_0 \\ c_1 z_1 + c_2 z_2 + \dots + c_N z_N = \tau_1 \quad \times \alpha_1 \\ c_1 z_1^2 + c_2 z_2^2 + \dots + c_N z_N^2 = \tau_2 \quad \times \alpha_2 \\ \vdots \\ c_1 z_1^{n-1} + c_2 z_2^{n-1} + \dots + c_N z_N^{n-1} = \tau_{n-1} \quad \times \alpha_{n-1} \\ c_1 z_1^n + c_2 z_2^n + \dots + c_N z_N^n = \tau_n \quad \times 1 \\ \vdots \\ c_1 z_1^{2N-1} + c_2 z_2^{2N-1} + \dots + c_N z_N^{2N-1} = \tau_{2N-1} \end{array} \right.$$

$$c_1 P(z_1) + c_2 P(z_2) + \dots + c_N P(z_N) =$$

$$\alpha_0 \tau_0 + \alpha_1 \tau_1 + \dots + \alpha_{N-1} \tau_{N-1} + \tau_N$$

$$\alpha_0 \tau_0 + \alpha_1 \tau_1 + \dots + \alpha_{n-1} \tau_{n-1} = -\tau_n.$$

# Solving Prony's system

$$P(Z) := \prod_{n=1}^N (Z - z_n) = Z^N + \sum_{n=0}^{N-1} \alpha_n Z^n.$$

$$\left\{ \begin{array}{l} c_1 + c_2 + \dots + c_N = \tau_0 \\ c_1 z_1 + c_2 z_2 + \dots + c_N z_N = \tau_1 \quad \times \alpha_0 \\ c_1 z_1^2 + c_2 z_2^2 + \dots + c_N z_N^2 = \tau_2 \quad \times \alpha_1 \\ \vdots \\ c_1 z_1^{N-1} + c_2 z_2^{N-1} + \dots + c_N z_N^{N-1} = \tau_{N-1} \quad \times \alpha_{N-2} \\ c_1 z_1^N + c_2 z_2^N + \dots + c_N z_N^N = \tau_N \quad \times \alpha_{N-1} \\ c_1 z_1^{N+1} + c_2 z_2^{N+1} + \dots + c_N z_N^{N+1} = \tau_{N+1} \quad \times 1 \\ \vdots \\ c_1 z_1^{2N-1} + c_2 z_2^{2N-1} + \dots + c_N z_N^{2N-1} = \tau_{2N-1} \end{array} \right.$$

$$c_1 z_1 P(z_1) + \dots + c_n z_n P(z_n) = \alpha_0 \tau_1 + \alpha_1 \tau_2 + \dots + \dots + \alpha_{n-1} \tau_n + \tau_{n+1}$$

$$\alpha_0 \tau_1 + \alpha_1 \tau_2 + \dots + \dots + \alpha_{n-1} \tau_n = -\tau_{n+1}.$$

# Solving Prony's system

$$\begin{pmatrix} \tau_0 & \tau_1 & \cdots & \tau_{n-1} \\ \tau_1 & \tau_2 & \cdots & \tau_n \\ \vdots & \vdots & \ddots & \vdots \\ \tau_{n-1} & \tau_n & \cdots & \tau_{2n-2} \end{pmatrix} \begin{pmatrix} \alpha_0 \\ \alpha_1 \\ \vdots \\ \alpha_{n-1} \end{pmatrix} = - \begin{pmatrix} \tau_n \\ \tau_{n+1} \\ \vdots \\ \tau_{2n-1} \end{pmatrix} \iff \mathbb{T}_0 \boldsymbol{\alpha} = -\boldsymbol{\tau}'.$$